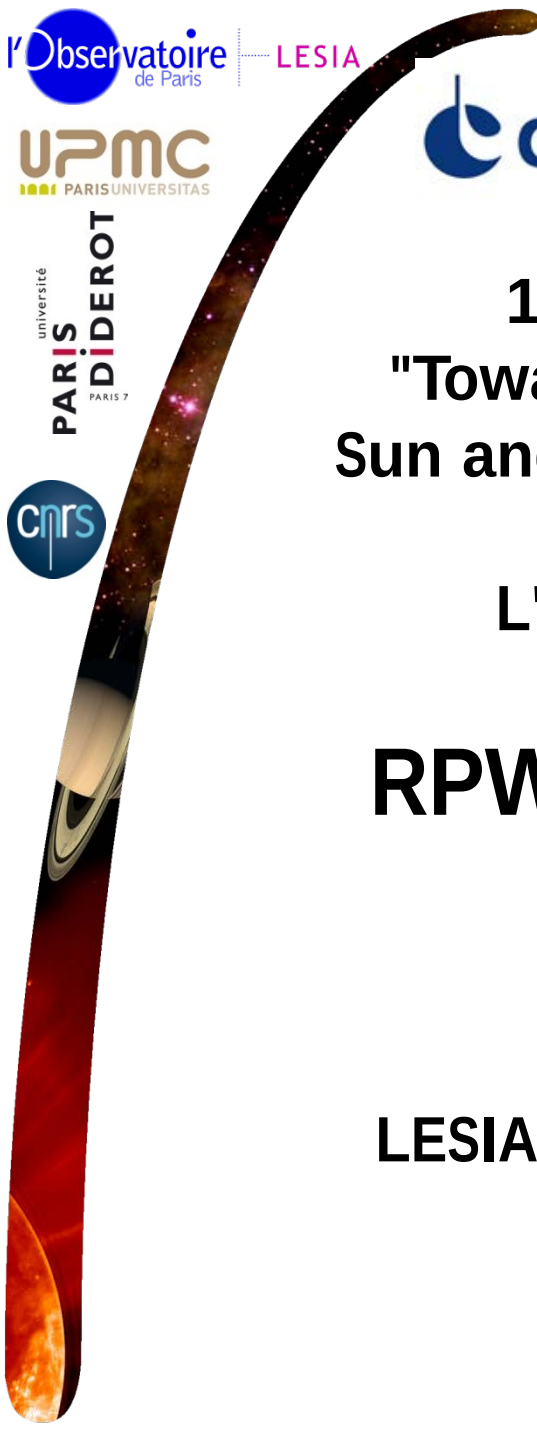


1st Solar Orbiter Summer School "Towards a Deeper Understanding of the Sun and the Heliosphere with Solar Orbiter"

L'Aquila, September 22-25, 2014

RPW: Measuring Solar Radio and Plasma Waves

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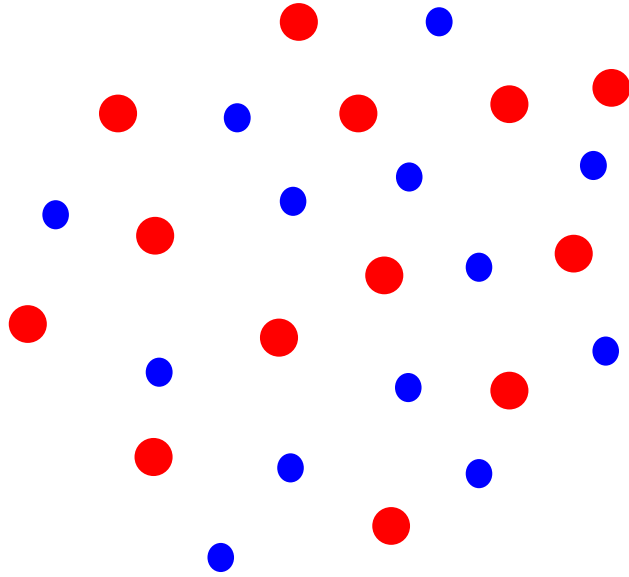
Outline



- ❑ Plasma waves versus Radio waves
- ❑ Plasma waves : what do we observed and understand ?
- ❑ Radio waves : what do we observed and understand ?
- ❑ « *Extra* » science : Dust measurements
- ❑ On the difficulty of measuring the DC/LF electric component of plasma waves
- ❑ RPW : a brief description

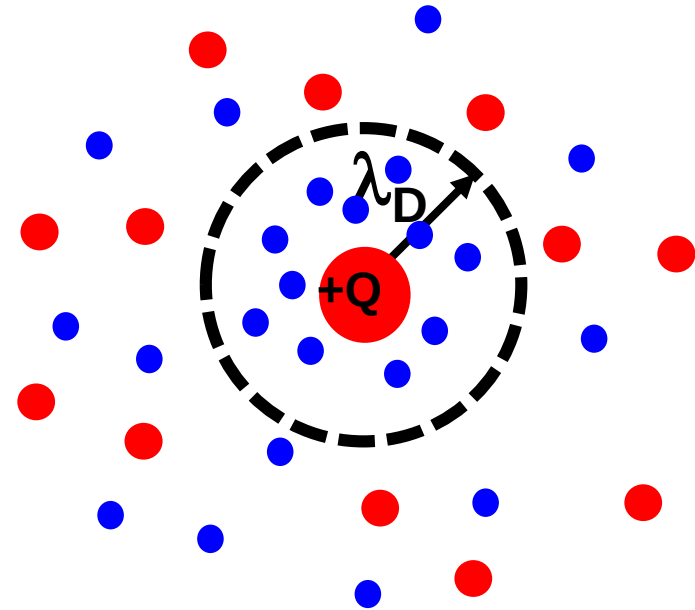
Plasma waves versus radio waves

Debye Screening



$$\lambda_D = \left(\frac{\epsilon_0 k_B T}{n e^2} \right)^{1/2}$$

$$\lambda_D \approx 69 \sqrt{T/n} \quad \text{in S.I. units}$$



$$\Phi(r) = \frac{Q}{4\pi\epsilon_0 r} e^{-r\sqrt{2}/\lambda_D}$$

- $\lambda_D \approx 10 \text{ m @ } 1 \text{ AU}$
- $\lambda_D \approx 5 \text{ m @ } 0.3 \text{ AU}$
- $\lambda_D \approx 1 \text{ m @ } 10 \text{ Rs}$

- v must be compared to both $V_{thi} \approx \sqrt{k_B T / m_i}$ and $V_{the} \approx \sqrt{k_B T / m_e}$
- If $v < V_{thi} < V_{the}$ then both electrons and ions have time « to participate » to the screening of $+Q$
- If $V_{thi} < v < V_{the}$ then only the electrons have time « to participate » to the screening of $+Q$
- If $V_{thi} < V_{the} < v$ then there is no screening of $+Q$

Timescale for screening ?

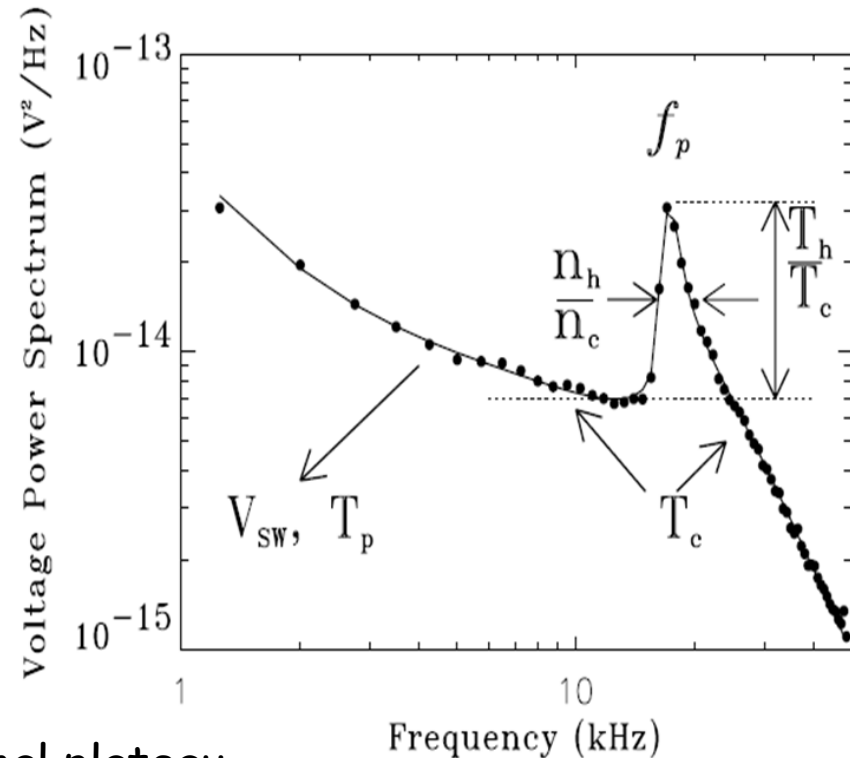
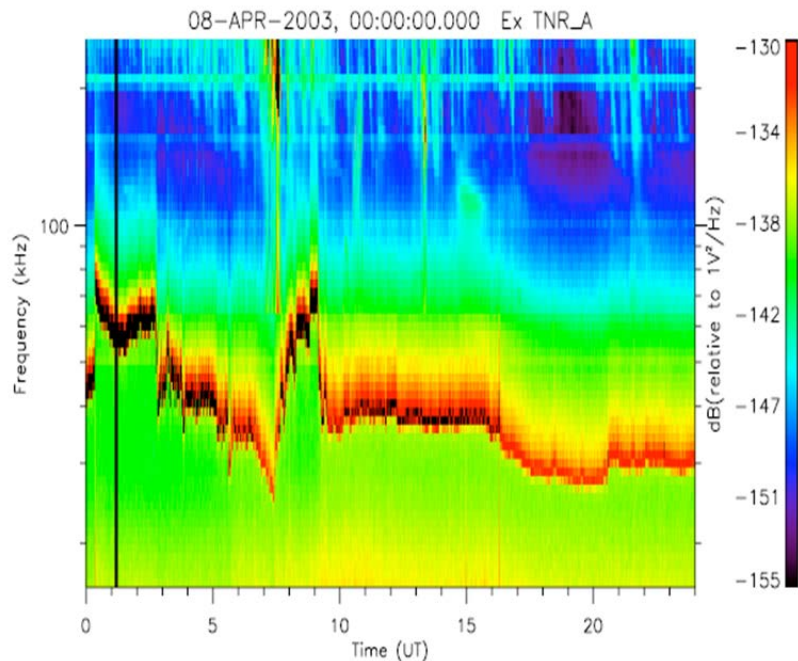
- Let's put suddenly a charge $+Q$ in a plasma at equilibrium
- The electrons, which move faster, will need a time τ to travel a distance λ_D at the most probable speed V_{the}

$$\tau \approx \left(\frac{\epsilon_0 m_e}{n e^2} \right)^{1/2} \equiv \frac{1}{\omega_p}$$

$\omega_p = 2\pi F_p$ is the angular plasma frequency

- $F_p = 9\sqrt{n}$ in S.I. units
- $F_p \approx 20$ kHz @ 1 AU
- $F_p \approx 60$ kHz @ 0.3 AU
- $F_p \approx 0.5$ MHz @ 10 Rs

Electron Density and Temperature from Quasi-Thermal Noise Spectroscopy



Antenna geometry and size matter

- e^- passing closer than $\lambda_D \rightarrow F < F_p \rightarrow$ thermal plateau
- e^- passing further than $\lambda_D \rightarrow$ Langmuir waves @ $F < F_p$
Power $\propto (F_p/F)^3$
- In practice $L/\lambda_D \geq 1$ is needed : satisfied at perihelion and in dense SW
- RPW should measure n_e & T_e with accuracies respectively of a few % and 10 %

2 Maxwellians:

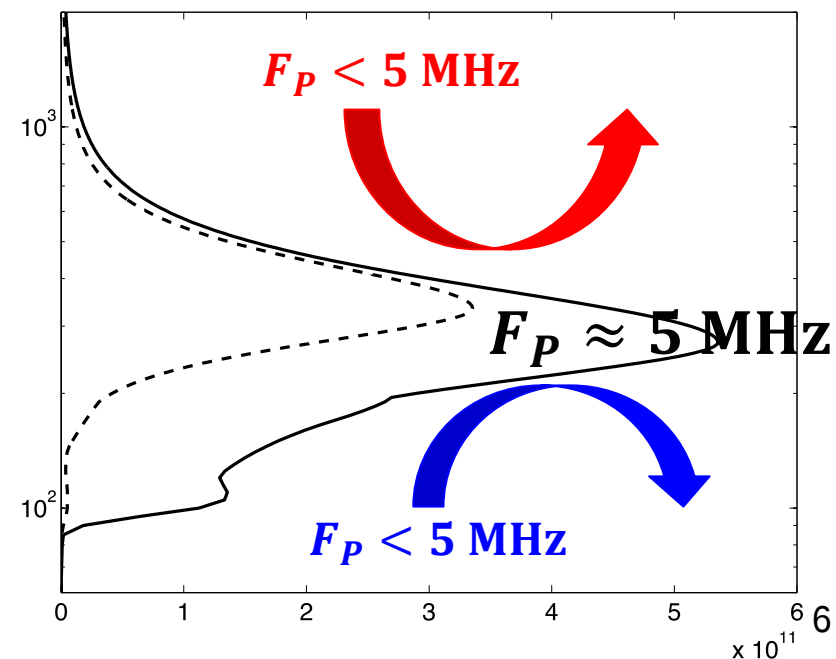
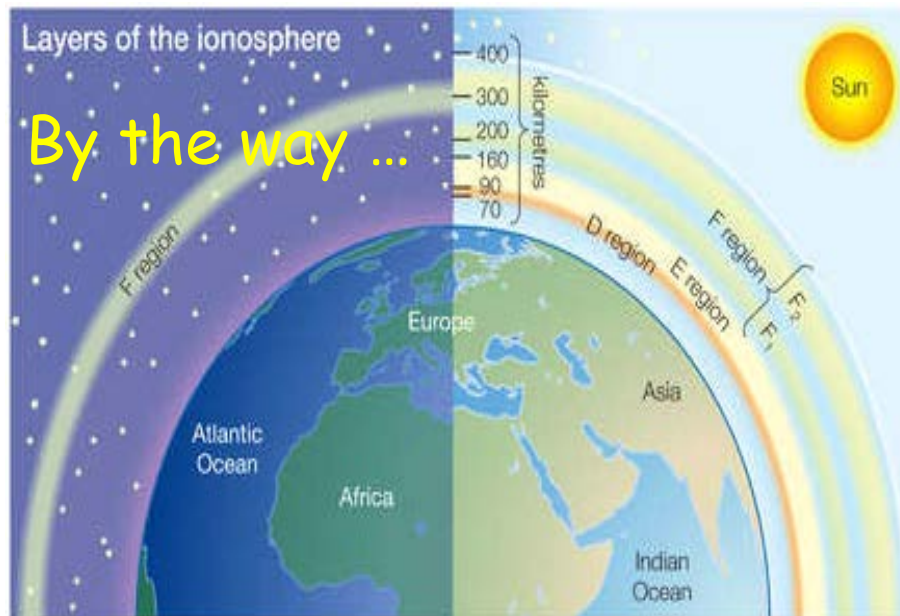
cold : n_c, T_c

hot : n_h, T_h

- Waves with $F < F_{Pi} < F_{Pe}$ are « seen » by both ions and electrons. They are screened and do not propagate
- Waves with $F_{Pi} < F < F_{Pe}$ are not « seen » by ions but are still screened by the electrons
- Waves with $F_{Pi} < F_{Pe} < F$ are not seen by the plasma and propagate freely

Plasma Waves

Radio Waves

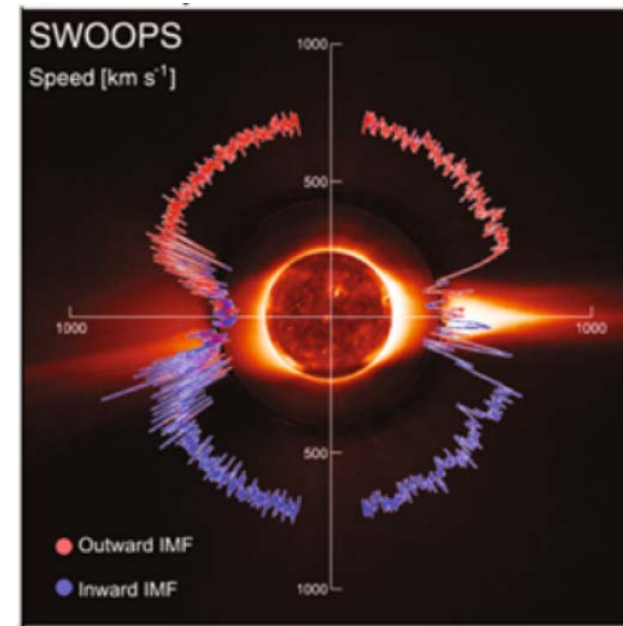
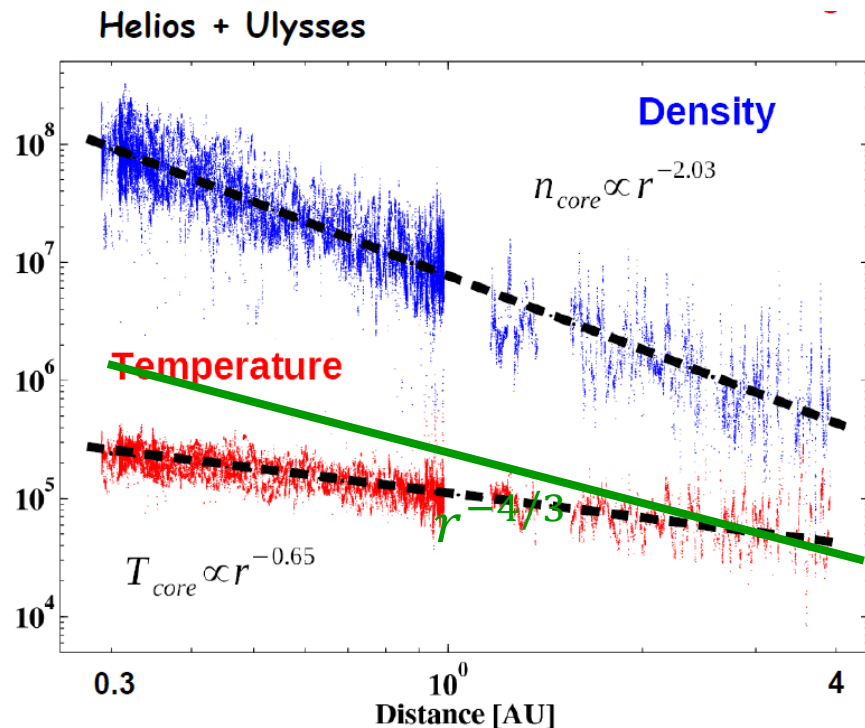


