Kappa Distributions: Turbulent Equilibria 2. Beam-Plasma Instability & Langmuir Turbulence

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- 2 Beam-Plasma Instability and Langmuir Turbulence
- 3 Turbulent Equilibrium and Kappa Distribution
 - Quiet Time Solar Wind Electron Distribution

5 Conclusion







Dispersion Relation

$$\operatorname{Re} \epsilon(k,\omega) < E^2 >_{k\omega} = 0.$$

Particle Kinetic Equation

$$\frac{\partial F_{a}}{\partial t} = \operatorname{Re} i \, \frac{e_{a}^{2}}{m_{a}^{2}} \frac{\partial}{\partial v} \int dk \int d\omega \, \frac{\langle E^{2} \rangle_{k\omega}}{\omega - kv + i0} \frac{\partial F_{a}}{\partial v}.$$







Wave Kinetic Equation

$$\frac{\partial}{\partial t} < E^{2} >_{k\omega} = -\frac{2 \operatorname{Im} \epsilon(k, \omega)}{\partial \operatorname{Re} \epsilon(k, \omega)/\partial \omega} < E^{2} >_{k\omega}$$

$$-\frac{4}{\partial \operatorname{Re} \epsilon(k, \omega)/\partial \omega} \operatorname{Im} \int dk' \int d\omega'$$

$$\times \left[\{\chi^{(2)}(k', \omega'|k - k', \omega - \omega')\}^{2} \left(\frac{< E^{2} >_{k-k', \omega - \omega'}}{\epsilon(k', \omega')} \right.$$

$$+ \frac{< E^{2} >_{k'\omega'}}{\epsilon(k - k', \omega - \omega')} \right] < E^{2} >_{k\omega}$$

$$- \frac{|\chi^{(2)}(k', \omega'|k - k', \omega - \omega')|^{2}}{\epsilon^{*}(k, \omega)} < E^{2} >_{k'\omega'} < E^{2} >_{k-k', \omega - \omega'} \right].$$

$$\operatorname{Re}\epsilon(k,\omega) < E^2 >_{k\omega} = 0, \qquad \Rightarrow \qquad \omega = \omega_k^\alpha \quad (\alpha = L,S).$$

$$< E^{2} >_{k\omega} = \sum_{\alpha} \left[I_{k}^{+\alpha} \,\delta(\omega - \omega_{k}^{\alpha}) + I_{k}^{-\alpha} \,\delta(\omega + \omega_{k}^{\alpha}) \right]$$
$$= \sum_{\sigma=\pm 1} \sum_{\alpha} I_{k}^{\sigma\alpha} \,\delta(\omega - \sigma\omega_{k}^{\alpha}).$$







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Langmuir ($\alpha = L$) and ion-sound wave ($\alpha = S$)

$$\begin{split} \omega_{k}^{L} &= \omega_{pe} \left(1 + \frac{3}{2} \, k^{2} \lambda_{De}^{2} \right) = \omega_{pe} \left(1 + \frac{3}{4} \frac{k^{2} v_{Te}^{2}}{\omega_{pe}^{2}} \right). \\ \omega_{-k}^{L} &= -\omega_{k}^{L}, \\ \omega_{k}^{S} &= \frac{kc_{S} \, (1 + 3T_{i}/T_{e})^{1/2}}{(1 + k^{2} \lambda_{De}^{2})^{1/2}}, \qquad \omega_{-k}^{S} = -\omega_{k}^{S}. \end{split}$$







Particle kinetic equation,

$$\begin{split} &\frac{\partial F_{a}}{\partial t} = \frac{\partial}{\partial v} \left(D \frac{\partial F_{a}}{\partial v} \right), \\ &D = \frac{\pi e_{a}^{2}}{m_{a}^{2}} \sum_{\sigma = \pm 1} \sum_{\alpha} \int dk \ I_{k}^{\sigma \alpha} \delta(\sigma \omega_{k}^{\alpha} - kv). \end{split}$$







