Tidal heating and orbital evolution of terrestrial exoplanets

Michaela Walterová and Marie Běhounková Faculty of Mathematics and Physics Charles University **numerical modelling** of tidal effects (tidal deformation, heating, orbital and rotational evolution) on rocky exoplanets without atmosphere

Maxwell or Andrade viscoelastic rheology

parametric studies and qualitative comparison with recent analytical models (Correia et al., 2014)

Observational motivation

tidally evolved moons in the Solar system

increasing number of detected exoplanets (preferentially close-in)

increasing number of detected terrestrial exoplanets

Theoretical motivation

analytical models describe main features of orbital evolution

internal structure and more complex rheologies can be taken into account by numerical models

possible coupling of orbital and internal evolution

Model and methods

governing equations for the planetary mantle are discretized using **spherical harmonic decomposition** in the lateral directions and **staggered finite difference scheme** in the radial direction

computation in the time domain

tidal deformation \rightarrow tidal heating and boundary deflections \rightarrow disturbing potential and disturbing force \rightarrow orbital evolution (Gauss's planetary equations)

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I. Introduction

Tides play a major role in the long-term evolution of dose-in excolanets and planetary satellites. Tidal effects include eilanced heating of the planetary interior, secular despinning with eventual locking into a spin-orbit resonance, and secula changes in orbital elements. Traditionally, the response of a planetary body to the tidal potential is modeled using simplified theological assumptions that are, howe valid only in special cases (e.g. Kaula, 1964; Mignard, 1979). A more general description of tides in terrestrial planets and moons, taking into account viscoelastic effects, has been developed only secondly by several authors (e.g. Henning et al., 2009, Makazov and Efformsky, 2018; Ferraz-Mello, 2018; Cometa et al., 2014).

Aforementioned studies present analytical methods of addressing the main features of tides on a homogeneous planetary body. Alternatively, the tidal interaction can be described numerically. A numerical model (such as Kelvin-Voigt mass-spring model of Proused et al., 2016) would erable implementation of more complex theologies and internal structures and description of phaenomena that are not easy to approach analytically. Here, we present a numerical model based on the direct evaluation of tidal deformation of a planetary mantle governed either by the Maxwell or the Andrade sheelogy. The model is then utilized for the evaluation of tidal heating, tidal deformation and the long-term orbital evolution.

Once we have computed the tidal deformation inside of the planet, it is

possible to evaluate the disturbing

potential and the disturbing force

associated with boundary deflections.

While the disturbing potential at the

surface enables us to compute the

potential Love number k the disturbing

force at the host star's position may be

inserted into the Gauss's planetary

equations for the sent-major axis and the

additional evolution equation for the spin

rate (the total angular momentum has to

be conserved). The obliquity of the planet

is assumed to be zero. We compute the set

of evolution equations explicitly on two time-scales: on the shorter time-scale we

seek the equilibrium spin rate

corresponding to the actual orbital

parameters, on the longer time scale we

take a step in the semi-major axis and in

Exo-Barth orbiting a host star with 0.1 solar

Semi-major axis for the tidal heating and

Love number computation: 0.1 AU.

entricity, supplemented with an

d Orbital evaluation

II. Model and Methods at Tidal deformation

The planet is mobiled as a spherical hody composed of a solid inser core, liquid outer core and a viceolistic market. We compute the tidal delormation and disspation oblig in the market, using a spherical hearness decomposition in the latest directore and a staggered finite difference elsewin the radial direction (Toble et al., 2006). The fights of eligiparity, pressure within the liquid core and creatingsi force are indiated in the model.

b) Tidal beating

The average power associated with the tidal leating of the mantle is computed as

$$P_t = \int_{V_t} \frac{1}{T_t} \int_{\tau}^{\tau+T_t} \delta(\tau) : \mathbf{B}(\tau) \, d\tau \, dV$$

where D symbolizes the deviance part of the Caudy stress theorized of in the trainme tensor. We evaluate the heating once one orbital pentid and investigate its parameter dependence. In the case of the Mannell Heating we compare the fidal heating for effective (rafal) viscosities, which are much lower than standard viscosities used in markle convection models (e.g. Bithoushout et al., 2018).



- 6.4

Sector that evolution of the sect-maps and (top left), eccentrary of the orbit (bottom left), optio-orbit ratio (top 0 right) and that heating (bottom right). Initial sectementary of the model planet, governed by the Maxwell theology, was either q= 0.2 (black dashed line) or q= 0.4 (nd solid line). Effective that is successful of the planet that 0² Plan.

III. Results and Conclusions

a) Love numbers

Protocial Love numbers of visco-lattic bodies are frequency-dependent quantities. At the table likeling of planets on accentre others operates on multiple frequencies, the Loverumbers, as well as the phase lag between the table and the distanting potential, vary during the other. Similarly, if we keep a constant frequency, the Love runnless vary with increasing viscosity of the match, as is down on the figures beine (cf. analytical results of Corvena et al., 2014).

b) Tidal heating

We evaluated the average tidal heating as a function of the spin-orbit ratio and the eccentricity. Due to different frequency dependencies of dissipation in the Marwell

and the Andrade model, the tidal heating in the first case achieves deep and narrow local minima, while in the laner case the local minima are broad and duallow.

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d Orbital evaluation

Table evolution of the place leads to the conclusions of its orbit. As the eccentricity guidably decays, the placet advance neuronal opailblotum open states: the apple-orbits and the placet of the state open orbits, and a 11 monotone or the state operformations. Rach manifestion between two open-orbits resonances the precoded by an era of penesuranel tidal heating. This may have consequences for the insert open-of evolution, including the type of states.

(g & a consequences for the internal evolution of the including the type of mantle convection possible volcanous on the planet. accell



the econdricity.

d) Model planet

Andrade rheology Frequency dependence of tidal beating similar to the fluid limit, minima are shallow.



Average tidal heating of the mande as a function of the spin-orbit ratio $(t_i)_{\gamma}$ and the eccentricity e for both theologies and three different viscosities. Prominent features - local minima of tidal heating - are associated

with anim-orbit resonances, where the notational frequency equals integer or half-integer multiple of the

orbital frequency. Relative depth of distinct minima depends on the eccentricity.



India log (left) and Love number t_0 (right) of a Maximum oxy at a number of effective tidal variety. The spin-orbit ratio of the model planet with e = 0was set unstant, $\Omega/n = 2.1$.

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