- There is a Magnetic Field Magnetic in the Solar Wind which imposes new time and length scales (gyro-radii and -frequencies)
- Full zoo of waves (see a good Plasma Physics textbooks)

EM character	oscillating species	conditions	dispersion relation	name
electrostatic	electrons	$\vec{B}_0 = 0$ or $\vec{k} \parallel \vec{B}_0$	$\omega^2 = \omega_p^2 + 3k^2 v_{th}^2$	plasma oscillation (or Langmuir wave)
		$\vec{k} \perp \vec{B}_0$	$\omega^2 = \omega_p^2 + \omega_c^2 = \omega_h^2$	upper hybrid oscillation
	ions	$\vec{B}_0 = 0$ or $\vec{k} \parallel \vec{B}_0$	$\omega^2 = k^2 v_s^2 = k^2 \frac{\gamma_e K T_e + \gamma_i K T_i}{M}$	ion acoustic wave
		$\vec{k} \perp \vec{B}_0$ (nearly)	$\omega^2 = \Omega_c^2 + k^2 v_s^2$	electrostatic ion cyclotron wave
		$\vec{k} \perp \vec{B}_0$ (exactly)	$\omega^2 = [(\Omega_c \omega_c)^{-1} + \omega_i^{-2}]^{-1}$	lower hybrid oscillation
electromagnetic	electrons	$\vec{B}_0 = 0$	$\omega^2 = \omega_p^2 + k^2 c^2$	light wave
		$\vec{k} \perp \vec{B}_0, \vec{E}_1 \parallel \vec{B}_0$	$\frac{c^2 k^2}{\omega^2} = 1 - \frac{\omega_p^2}{\omega^2}$	O wave
		$\vec{k} \perp \vec{B}_0, \vec{E}_1 \perp \vec{B}_0$	$\frac{c^2 k^2}{\omega^2} = 1 - \frac{\omega_p^2}{\omega^2} \frac{\omega^2 - \omega_p^2}{\omega^2 - \omega_h^2}$	X wave
		$\vec{k} \parallel \vec{B}_0$ (right circ. pol.)	$\frac{c^2 k^2}{\omega^2} = 1 - \frac{\omega_p^2 / \omega^2}{1 - (\omega_c / \omega)}$	R wave (whistler mode)
		$\vec{k} \parallel \vec{B}_0$ (left circ. pol.)	$\frac{c^2 k^2}{\omega^2} = 1 - \frac{\omega_p^2 / \omega^2}{1 + (\omega_c / \omega)}$	L wave
	ions	$\vec{B}_0 = 0$		none
		$\vec{k} \parallel \vec{B}_0$	$\omega^2 = k^2 v_A^2$	Alfvén wave
		$\vec{k} \perp \vec{B}_0$	$\frac{\omega^2}{k^2} = c^2 \frac{v_s^2 + v_A^2}{c^2 + v_A^2}$	magnetosonic wave

In-situ Plasma waves : what do we observed and understand ?

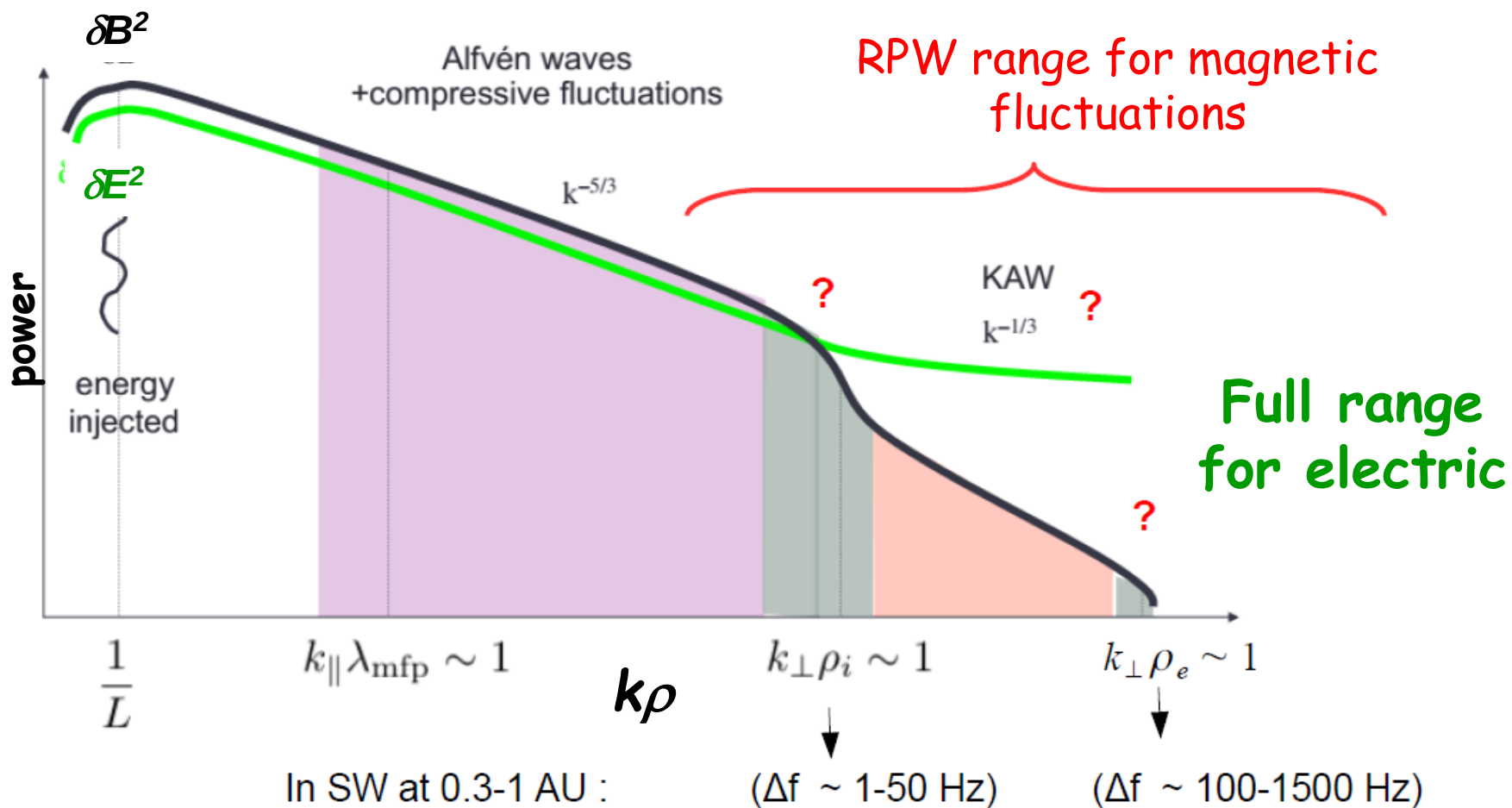
Waves are believed to both heat & accelerate the Solar Wind



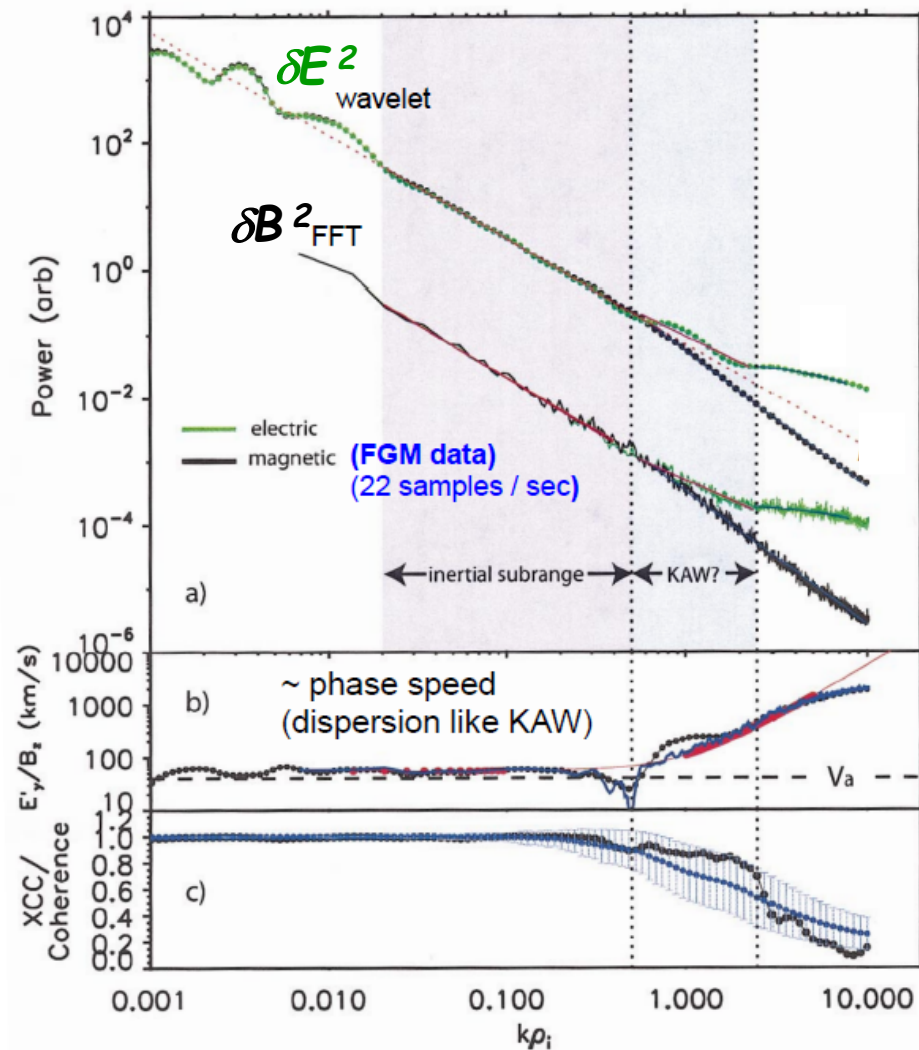
Turbulent heating/dissipation in the Solar Wind

Kinetic Alfvén Waves ? Whistler Waves ?
Coherent structures ? Other ?

Turbulent energy cascade from large to small scales



Evidences for Kinetic Alven Waves turbulence ?



Cluster

- Phase speed is $\sim$ Alfven speed
- Dispersion at short wavelength is like Kinetic Alfvén Wave
- Strong evidence/consistency with KAW turbulence

Evidence for a turbulent dissipation at the proton gyroscale?

What above $k\rho_p \sim 1$?

Cascade up to the electron scale?

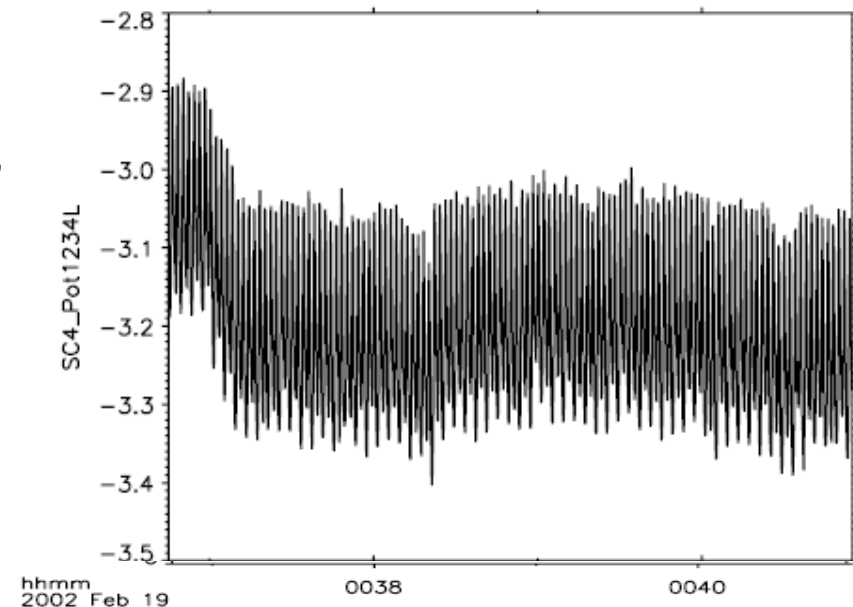
Bale et al., PRL, 2005
(also Sahraoui et al., PRL, 2009)

Caveats

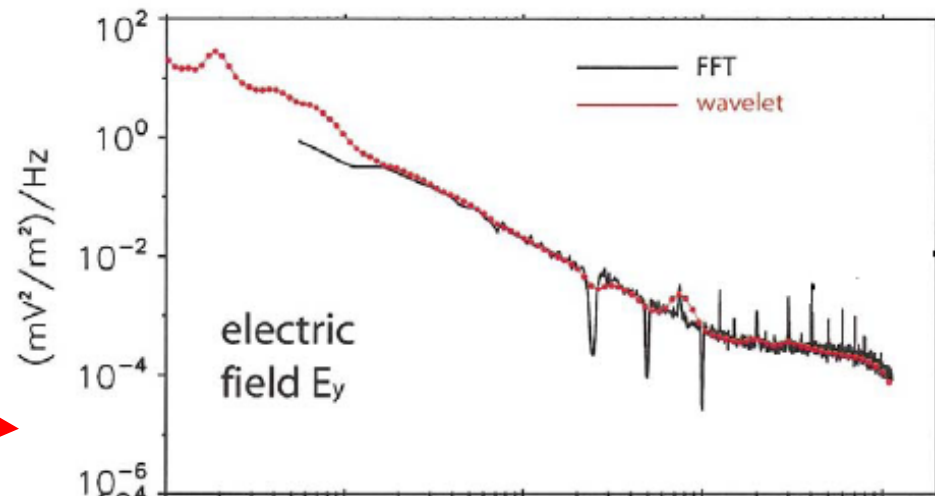
On Cluster electric measurements not easy because of S/C charging

Spin tone in the raw E data due to :

- illumination variations on the probes
- wake effects



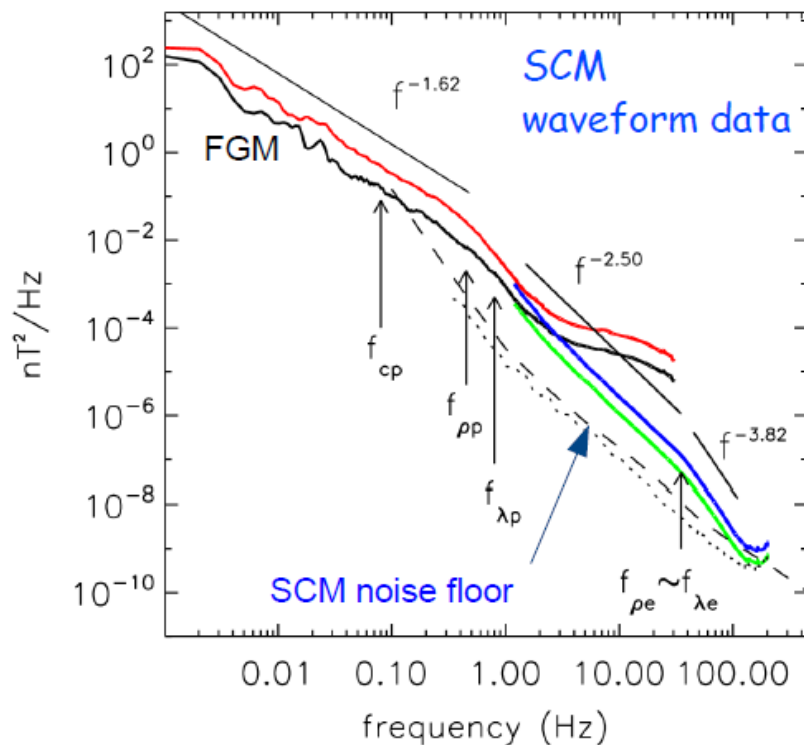
Spin tone needs to be filtered



Evidences for a Cascade well above the proton scales up to the electron typical scales (gyroscale and above ?)