Wave kinetic equation

$$\begin{split} &< E^2 >_{k\omega} = \sum_{\sigma = \pm 1} \sum_{\alpha} I_k^{\sigma \alpha} \, \delta(\omega - \sigma \omega_k^{\alpha}), \\ &< E^2 >_{k'\omega'} = \sum_{\sigma' = \pm 1} \sum_{\beta} I_{k'}^{\sigma' \beta} \, \delta(\omega' - \sigma' \omega_{k'}^{\beta}), \\ &< E^2 >_{k-k',\omega-\omega'} = \sum_{\sigma'' = \pm 1} \sum_{\gamma} I_{k-k'}^{\sigma'' \gamma} \, \delta(\omega - \omega' - \sigma'' \omega_{k-k'}^{\gamma}), \end{split}$$







$$\begin{split} \frac{1}{\epsilon(k,\omega)} &= \mathcal{P} \, \frac{1}{\epsilon(k,\omega)} - \sum_{\alpha} \sum_{\sigma=\pm 1} \frac{i\pi \, \delta(\omega - \sigma \omega_k^{\alpha})}{\partial \text{Re} \, \epsilon(k,\sigma \omega_k^{\alpha})/\partial \sigma \omega_k^{\alpha}}, \\ \frac{1}{\epsilon^*(k,\omega)} &= \mathcal{P} \, \frac{1}{\epsilon^*(k,\omega)} + \sum_{\alpha} \sum_{\sigma=\pm 1} \frac{i\pi \, \delta(\omega - \sigma \omega_k^{\alpha})}{\partial \text{Re} \, \epsilon(k,\sigma \omega_k^{\alpha})/\partial \sigma \omega_k^{\alpha}}. \end{split}$$







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Kappa Distribution

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$$\begin{split} &\frac{\partial l_{k}^{\sigma\alpha}}{\partial t} = -\frac{2\operatorname{Im}\epsilon(k,\sigma\omega_{k}^{\alpha})}{\partial\operatorname{Re}\epsilon(k,\sigma\omega_{k}^{\alpha})/\partial\sigma\omega_{k}^{\alpha}} l_{k}^{\sigma\alpha} \\ &-\frac{4}{\partial\operatorname{Re}\epsilon(k,\sigma\omega_{k}^{\alpha})/\partial\sigma\omega_{k}^{\alpha}} \operatorname{Im} \int dk' \\ &\times \left[\sum_{\sigma''=\pm 1} \sum_{\gamma} \mathcal{P} \frac{\{\chi^{(2)}(k',\sigma\omega_{k}^{\alpha}-\sigma''\omega_{k-k'}^{\gamma}|k-k',\sigma''\omega_{k-k'}^{\gamma})\}^{2}}{\epsilon(k',\sigma\omega_{k}^{\alpha}-\sigma''\omega_{k-k'}^{\gamma})} \ l_{k-k'}^{\sigma''\gamma} l_{k}^{\sigma\alpha} \right] \\ &+ \sum_{\sigma'=\pm 1} \sum_{\beta} \mathcal{P} \frac{\{\chi^{(2)}(k',\sigma'\omega_{k'}^{\beta}|k-k',\sigma\omega_{k}^{\alpha}-\sigma'\omega_{k'}^{\beta})\}^{2}}{\epsilon(k-k',\sigma\omega_{k}^{\alpha}-\sigma'\omega_{k'}^{\beta})} \ l_{k'}^{\sigma'\beta} l_{k}^{\sigma\alpha} \right] \end{split}$$



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$$+ \frac{4\pi}{\partial \operatorname{Re} \epsilon(k, \sigma \omega_{k}^{\alpha}) / \partial \sigma \omega_{k}^{\alpha}} \sum_{\sigma', \sigma''=\pm 1} \sum_{\beta, \gamma} \\ \times \operatorname{Im} \int dk' \left[\{\chi^{(2)}(k', \sigma' \omega_{k'}^{\beta} | k - k', \sigma'' \omega_{k-k'}^{\gamma})\}^{2} \\ \times \left(\frac{I_{k-k'}^{\sigma'' \gamma} I_{k}^{\sigma \alpha}}{\partial \operatorname{Re} \epsilon(k', \sigma' \omega_{k'}^{\beta}) / \partial \sigma' \omega_{k'}^{\beta}} + \frac{I_{k'}^{\sigma' \beta} I_{k}^{\sigma \alpha}}{\partial \operatorname{Re} \epsilon(k - k', \sigma'' \omega_{k-k'}^{\gamma}) / \partial \sigma'' \omega_{k-k'}^{\gamma}} \right) \\ + \frac{|\chi^{(2)}(k', \sigma' \omega_{k'}^{\beta} | k - k', \sigma'' \omega_{k-k'}^{\gamma})|^{2}}{\partial \operatorname{Re} \epsilon(k, \sigma \omega_{k}^{\alpha}) / \partial \sigma \omega_{k}^{\alpha}} I_{k'}^{\sigma' \beta} I_{k-k'}^{\sigma'' \gamma} \right] \\ \times \delta(\sigma \omega_{k}^{\alpha} - \sigma' \omega_{k'}^{\beta} - \sigma'' \omega_{k-k'}^{\gamma}).$$



Evaluations of nonlinear susceptibilites must be done. See the following references:

- P. H. Yoon, Generalized weak turbulence theory, POP, 7, 4858 (2000)
- L. F. Ziebell, R. Gaelzer, and P. H. Yoon, Nonlinear development of weak beam-plasma instability, POP, 8, 3982 (2001)
- P. H. Yoon, Effects of spontaneous fluctuations on the generalized weak turbulence theory, POP, 12, 042306 (2005)
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Induced Emission: Linear Wave-Particle Resonance

$$\begin{aligned} \frac{\partial I_k^{\sigma L}}{\partial t} \bigg|_{\text{ind. emiss.}} &= \pi \sigma \omega_k^L \frac{\omega_{pe}^2}{k} \int dv \, \delta(\sigma \omega_k^L - kv) \frac{\partial F_e}{\partial v} I_k^{\sigma L}, \\ \frac{\partial}{\partial t} \bigg|_{\text{ind. emiss.}} \frac{I_k^{\sigma S}}{\mu_k} &= \pi \mu_k \sigma \omega_k^L \frac{\omega_{pe}^2}{k} \int dv \, \delta(\sigma \omega_k^S - kv) \frac{\partial}{\partial v} \\ & \times \left(F_e + \frac{m_e}{m_i} F_i\right) \frac{I_k^{\sigma S}}{\mu_k}, \\ \mu_{\mathbf{k}} &= k^3 \lambda_{De}^3 \sqrt{\frac{m_e}{m_i}} \sqrt{1 + \frac{3T_i}{T_e}}. \end{aligned}$$







Decay/Coalescence: Nonlinear Three-Wave Resonance

$$\begin{split} \frac{\partial I_{k}^{\sigma L}}{\partial t} \bigg|_{\text{decay}} &= \frac{\pi}{2} \frac{e^{2}}{T_{e}^{2}} \sigma \omega_{k}^{L} \sum_{\sigma',\sigma''=\pm 1} \int dk' \frac{\mu_{k-k'}}{(k-k')^{2}} \\ &\times \left[\sigma \omega_{k}^{L} I_{k'}^{\sigma' L} \frac{I_{k-k'}^{\sigma'' S}}{\mu_{k-k'}} - \left(\sigma' \omega_{k'}^{L} \frac{I_{k-k'}^{\sigma'' S}}{\mu_{k-k'}} + \sigma'' \omega_{k-k'}^{L} I_{k'}^{\sigma' L} \right) I_{k}^{\sigma L} \right] \\ &\times \delta(\sigma \omega_{k}^{L} - \sigma' \omega_{k'}^{L} - \sigma'' \omega_{k-k'}^{S}), \end{split}$$







$$\begin{split} & \frac{\partial}{\partial t} \bigg|_{\text{decay}} \frac{I_k^{\sigma S}}{\mu_k} = \frac{\pi}{4} \frac{e^2}{T_e^2} \sigma \omega_k^S \sum_{\sigma', \sigma''=\pm 1} \int dk' \frac{\mu_k}{k^2} \\ & \times \left[\sigma \omega_k^L I_{k'}^{\sigma'L} I_{k-k'}^{\sigma''L} - \left(\sigma' \omega_{k'}^L I_{k-k'}^{\sigma''L} + \sigma'' \omega_{k-k'}^L I_{k'}^{\sigma'L} \right) \frac{I_k^{\sigma S}}{\mu_k} \right] \\ & \times \delta(\sigma \omega_k^S - \sigma' \omega_{k'}^L - \sigma'' \omega_{k-k'}^L). \end{split}$$