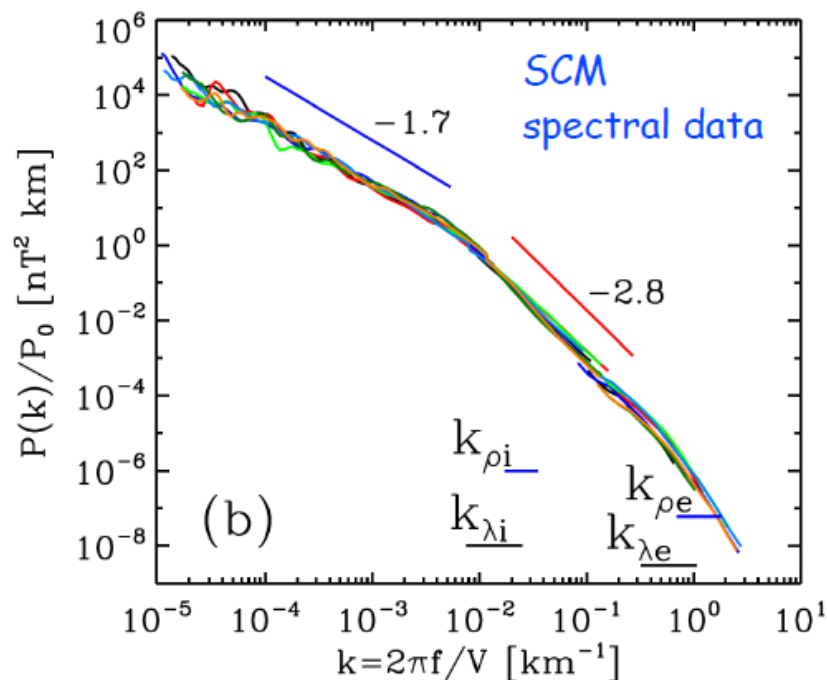
A first case study

Cluster



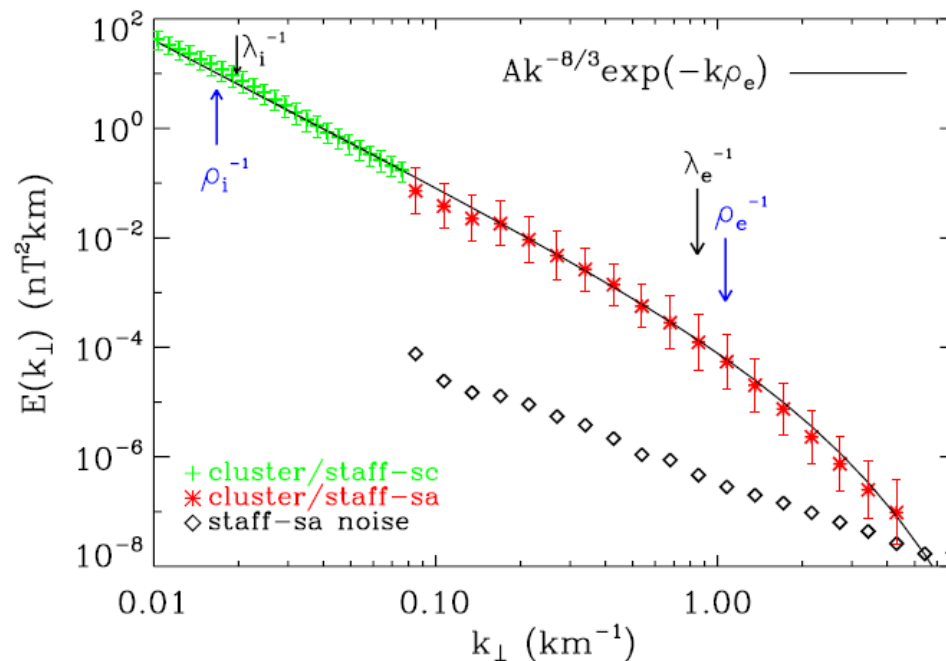
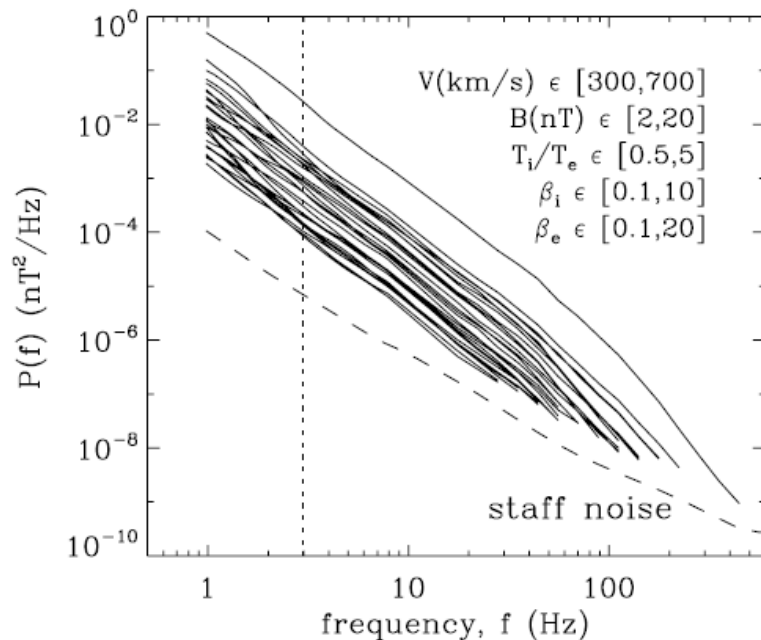
Sahraoui et al., PRL, 2009

Another similar study (several cases)

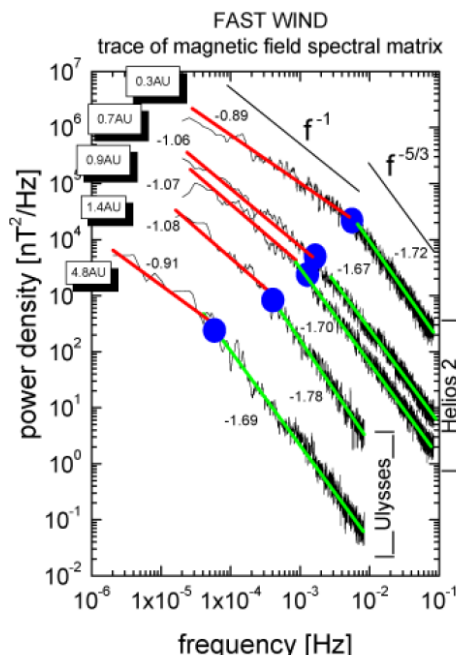


Alexandrova et al., PRL, 2009

- Exponential decrease or power law ?
- Observational limitations : noise floor both for dE2 and dB2

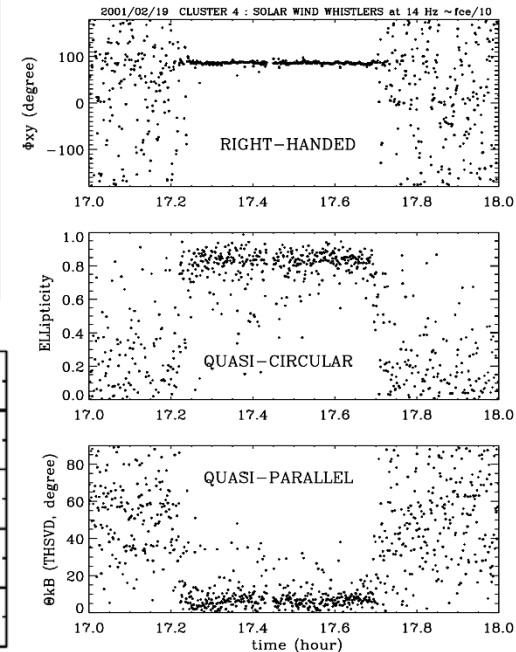
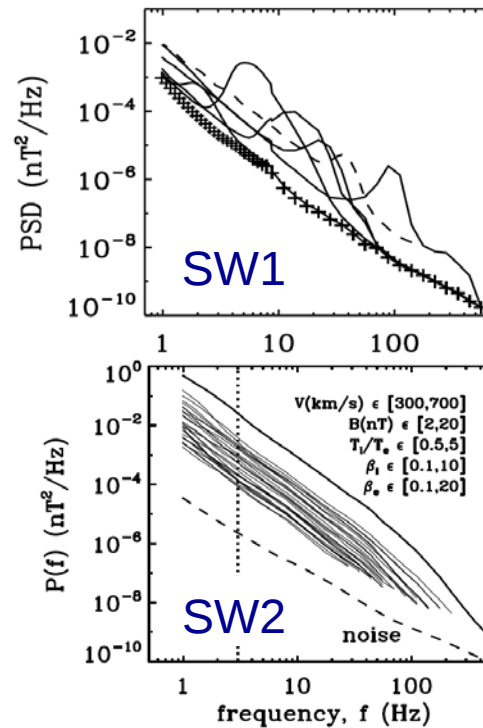
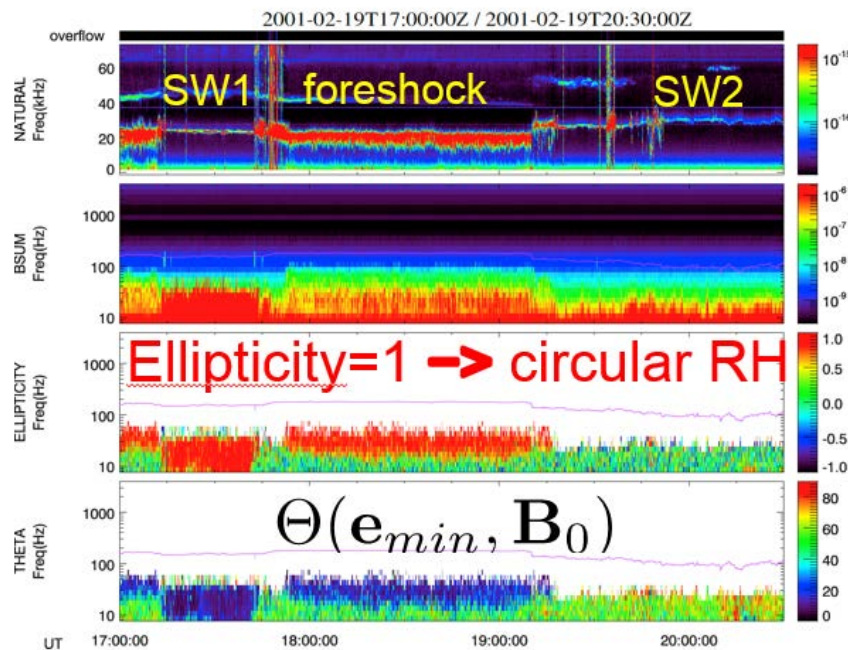


Even with the same search coil sensitivity as on Cluster, RPW will make important improvements

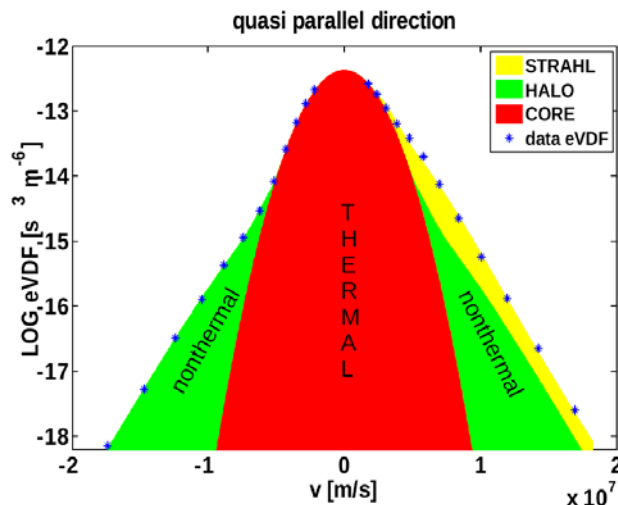


- A function of the ion thermal pressure
- Similar to the behaviour in the « far » dissipation range in usual fluid turbulence
 $E(k) \propto k^3 e^{-ckl_d} \text{ ?}$

Coherent waves are also observed for e.g Whistler waves (Lacombe et al., 2014)

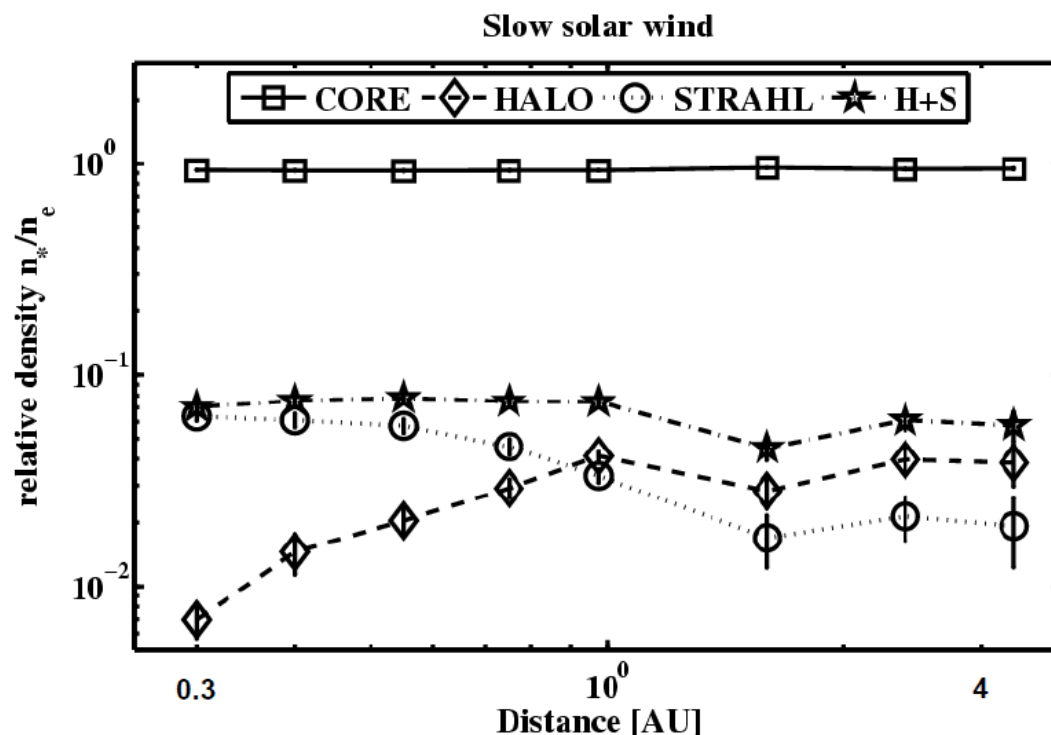


- Polarized fluctuations => spectra with bumps (10% of data)
- Non-polarized fluctuations => permanent (or background) turbulence (90% of data, Alexandrova et al. 2012, 2013)
- Permanent turbulence + sporadic polarized fluctuations => "intermediate" spectral shape (breaks, small bumps, ...)



Can these Whistler waves explain the radial evolution of electron VDFs ?

Helios, Cluster, Ulysses



Data set : ~ 240,000 samples
from 0.3 to 4 AU

Stverak et al., JGR, 2009

Maksimovic et al., 2005

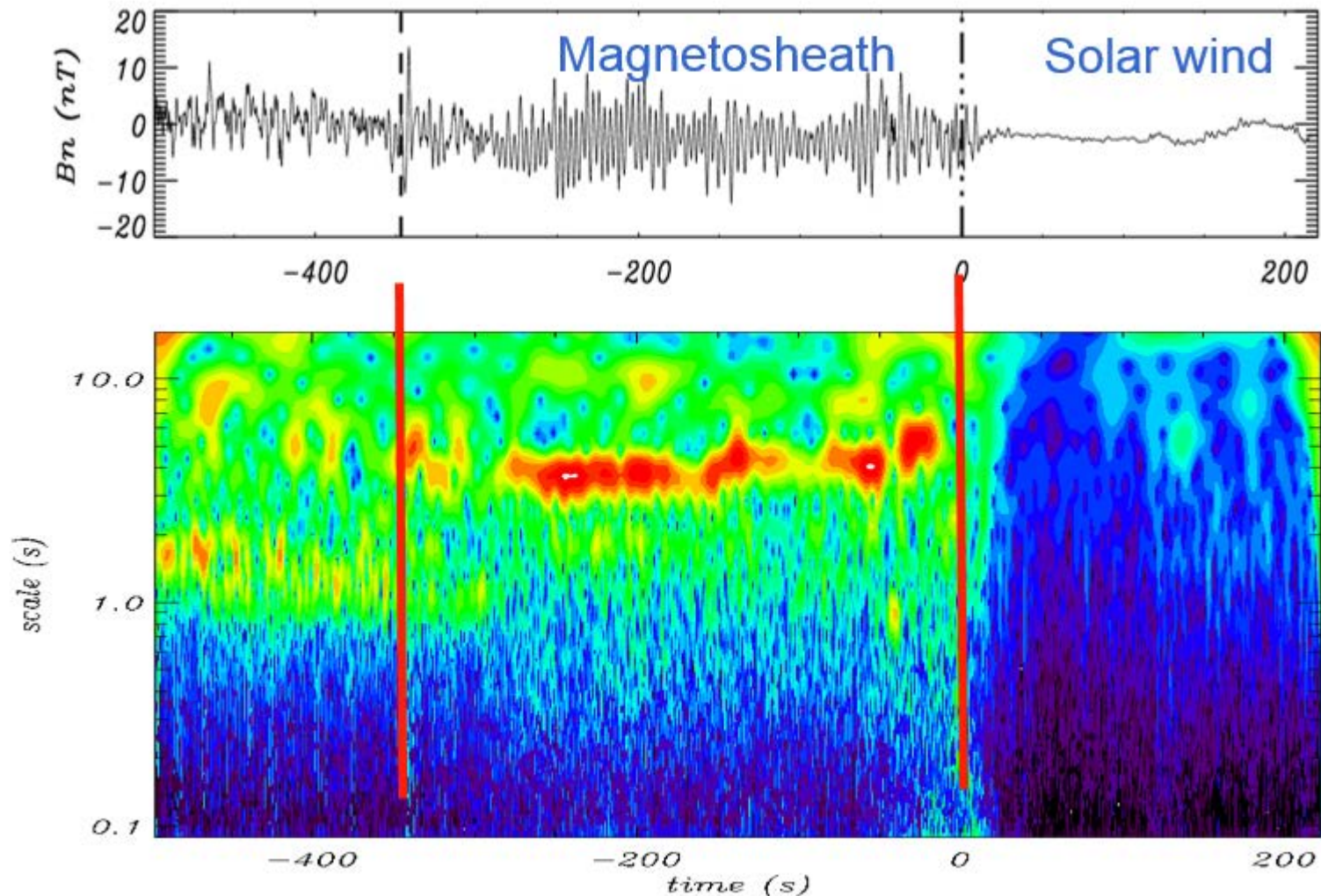
During expansion,
strahl is transformed
into halo by wave-
particle interactions ?