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Induced Scattering: Nonlinear Wave-Particle Resonance

$$\begin{aligned} \frac{\partial I_{k}^{\sigma L}}{\partial t} \bigg|_{\text{scatt.}} &= \sigma \omega_{k}^{L} \frac{\pi}{\omega_{pe}^{2}} \frac{e^{2}}{m_{e} m_{i}} \sum_{\sigma'=\pm 1} \int dk' \int dv \\ &\times (k-k') \frac{\partial F_{i}}{\partial v} \,\delta[\sigma \omega_{k}^{L} - \sigma' \omega_{k'}^{L} - (k-k') \,v] \,I_{k'}^{\sigma' L} \,I_{k}^{\sigma L}. \end{aligned}$$







Adding effects of spontaneous thermal fluctuation

This tutorial is too short to discuss in detail but in general, induced processes must be balanced by spontaneous processes.

$$\begin{aligned} \frac{\partial F_{e}}{\partial t} &= \frac{\partial}{\partial v_{i}} \left(A_{i} F_{e} + D_{ij} \frac{\partial F_{e}}{\partial v_{j}} \right), \\ A_{i} &= \frac{e^{2}}{4\pi m_{e}} \int d\mathbf{k} \frac{k_{i}}{k^{2}} \sum_{\sigma=\pm 1} \sigma \omega_{\mathbf{k}}^{L} \,\delta(\sigma \omega_{\mathbf{k}}^{L} - \mathbf{k} \cdot \mathbf{v}), \\ D_{ij} &= \frac{\pi e^{2}}{m_{e}^{2}} \int d\mathbf{k} \frac{k_{i} k_{j}}{k^{2}} \sum_{\sigma=\pm 1} \delta(\sigma \omega_{\mathbf{k}}^{L} - \mathbf{k} \cdot \mathbf{v}) \, I_{\mathbf{k}}^{\sigma L}. \end{aligned}$$



Langmuir Wave Kinetic Equation

$$\frac{\partial l_{\mathbf{k}}^{\sigma L}}{\partial t} = \frac{\pi \omega_{pe}^{2}}{k^{2}} \int d\mathbf{v} \, \delta(\sigma \omega_{\mathbf{k}}^{L} - \mathbf{k} \cdot \mathbf{v}) \begin{pmatrix} \hat{n} e^{2} \\ \pi \end{pmatrix} F_{e} + \sigma \omega_{\mathbf{k}}^{L} l_{\mathbf{k}}^{\sigma L} \mathbf{k} \cdot \frac{\partial F_{e}}{\partial \mathbf{v}} \end{pmatrix}$$
spont. emission induced emission
$$+ 2 \sum_{\sigma',\sigma''=\pm 1} \sigma \omega_{\mathbf{k}}^{L} \int d\mathbf{k}' \frac{\pi}{4} \frac{e^{2}}{T_{e}^{2}} \frac{\mu_{\mathbf{k}-\mathbf{k}'} (\mathbf{k} \cdot \mathbf{k}')^{2}}{k^{2} k'^{2} |\mathbf{k} - \mathbf{k}'|^{2}}$$

$$\times \delta(\sigma \omega_{\mathbf{k}}^{L} - \sigma' \omega_{\mathbf{k}'}^{L} - \sigma'' \omega_{\mathbf{k}-\mathbf{k}'}^{S})$$