And even solitary structures ...

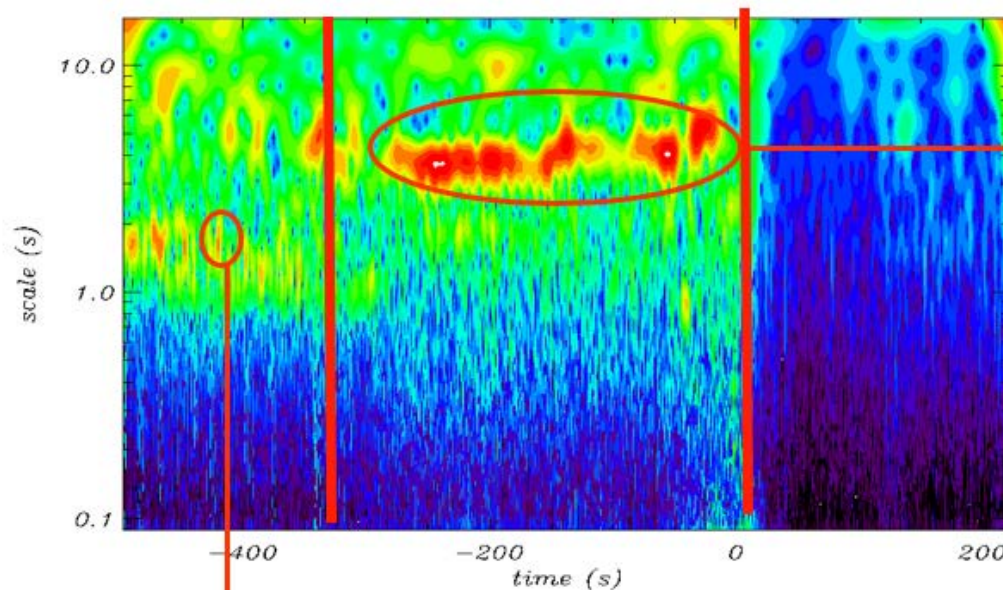
For e.g. Alexandrova et al. 2006, JGR



Wavelet analysis of CLUSTER SCM waveforms



Presence of time-localized events in the vicinity of Alfvén Ion Cyclotron wave



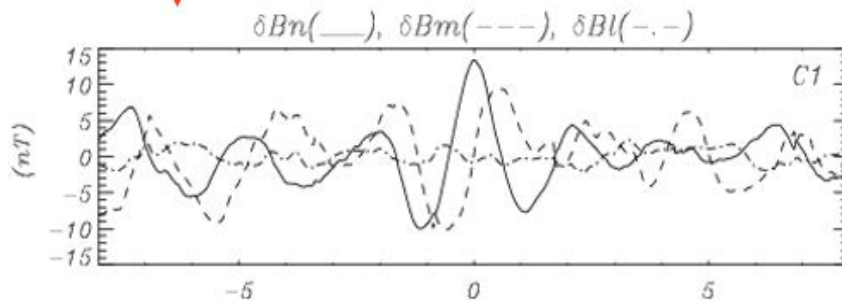
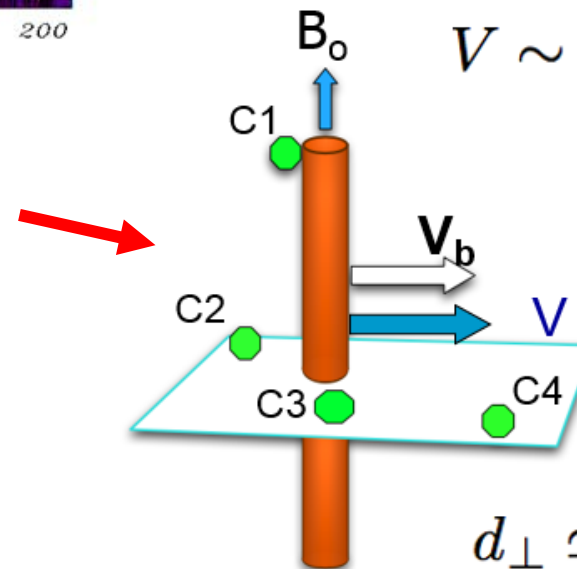
Alfvén Ion
Cyclotron wave

Alexandrova
et al.
2006, JGR

$$\lambda \propto 10 c / \omega_{pi}$$

$$f \propto 0.4 f_{ci}$$

In plasma frame:
 $V \sim [0, 0.3] V_A$



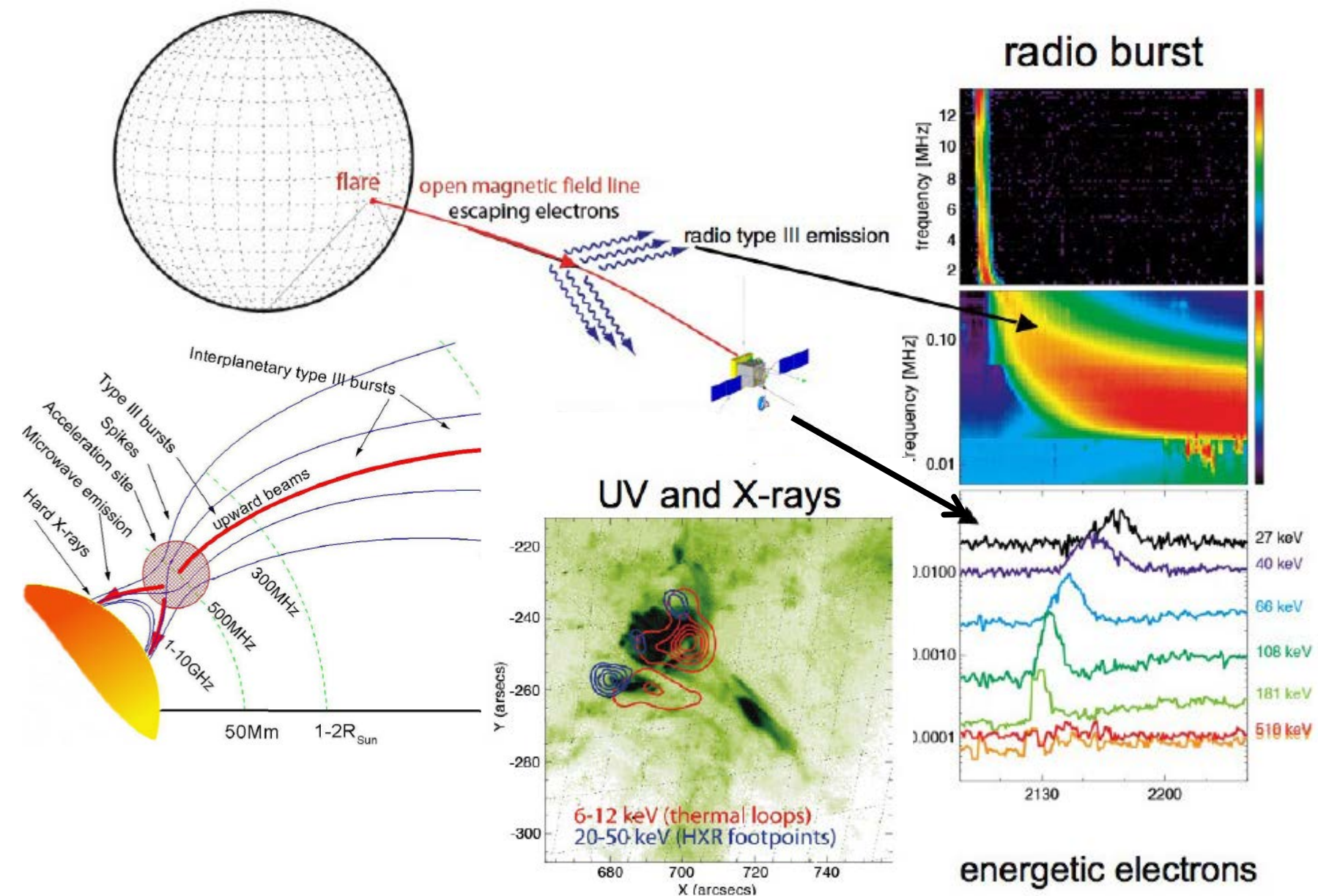
Full characterization with the 4
CLUSTER S/C of localized Alfvén
Vortices : role for the dissipation ?

$$d_{\perp} \simeq 500 km$$

$$d_{\perp} \sim 10 c / \omega_{pi}$$

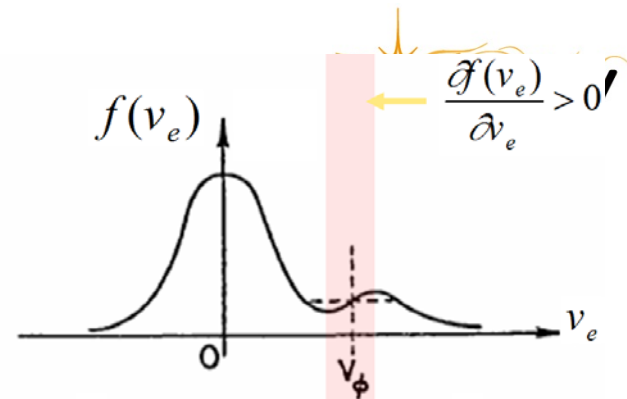
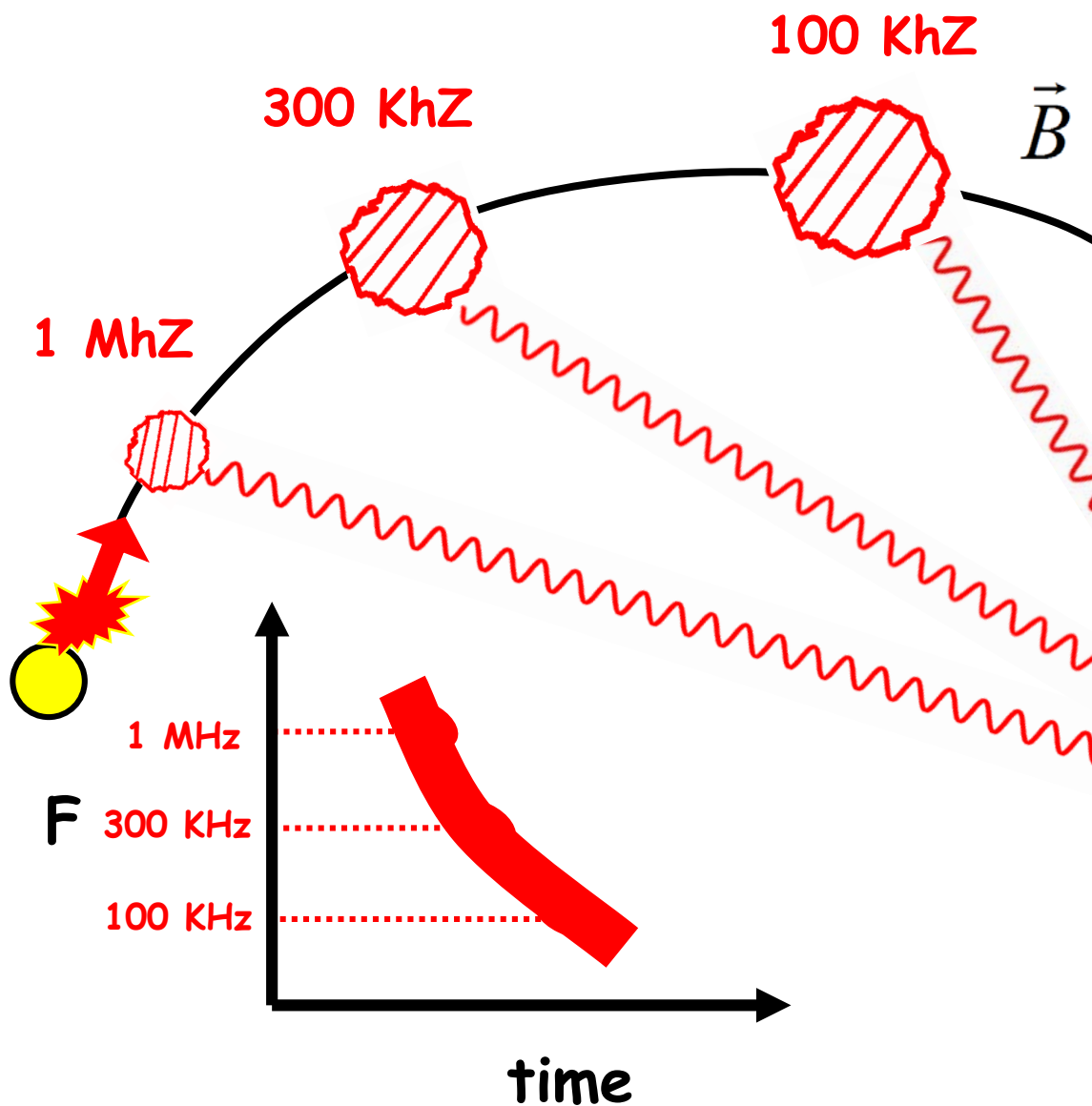
Radio waves :
what do we observed and
understand ?

Solar Radio Bursts



energetic electrons

Solar Type III radio emissions



Electrostatic
Langmuir waves
→ radio emission

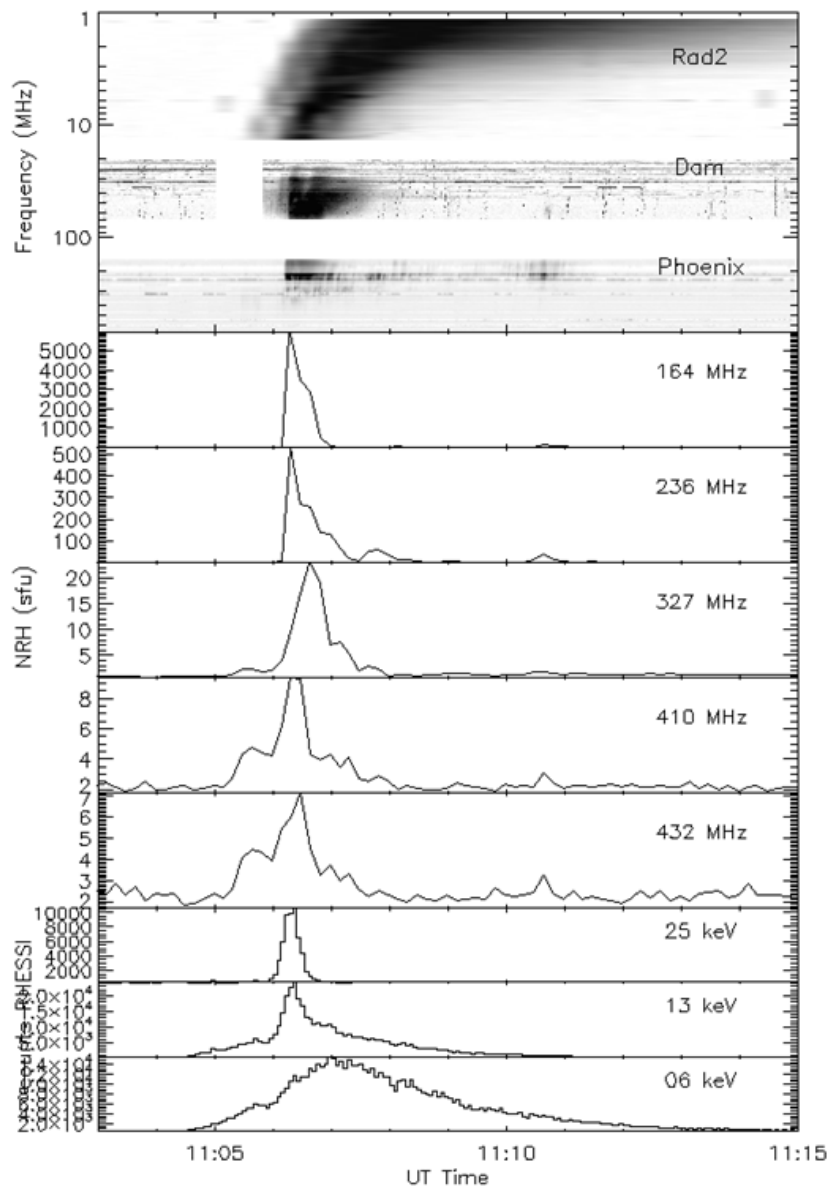
$$\left. \begin{aligned} F_p (\text{kHz}) &\propto \sqrt{N_e (\text{cm}^{-3})} \\ N_e &\propto 1/R^2 (\text{au}) \end{aligned} \right\}$$

$$\rightarrow F_p \propto \frac{1}{R}$$

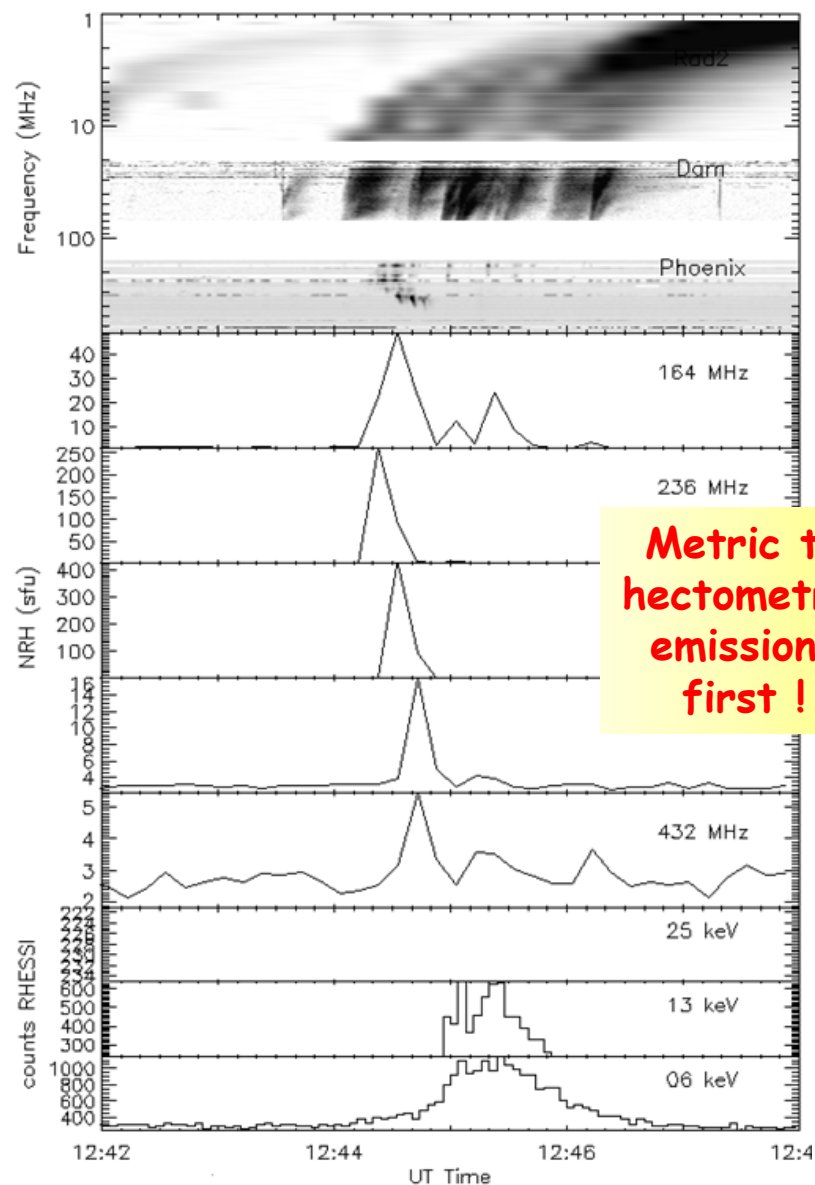
Combined RHESSI / NRH / Phoenix/Dam/WIND-WAVES observations



20 FEB 2002

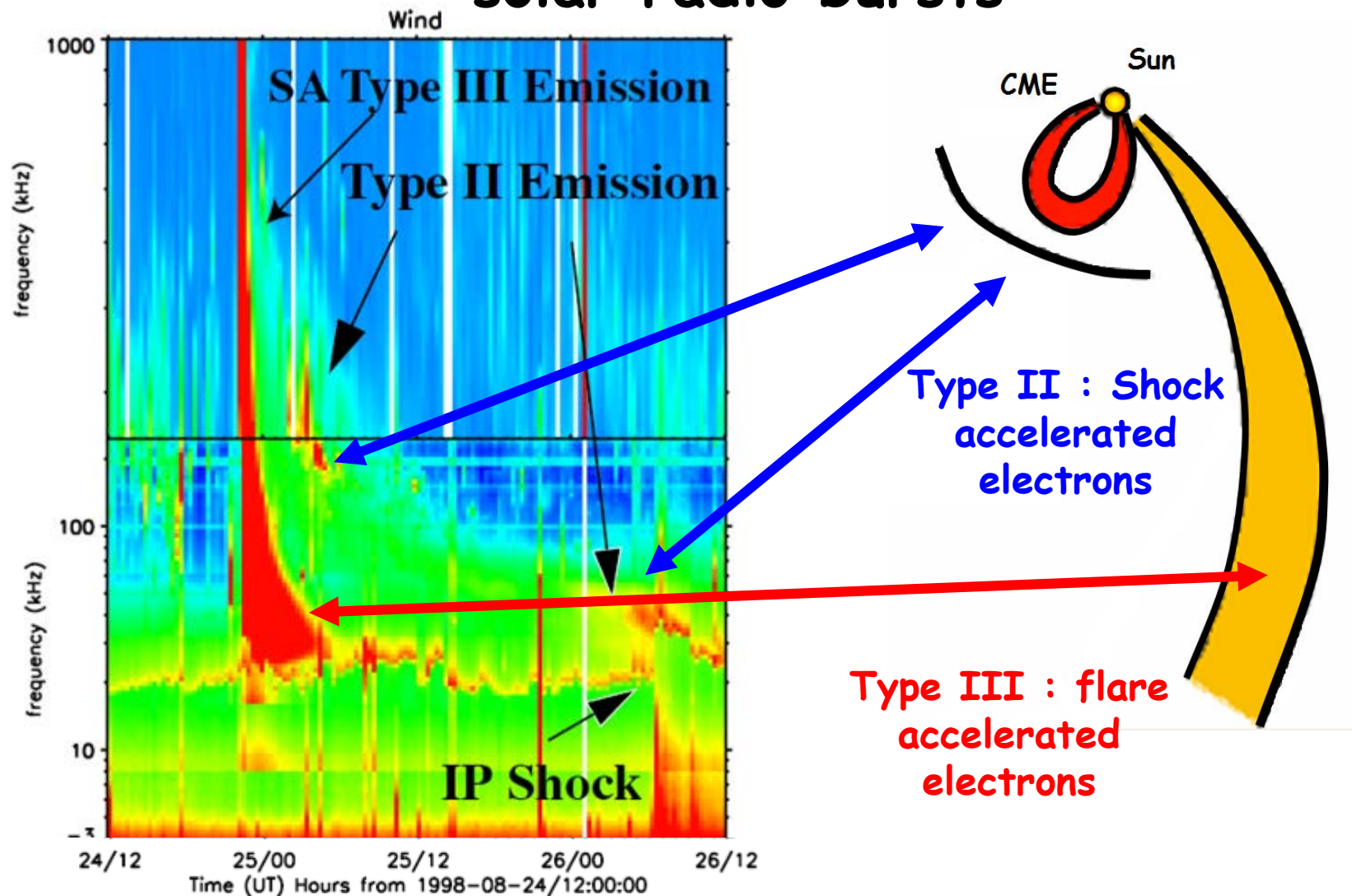


07 APR 2002

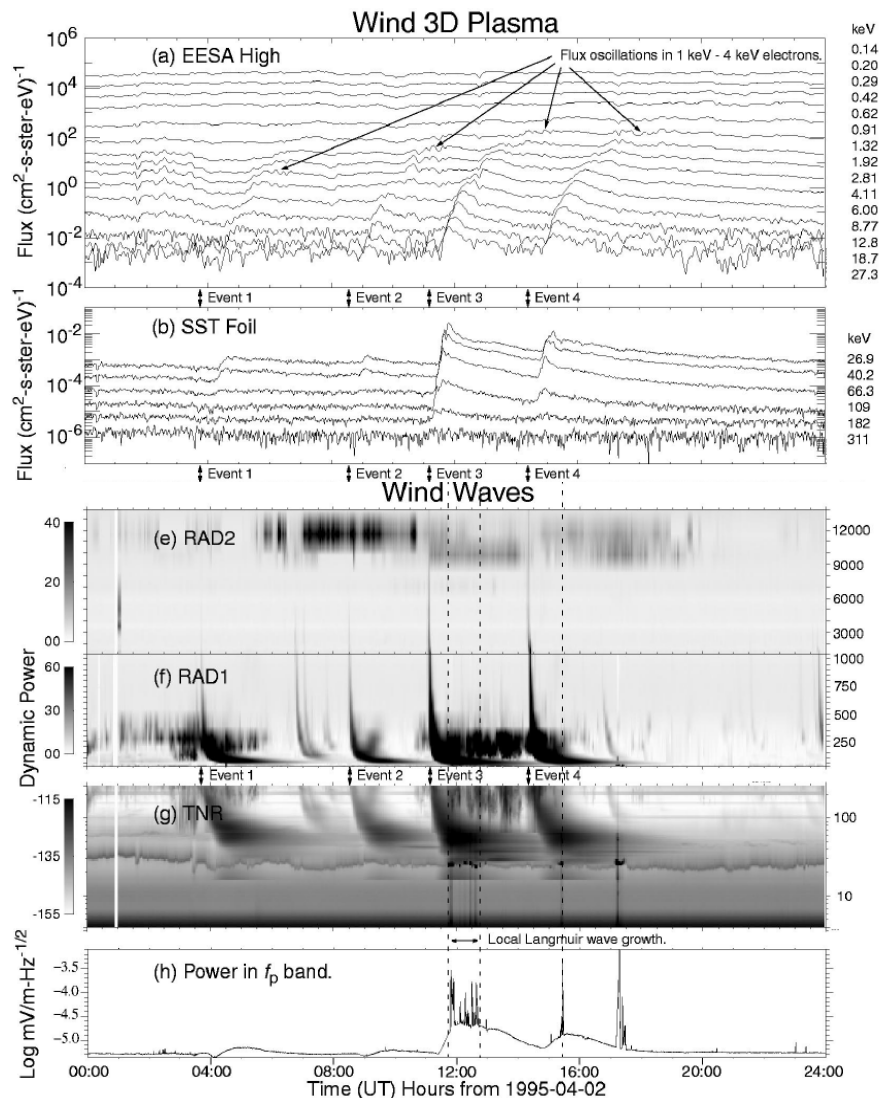


**Metric to
hectometric
emissions
first !**

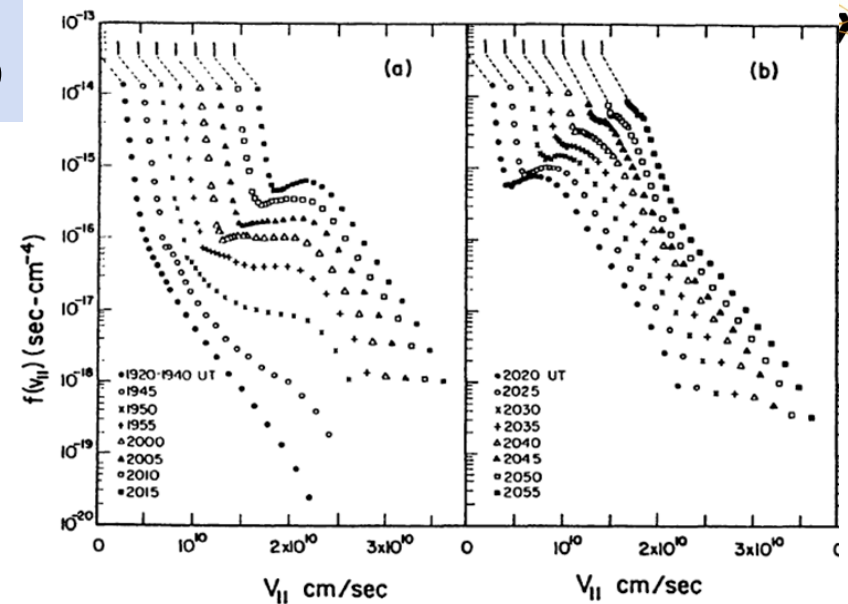
Type III & Type II solar radio bursts



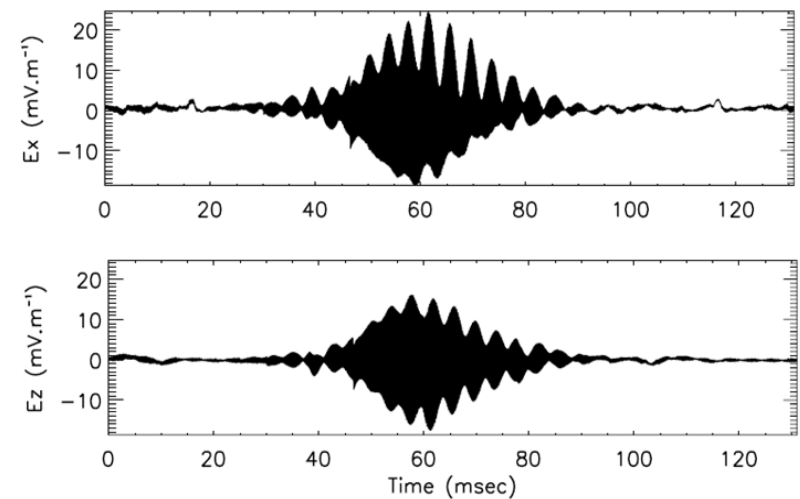
In-situ Type III and II measurements will be available on SO



Adapted from [Ergun et al., 1998]



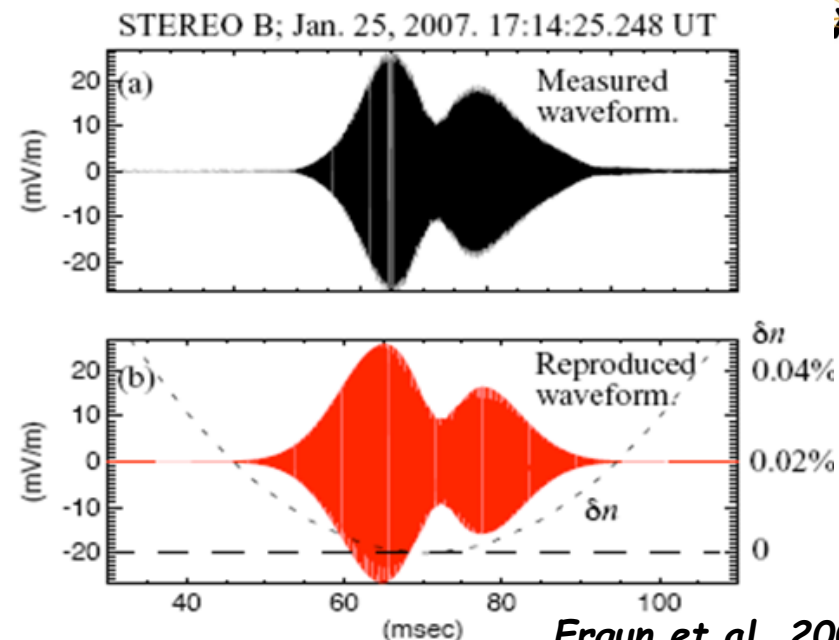
Adapted from [Lin et al. 1981]



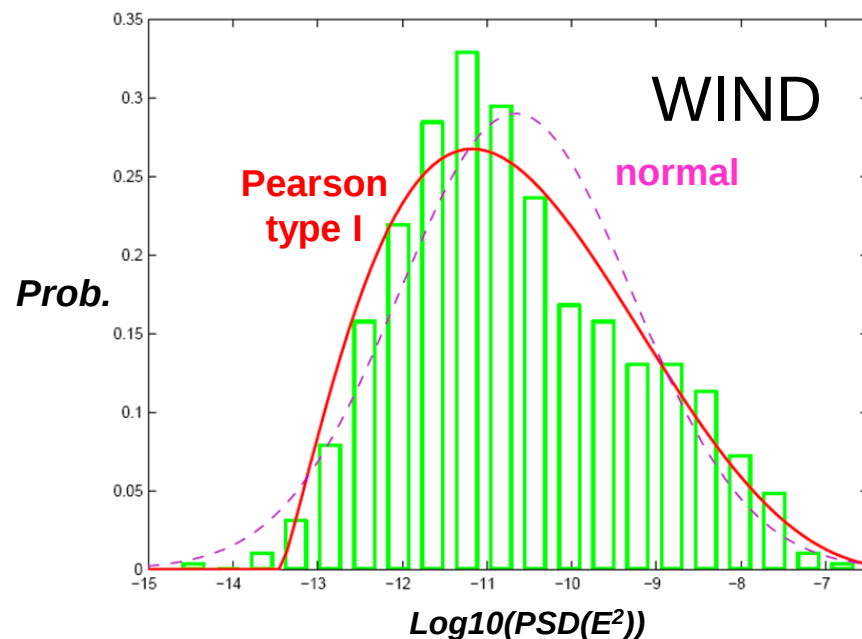
Swaves

RPW improvements for in-situ waves measurements

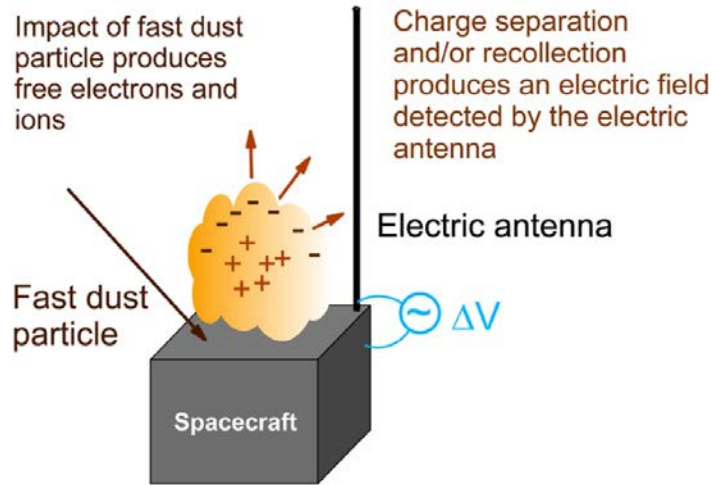
- RPW will measure both Langmuir waves and density fluctuations (from S/C pot. fluctuations, biased antennas)
DC E field and cross-shock potential
- RPW will measure simultaneously 2-axis E + 1 axis B up to 500 kHz → mode conversion
- Onboard statistics of LW power
Distribution of LW power more Pearson like than Normal ?
(Vidojevic et al., 2010)