$$\times \left(\underbrace{\sigma \omega_{\mathbf{k}}^{L} l_{\mathbf{k}'}^{\sigma' L} l_{\mathbf{k}-\mathbf{k}'}^{\sigma'' S}}_{\text{spont. decay}} - \sigma' \omega_{\mathbf{k}'}^{L} l_{\mathbf{k}-\mathbf{k}'}^{\sigma' L} l_{\mathbf{k}}^{\sigma L} - \sigma'' \omega_{\mathbf{k}-\mathbf{k}'}^{L} l_{\mathbf{k}'}^{\sigma' L} l_{\mathbf{k}'}^{\sigma L} \right)$$

$$-\frac{\pi e^{2}}{m_{e}^{2}\omega_{pe}^{2}}\sum_{\sigma'=\pm 1}\int d\mathbf{k}'\int d\mathbf{v} \frac{(\mathbf{k}\cdot\mathbf{k}')^{2}}{k^{2}k'^{2}}\delta[\sigma\omega_{\mathbf{k}}^{L}-\sigma'\omega_{\mathbf{k}'}^{L}-(\mathbf{k}-\mathbf{k}')\cdot\mathbf{v}]$$

$$\times \left[\frac{\hat{n}e^{2}}{\pi\omega_{pe}^{2}}\sigma\omega_{\mathbf{k}}^{L}\left(\sigma'\omega_{\mathbf{k}'}^{L}I_{\mathbf{k}}^{\sigma L}-\sigma\omega_{\mathbf{k}}^{L}I_{\mathbf{k}'}^{\sigma'L}\right)(F_{e}+F_{i})\right]$$
spontaneous scattering
$$+I_{\mathbf{k}'}^{\sigma'L}I_{\mathbf{k}}^{\sigma L}(\mathbf{k}-\mathbf{k}')\cdot\frac{\partial}{\partial\mathbf{v}}\left((\sigma\omega_{\mathbf{k}}^{L}-\sigma'\omega_{\mathbf{k}'}^{L})F_{e}-\frac{m_{e}}{m_{i}}(\sigma\omega_{\mathbf{k}}^{L})F_{i}\right)\right].$$
induced scattering

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Forward/backward-Ion-sound Wave Kinetic Equation

$$\begin{aligned} \frac{\partial l_{\mathbf{k}}^{\sigma S}}{\partial t} &= \frac{\pi \mu_{\mathbf{k}} \omega_{pe}^{2}}{k^{2}} \int d\mathbf{v} \ \delta(\sigma \omega_{\mathbf{k}}^{S} - \mathbf{k} \cdot \mathbf{v}) \\ \times \left[\underbrace{\frac{\hat{n} e^{2}}{\pi} \left(F_{e} + F_{i}\right)}_{\pi} + \underbrace{\sigma \omega_{\mathbf{k}}^{L} l_{\mathbf{k}}^{\sigma S} \left(\mathbf{k} \cdot \frac{\partial}{\partial \mathbf{v}}\right) \left(F_{e} + \frac{m_{e}}{m_{i}} F_{i}\right)}_{\text{spont. emission}} \right] \\ \text{spont. emission} \quad \text{induced emission} \\ \sum_{\sigma', \sigma'' = \pm 1} \sigma \omega_{\mathbf{k}}^{L} \int d\mathbf{k}' \ \frac{\pi}{4} \frac{e^{2}}{T_{e}^{2}} \frac{\mu_{\mathbf{k}} \left[\mathbf{k}' \cdot (\mathbf{k} - \mathbf{k}')\right]^{2}}{k^{2} k'^{2} |\mathbf{k} - \mathbf{k}'|^{2}} \ \delta(\sigma \omega_{\mathbf{k}}^{S} - \sigma' \omega_{\mathbf{k}'}^{L} - \sigma'' \omega_{\mathbf{k} - \mathbf{k}'}^{L}) \\ \times \left(\underbrace{\sigma \omega_{\mathbf{k}}^{L} l_{\mathbf{k}'}^{\sigma'L} l_{\mathbf{k} - \mathbf{k}'}^{\sigma''L}}_{\text{spont. decay}} - \underbrace{\sigma' \omega_{\mathbf{k}'}^{L} l_{\mathbf{k} - \mathbf{k}'}^{\sigma S} - \sigma'' \omega_{\mathbf{k} - \mathbf{k}'}^{L} l_{\mathbf{k}'}^{\sigma S}}_{\text{spont. decay}} \right). \end{aligned}$$

Quasi-stationary ions

$$F_i = rac{e^{-v^2/v_{T_i}^2}}{\pi^{1/2} v_{T_i}}.$$

Initial beam plus background electron distribution

$$F_e(v,0) = \frac{(1-\delta) e^{-v^2/v_{Te}^2}}{\pi^{1/2} v_{Te}} + \frac{\delta e^{-(v-v_0)^2/v_{Te}^2}}{\pi^{1/2} v_{Te}}.$$





Beam-Plasma Instability: Quasi-Linear Saturation Regime





Beam-Plasma Instability: Nonlinear Mode-Coupling Regime





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Beam-Plasma Instability: Nonlinear Mode-Coupling Regime





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Formation of energetic tail in the near the end of nonlinear regime (kappa VDF?)





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Physica Scripta. Vol. 59, 19-26, 1999

Kinetic Theoretical Foundation of Lorentzian Statistical Mechanics

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$$\begin{split} \frac{\partial f_e}{\partial t} &= \frac{\partial}{\partial v_i} \left(A_i f_e + D_{ij} \frac{\partial f_e}{\partial v_j} \right), \\ A_i &= \frac{e^2}{4\pi m_e} \int d\mathbf{k} \frac{k_i}{k^2} \sum_{\sigma=\pm 1} \sigma \omega_{\mathbf{k}}^L \delta(\sigma \omega_{\mathbf{k}}^L - \mathbf{k} \cdot \mathbf{v}), \\ D_{ij} &= \frac{\pi e^2}{m_e^2} \int d\mathbf{k} \frac{k_i k_j}{k^2} \sum_{\sigma=\pm 1} \delta(\sigma \omega_{\mathbf{k}}^L - \mathbf{k} \cdot \mathbf{v}) I_{\mathbf{k}}^{\sigma L}. \end{split}$$

Asymptotic solution ($\partial/\partial t = 0$)

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Steady-State Solution (Quasi-Equilibrium)

Electron kinetic equation

$$\begin{split} 0 &= \frac{\partial}{\partial v_i} \left(v_i G f_e + D \frac{v_i v_j}{v^2} \frac{\partial f_e}{\partial v_j} \right), \\ G &= \frac{e^2 \omega_{pe}^2}{4 \pi m_e v^2} \int \frac{d \mathbf{k}}{k^2} \delta(\omega_{pe} - \mathbf{k} \cdot \mathbf{v}), \\ D &= \frac{\pi e^2 \omega_{pe}^2}{m_e^2 v^2} \int \frac{d \mathbf{k}}{k^2} \delta(\omega_{pe} - \mathbf{k} \cdot \mathbf{v}) I(\mathbf{k}). \end{split}$$

Steady-state solution [Gurevich, 1960; Hasegawa et al., 1985]

$$f_e = C \exp\left(-\int dv \frac{vG}{D}\right).$$

Steady-State Wave Equation

$$0 = \frac{\pi \omega_{pe}^{2}}{k^{2}} \int d\mathbf{v} \delta(\omega_{k} - \mathbf{k} \cdot \mathbf{v}) \left(\frac{ne^{2}}{\pi} f_{e} + \omega_{pe} I(\mathbf{k}) \mathbf{k} \cdot \frac{\partial f_{e}}{\partial \mathbf{v}} \right)$$
$$- \frac{\omega_{pe}}{4\pi n T_{i}} \int d\mathbf{k}' \int d\mathbf{v} \delta[\omega_{k} - \omega_{k'} - (\mathbf{k} - \mathbf{k}') \cdot \mathbf{v}]$$
$$\times \left(\frac{T_{i}}{4\pi^{2}} [\omega_{k'} I(\mathbf{k}) - \omega_{k} I(\mathbf{k}')] + I(\mathbf{k}) I(\mathbf{k}') (\omega_{k} - \omega_{k'}) \right) f_{i}.$$