Ergun et al, 2008

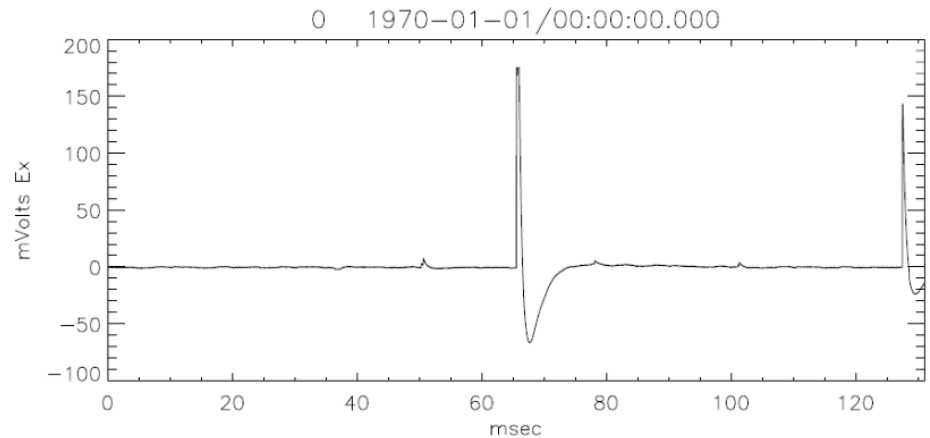
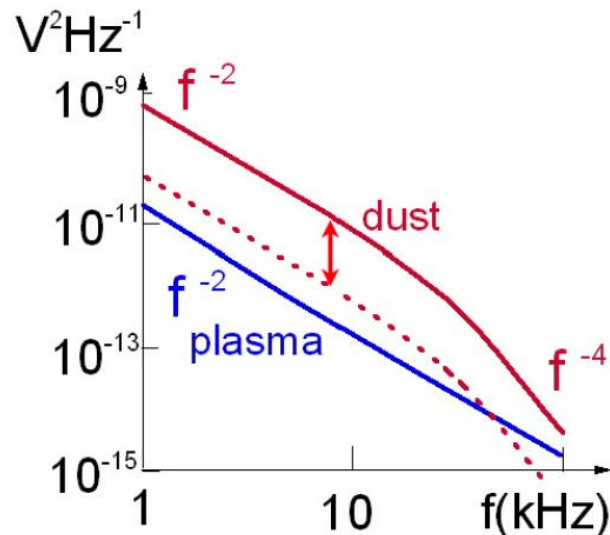


Interplanetary Dust



Released charge : $Q \simeq 0.7m^{1.02}v^{3.48}$

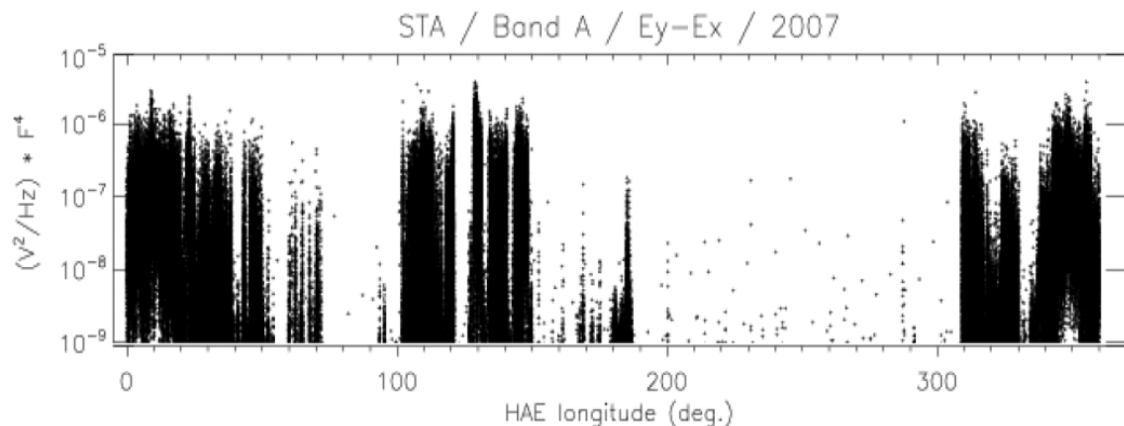
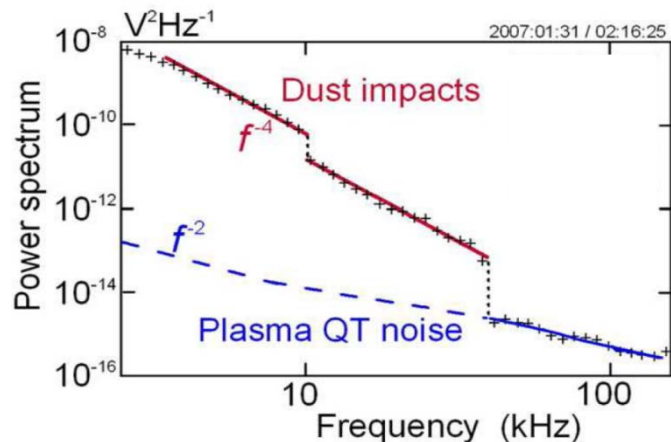
Induced voltage pulse on S/C of capacitance C : $\delta V \sim -Q/C$



Temporal domain

Spectral domain

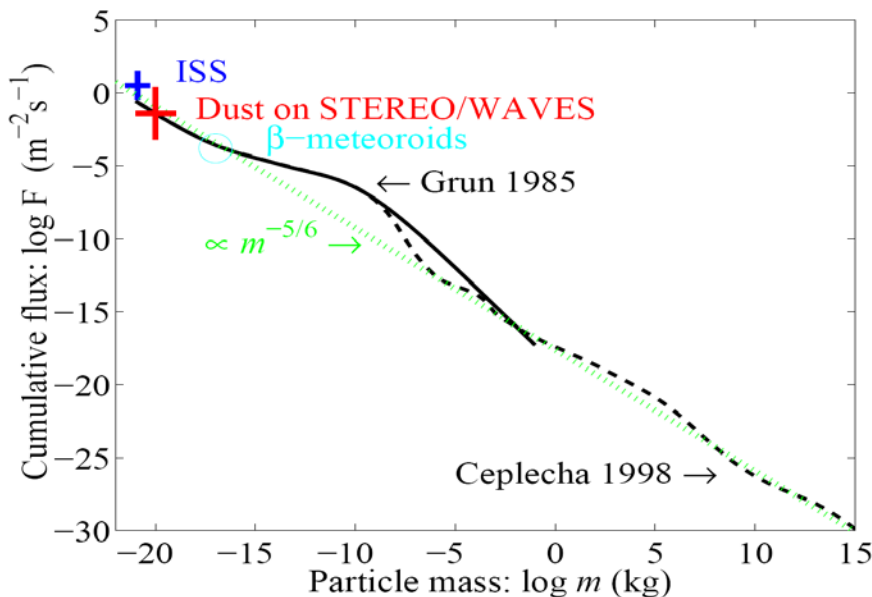
Discovery of a large flux of nanodust on Stereo Meyer-Vernet et al., 2012



Released charge : $Q \simeq 0.7 m^{1.02} v^{3.48}$

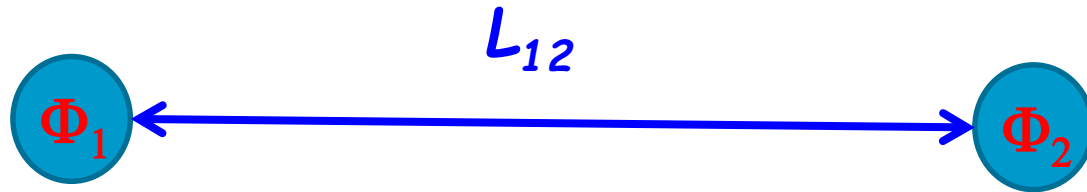
→ A nanoparticle @ 300 km/s ~ a grain of mass 10⁴ greater @ 20 km/s

Picked-up by the -VXB field

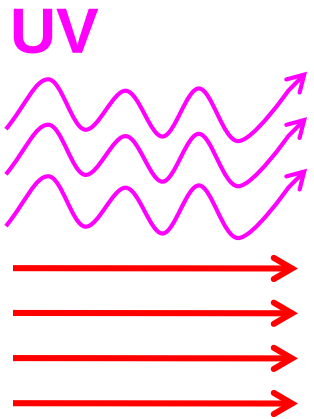


On the difficulty of measuring the DC/LF electric component of plasma waves

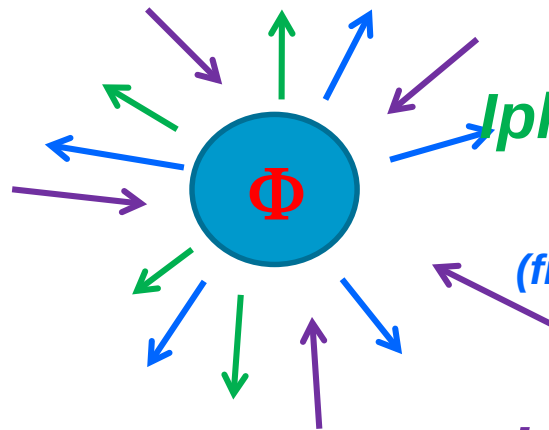
Basics for measuring DC/LF E Field



$$|E| = |\Phi_2 - \Phi_1| / L_{12}$$



Protons :
 $I_p \propto N_p V_p$
Negligible at 0th order



Photoelectrons

$I_{ph} \propto N_{ph} V_{ph}$ corrected for Φ

Secondary electrons

(from impact of external energetic electrons)

SW electrons

$I_e \propto N_e V_{the}$ corrected for Φ

Φ determined by $I_p + I_e + I_{ph} = 0$,
 actually $I_e = -I_{ph}$ (depending mostly on N_e)

$$J_{ph0} \approx \frac{10^{-5} \chi}{d_{AU}^2}$$

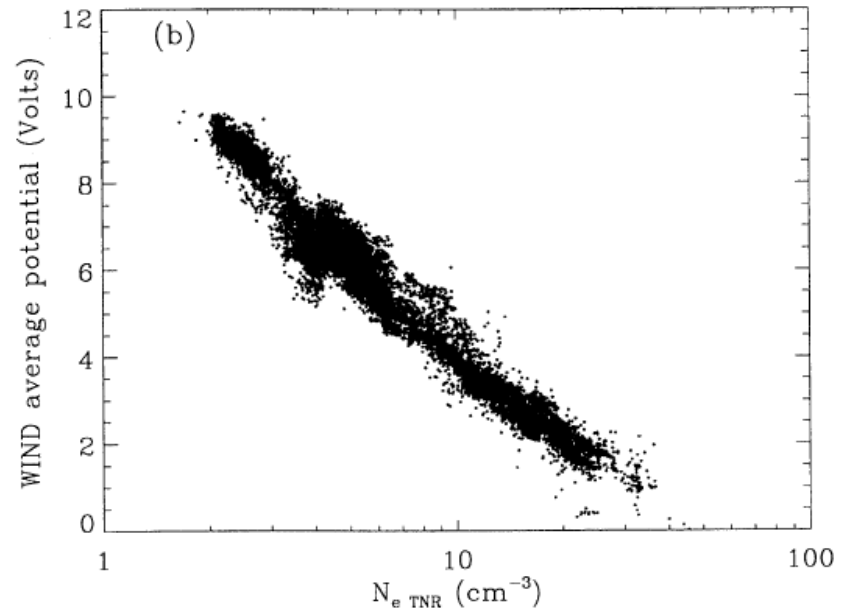
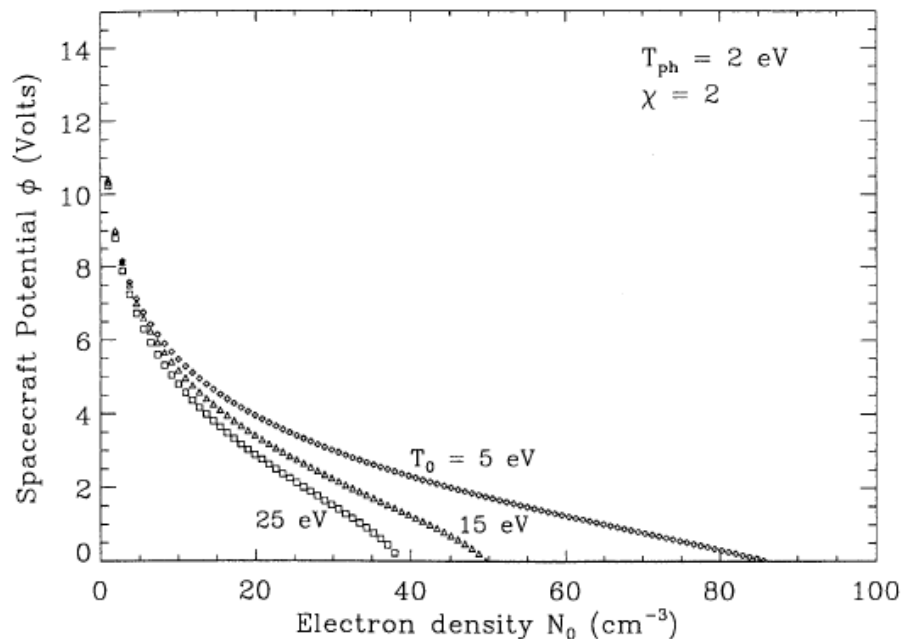
$$J_{ph} = J_{ph0} \left(1 + \frac{e\phi}{k_B T_{ph}}\right) e^{-e\phi/k_B T_{ph}},$$

$$J_e = e N_e v_e,$$

$$J_e = J_{e0} \left(1 + \frac{e\phi}{kT_{e0}}\right)$$

Φ varies with N_e !!

We want to measure $\sim mV/m$!!



We need to bias the antennas

$$I_e + I_{ph} + I_{bias} = 0$$

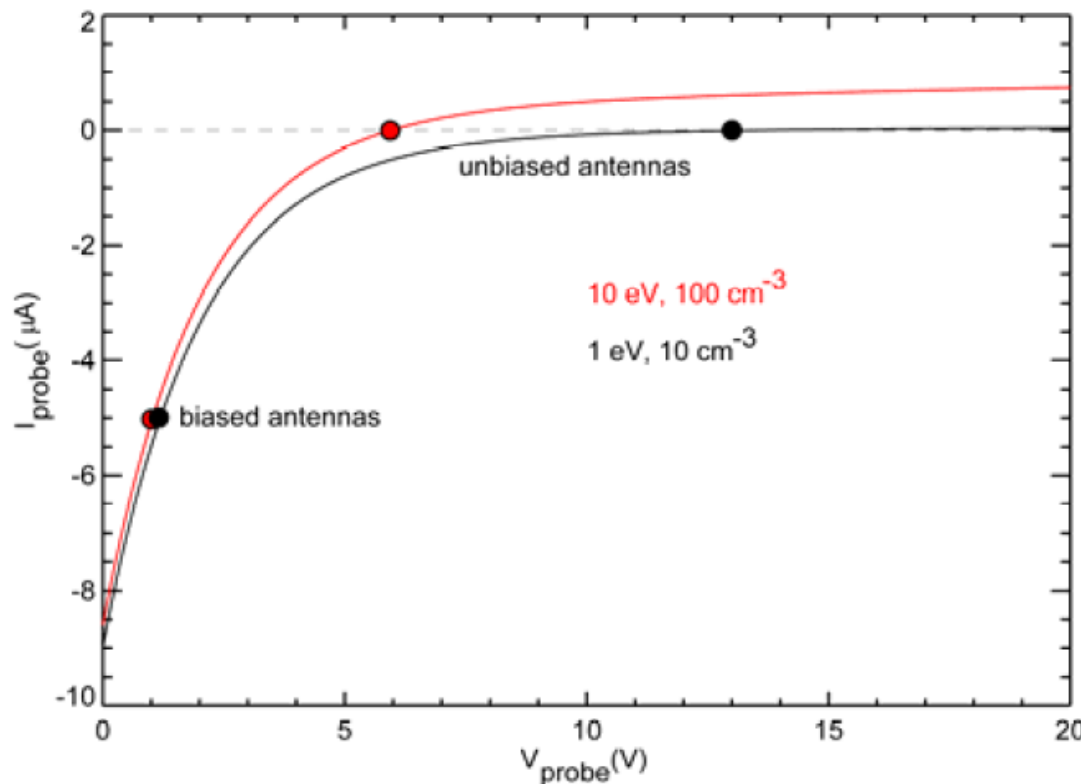
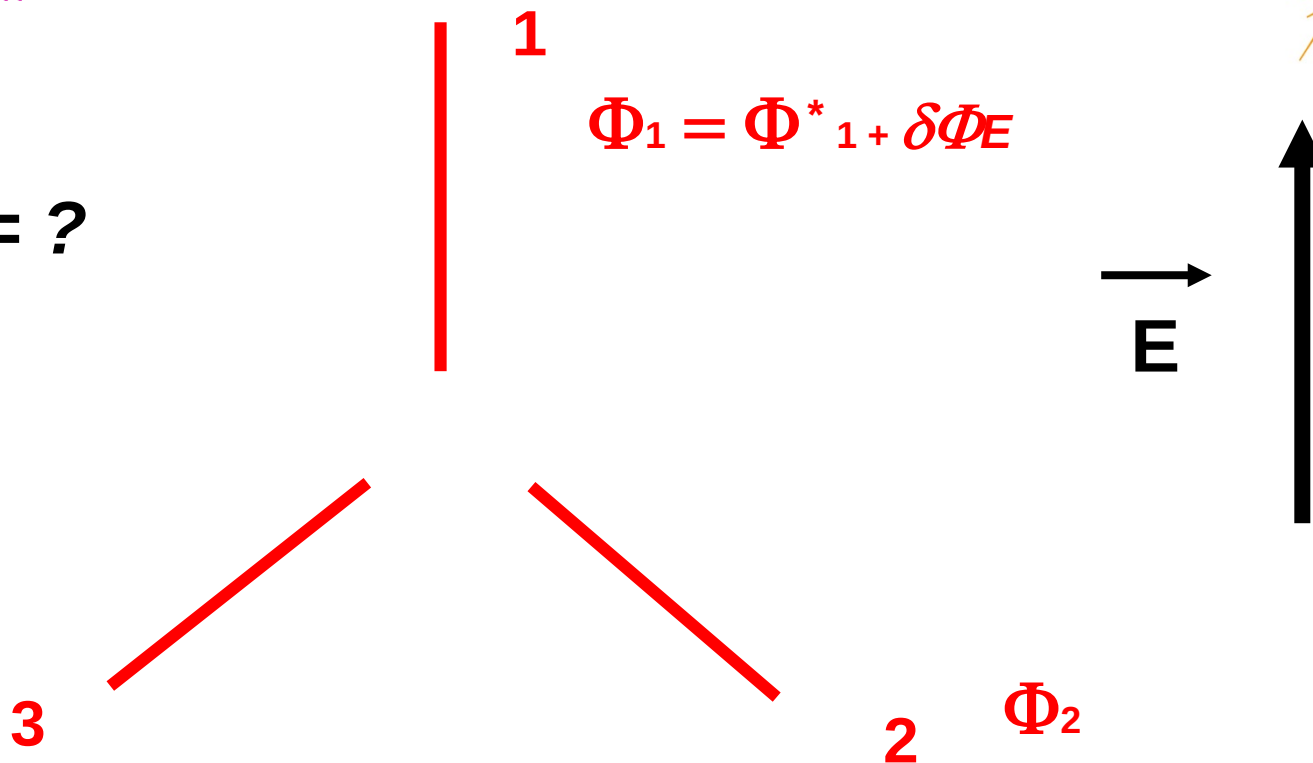


Figure 2.4 The simulated voltage-current curve of the antenna for two different plasma regimes. The potential variations of biased antenna are significantly less for different plasma parameters than for unbiased antenna.

courtesy C. Cully

$$\delta\Phi_E = ?$$



If $\Phi_1^* = \Phi_2$ then $\Phi_2 - \Phi_1 = \delta\Phi_E$

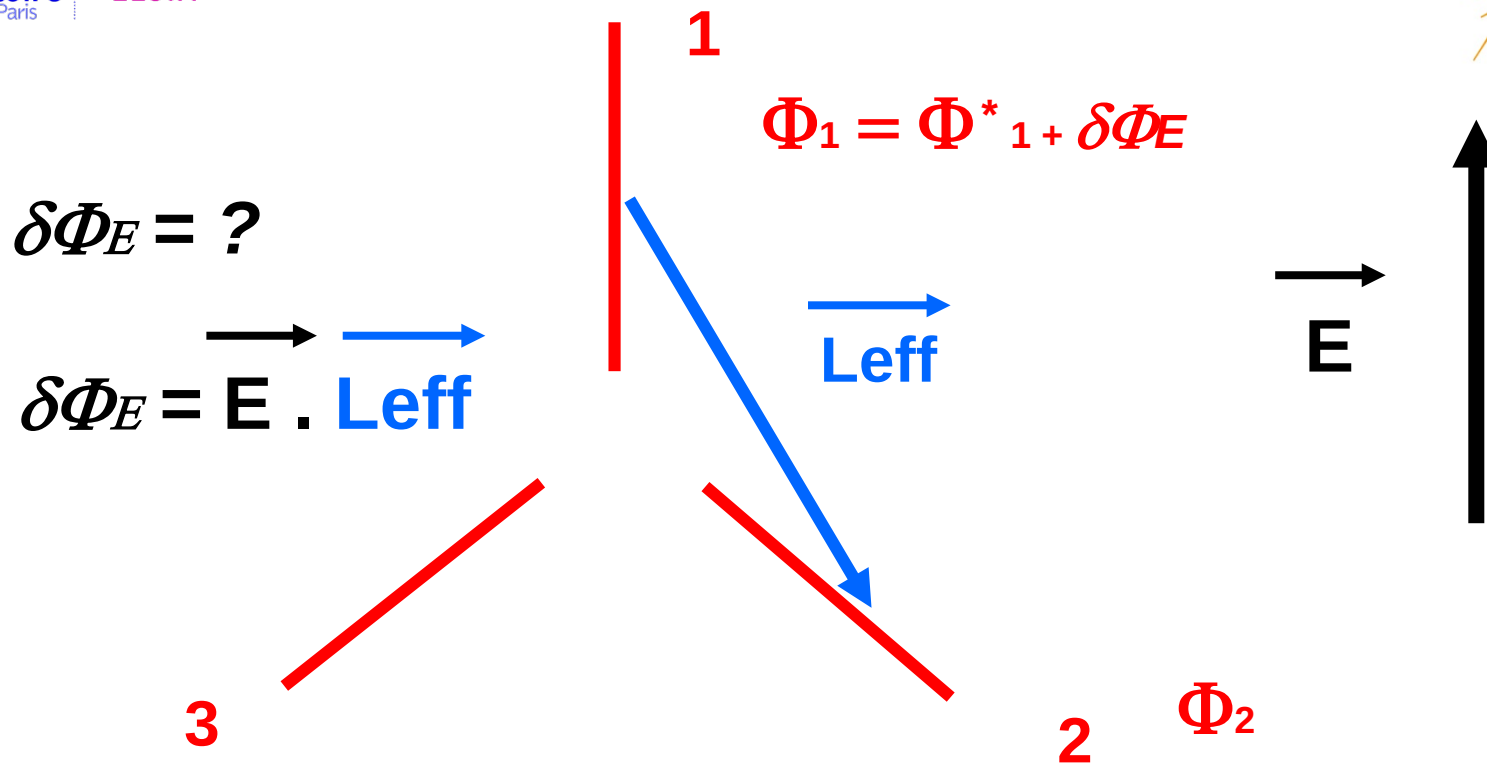
For that we need :

- Equal illumination for 1 & 2
- Symmetry with respect to the S/C
- Biasing the probes

Typically

If $\Phi_1^* = \Phi_2 \sim 0$ to 10 Volts

$E \sim$ few to few 100s of mV/m



If $\Phi_1^* = \Phi_2$ then $\Phi_2 - \Phi_1 = \delta\Phi_E$

For that we need :

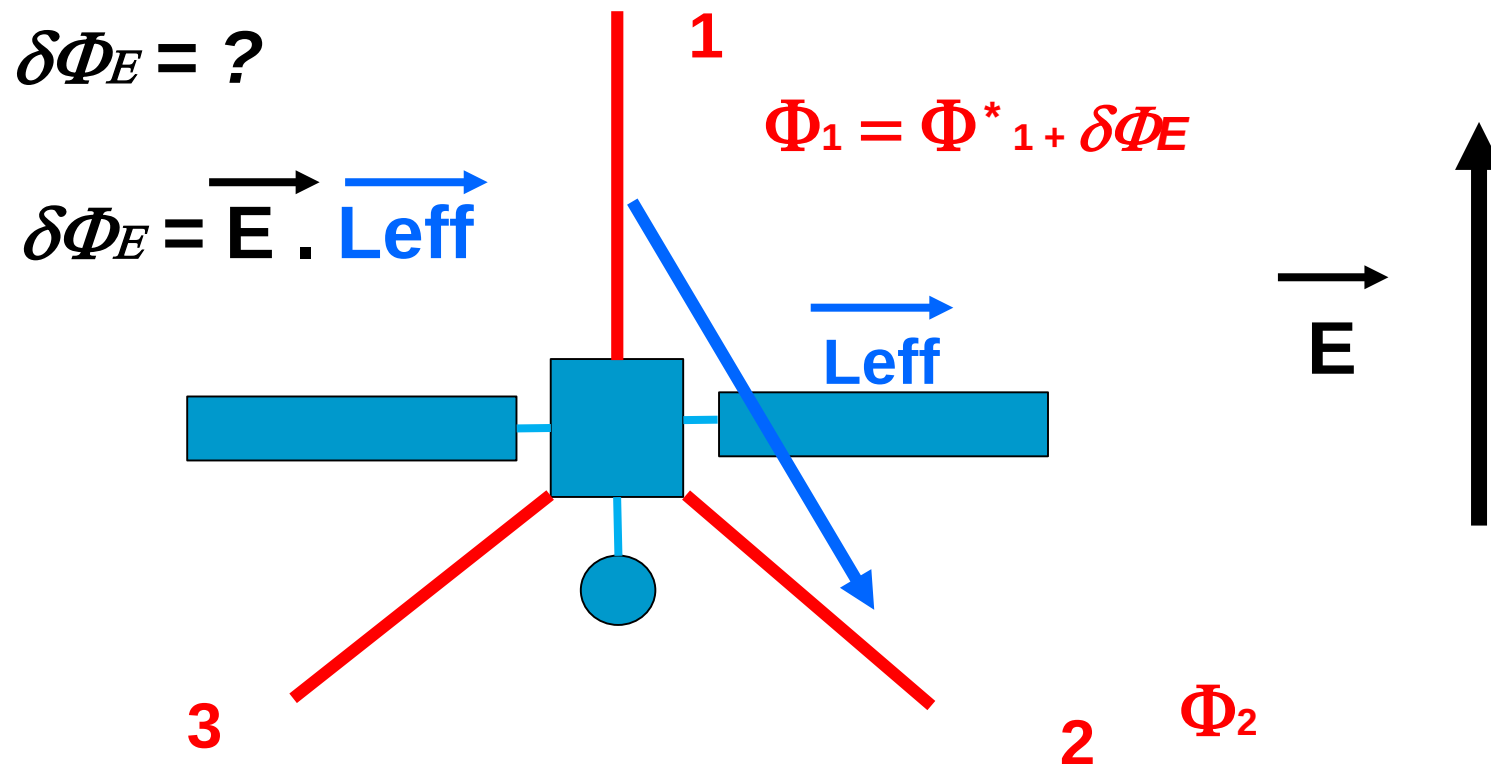
- Equal illumination for 1 & 2
- Symmetry with respect to the S/C
- Biasing the probes

Typically

If $\Phi_1^* = \Phi_2 \sim 0$ to 10 Volts

$E \sim$ few to few 100s of mV/m

$\vec{L}_{eff}$ can only be determined by simulation



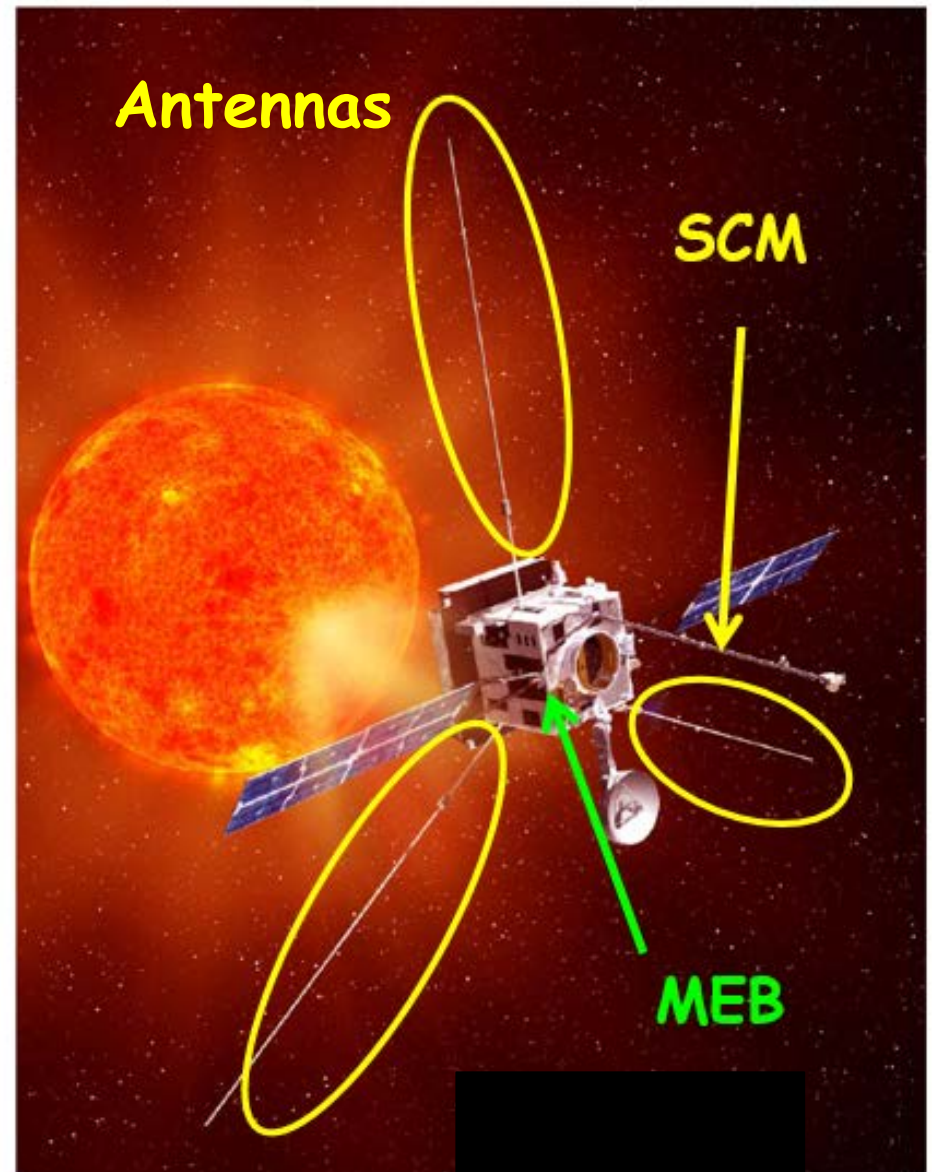
The problem is that there is a S/C in addition to the RPW antennas !!