Balance of spontaneous and induced emissions $0 = \frac{\pi \omega_{pe}^2}{k^2} \int d\mathbf{v} \delta(\omega_k - \mathbf{k} \cdot \mathbf{v}) \left(\frac{ne^2}{\pi} f_e + \omega_{pe} I(\mathbf{k}) \mathbf{k} \cdot \frac{\partial f_e}{\partial \mathbf{v}} \right) = \mathbf{0}$ $-\frac{\omega_{pe}}{4\pi nT_{i}}\int d\mathbf{k}'\int d\mathbf{v}\delta[\omega_{k}-\omega_{k'}-(\mathbf{k}-\mathbf{k}')\cdot\mathbf{v}]$ $\times \left(\frac{T_{i}}{4\pi^{2}}[\omega_{k'}I(\mathbf{k})-\omega_{k}I(\mathbf{k}')]+I(\mathbf{k})I(\mathbf{k}')(\omega_{k}-\omega_{k'})\right)f_{i}.$ =0

Balance of spontaneous and induced scatterings



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Kappa Distribution

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Balance of particle equation and linear wave equation leads to

$$f_e(\mathbf{v}) \propto \left[1 + \frac{\mathbf{v}^2}{\left(\kappa - \frac{3}{2}\right)\mathbf{v}_{Te}^2}\right]^{-\kappa - 1},$$

$$I(k) = \frac{T_e}{4\pi^2} \frac{\kappa - \frac{3}{2}}{\kappa + 1} \left(1 + \frac{\omega_{pe}^2}{\left(\kappa - \frac{3}{2}\right)\left(k\mathbf{v}_{Te}\right)^2}\right).$$

Balance of nonlinear wave equation separately leads to

$$I(k) = \frac{T_i}{4\pi^2} \left(1 + \frac{4}{3} \frac{\omega_{pe}^2}{k^2 v_{Te}^2} \right).$$





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The two alternative expressions for the wave intensity,

$$\begin{split} I(k) &= \frac{T_e}{4\pi^2} \frac{\kappa - \frac{3}{2}}{\kappa + 1} \left(1 + \frac{\omega_{pe}^2}{\left(\kappa - \frac{3}{2}\right) \left(k v_{Te}\right)^2} \right) \qquad \text{and} \\ I(k) &= \frac{T_i}{4\pi^2} \left(1 + \frac{4}{3} \frac{\omega_{pe}^2}{k^2 v_{Te}^2} \right), \end{split}$$

must be reconciled, which is possible for

$$T_e \frac{\kappa - 3/2}{\kappa + 1} = T_i, \quad \text{and} \quad \kappa - \frac{3}{2} = \frac{3}{4},$$
K()



or equivalently,

$$\kappa = \frac{9}{4} = 2.25,$$
 or $q = \frac{\kappa - 1}{\kappa} = \frac{5}{9}.$

To test this finding, we turn to the solar wind.







- P. H. Yoon, T. Rhee & C.-M. Ryu, Self-consistent generation of superthermal electrons by beam-plasma interaction, PRL, 95, 215003 (2005)
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Quiet Time Solar Wind Electrons





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SOLAR WIND ELECTRON DISTRIBUTION AT 1 AU



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- Kappa distributions may imply non-extensive statistical principles dictating the behavior of complex systems, and systems interacting with long-range force such as plasmas.
- Turbulent equilibrium describes a system of plasma particles constantly exchanging momentum and energy with turbulence, while maintaining a dynamical steady-state. Such a state happens to correspond to kappa distribution.
- Conclusion: turbulent equilibrium state for plasmas may be equivalent to the non-extensive equilibrium.



- The quiet-time solar wind also features inverse power-law proton distribution $f(v) \propto v^{-\alpha}$, where $\alpha \sim 5$ [Gloeckler & Fisk, Fisk & Gloeckler, ...].
- There is a permanent low-frequency solar wind turbulence of the kinetic Alfvénic variety.
- Question: Are the two dynamically coupled? In other words, are the solar wind protons and kinetic Alfvénic turbulence in dynamical equilibrium?















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