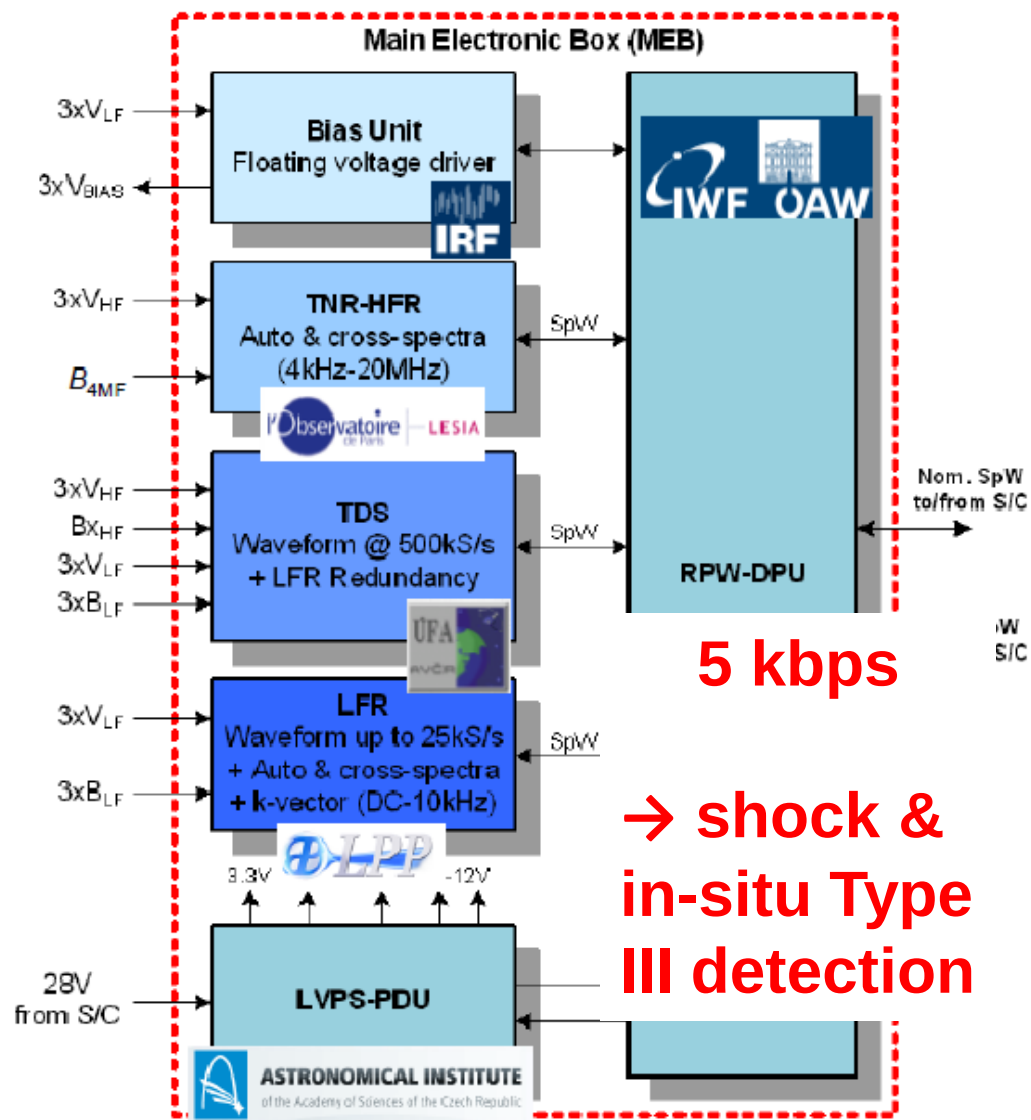
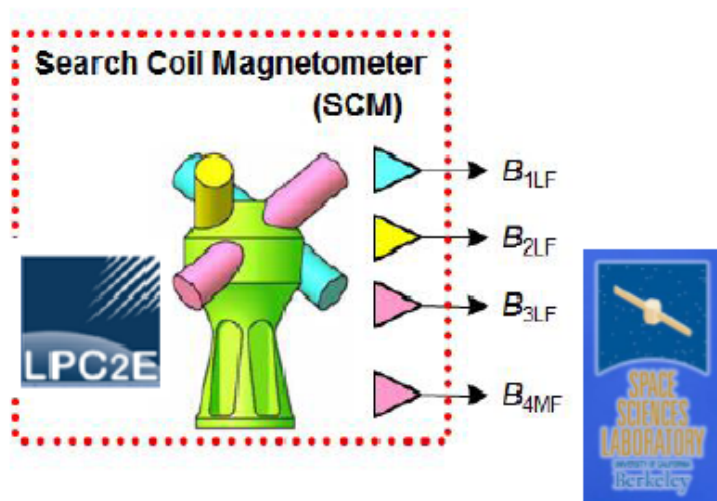
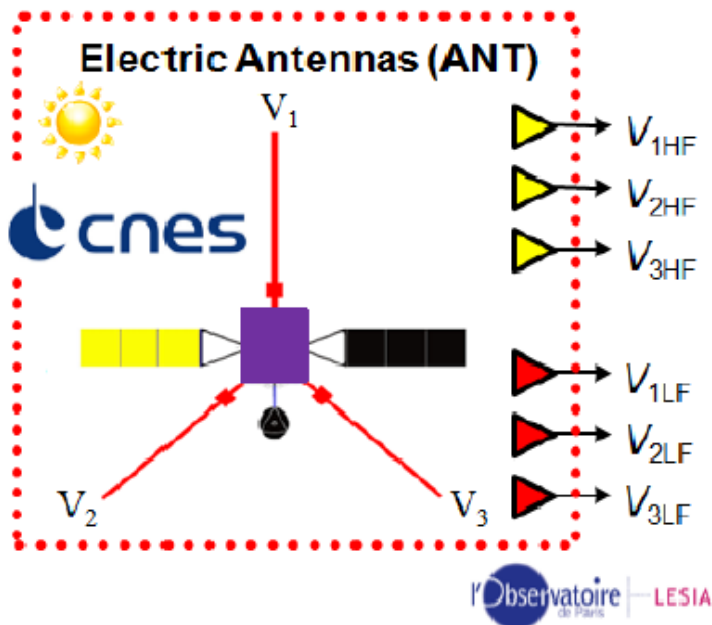
Need to simulate the effect of the S/C

RPW : a (very) brief description

Solar Orbiter

Exploring the Sun-Earth connection





SCM sensitivity + HELIOS variations

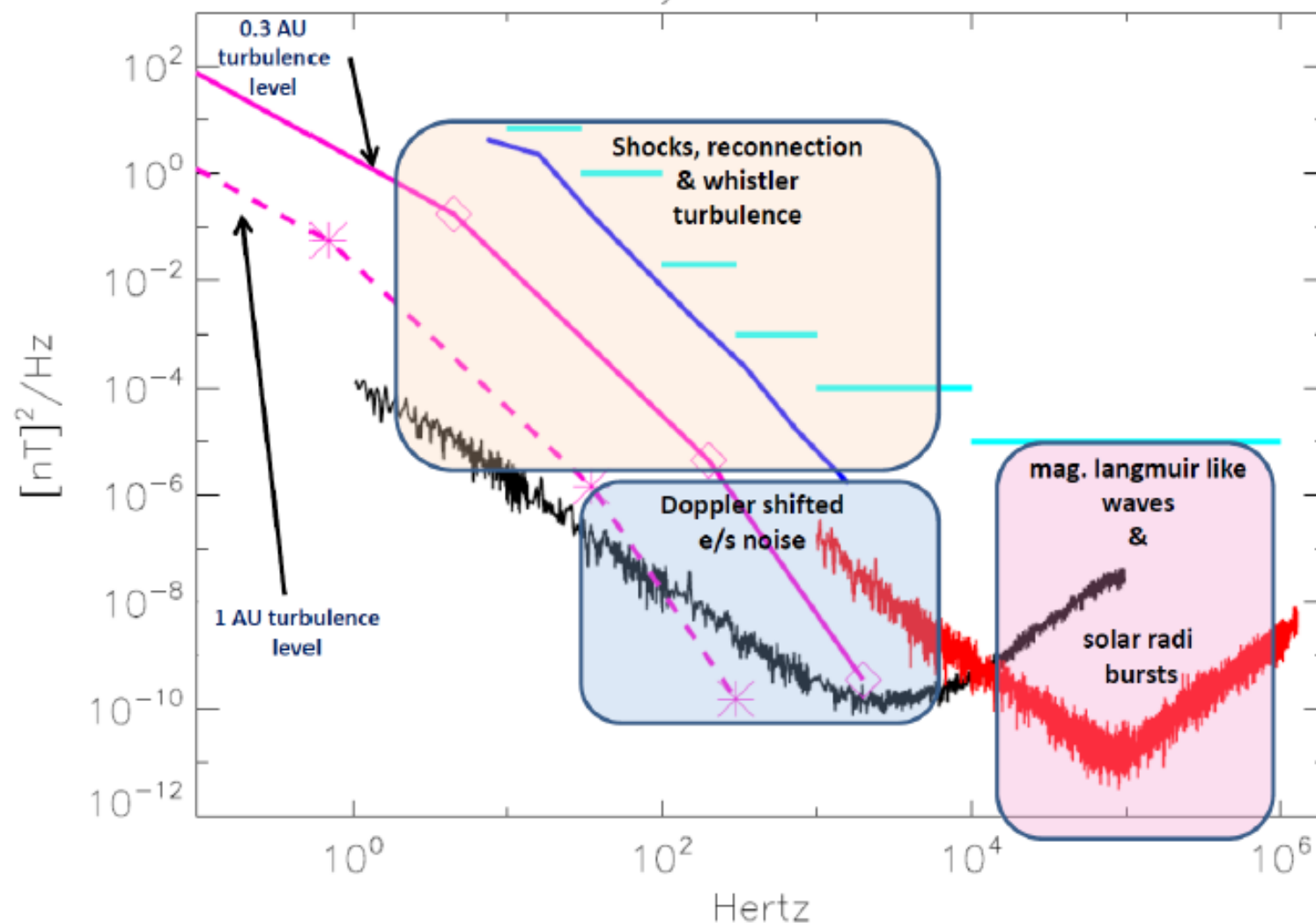
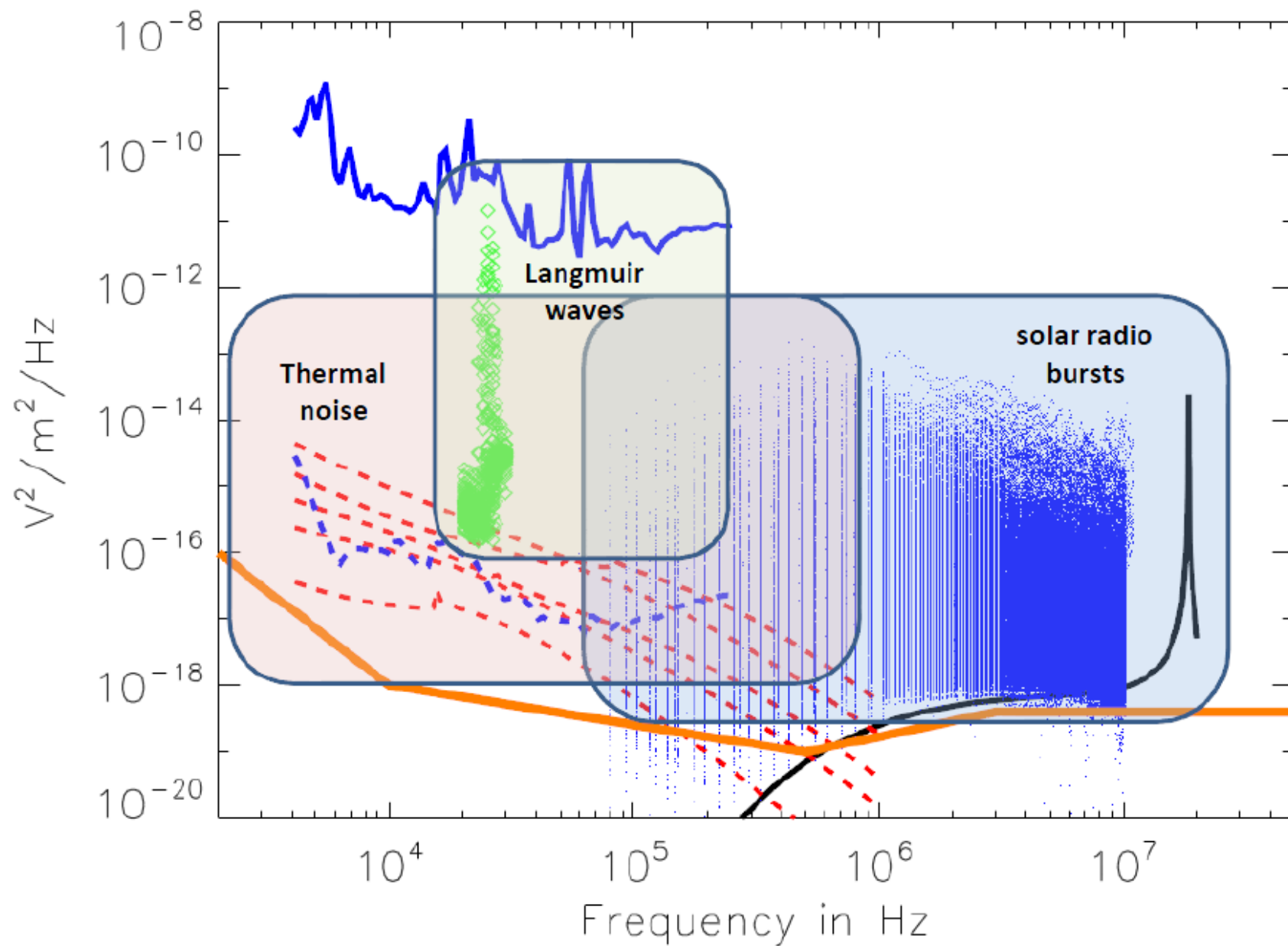
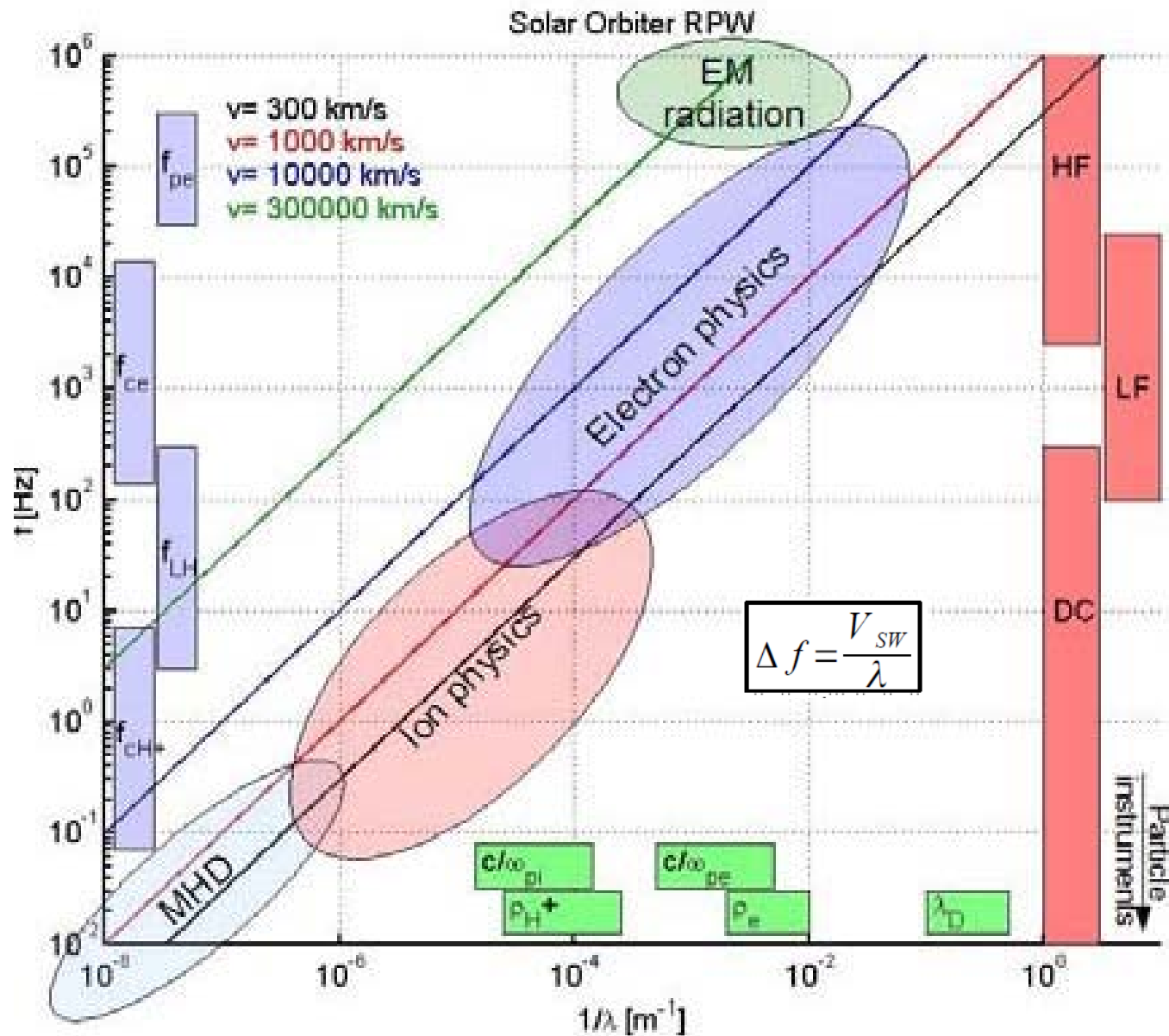


Figure 10: Magnetic field fluctuations of various natural phenomena that will be observed by the RPW instrument



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“Instantaneous” 5 x 5 spectral matrix
(256 FFT points)

$$\mathbf{SM}(\omega_j^{(m)}) = \begin{bmatrix} B_1 B_1^* & B_1 B_2^* & B_1 B_3^* & B_1 E_1^* & B_1 E_2^* \\ cc & B_2 B_2^* & B_2 B_3^* & B_2 E_1^* & B_2 E_2^* \\ cc & cc & B_3 B_3^* & B_3 E_1^* & B_3 E_2^* \\ cc & cc & cc & E_1 E_1^* & E_1 E_2^* \\ cc & cc & cc & cc & E_2 E_2^* \end{bmatrix}$$

Time Averaged Spectral Matrix (ASM)

$$\mathbf{ASM}(\omega_j^{(m)}) = \frac{1}{N_{SM}^{(m)}} \sum_{k=1}^{N_{SM}^{(m)}} \mathbf{SM}_k(\omega_j^{(m)}) = \langle \mathbf{SM} \rangle_{time}$$

Frequency average ...

$$\mathbf{S}(\omega_j^{(m)}) = \langle \mathbf{ASM} \rangle_{frequency}$$

... before computations of the BPs
(i.e. wave parameters)

Mono-**k**
assumption :

(Means, JGR, 1972) {

(Samson & Olson, GJRA, 1980) {

$$\mathbf{n} \times \mathbf{E} = \frac{\omega}{k} \mathbf{B} \longrightarrow$$

$$\frac{S_{ij}}{\sqrt{S_{ii} S_{jj}}} \longrightarrow$$

- BP1** set 1: Power spectrum of the magnetic field (**B**)
- BP1** set 2: Power spectrum of the electric field (**E**)
- BP1** set 3: Wave normal vector (from **B**)
- BP1** set 4: Wave ellipticity estimator (from **B**)
- BP1** set 5: Wave planarity estimator (from **B**)
- BP1** set 6: X_{so} (radial)-component of the Poynting vector
- BP1** set 7: Phase velocity estimator

- BP2** set 1: Autocorrelations
- BP2** set 2: Normalized cross correlations