

Simulations of flux emergence events I

Daniel Nóbrega-Siverio^{1,2,3,4}

1. Instituto de Astrofísica de Canarias, E-38205 La Laguna, Tenerife, Spain.
2. Universidad de La Laguna, Dept. Astrofísica, E-38206 La Laguna, Tenerife, Spain
3. Rosseland Centre for Solar Physics, University of Oslo, PO Box 1029 Blindern, 0315 Oslo, Norway.
4. Institute of Theoretical Astrophysics, University of Oslo, PO Box 1029 Blindern, 0315 Oslo, Norway.

Email: dnobrega@iac.es



Rosseland
Centre
for Solar
Physics

Introduction



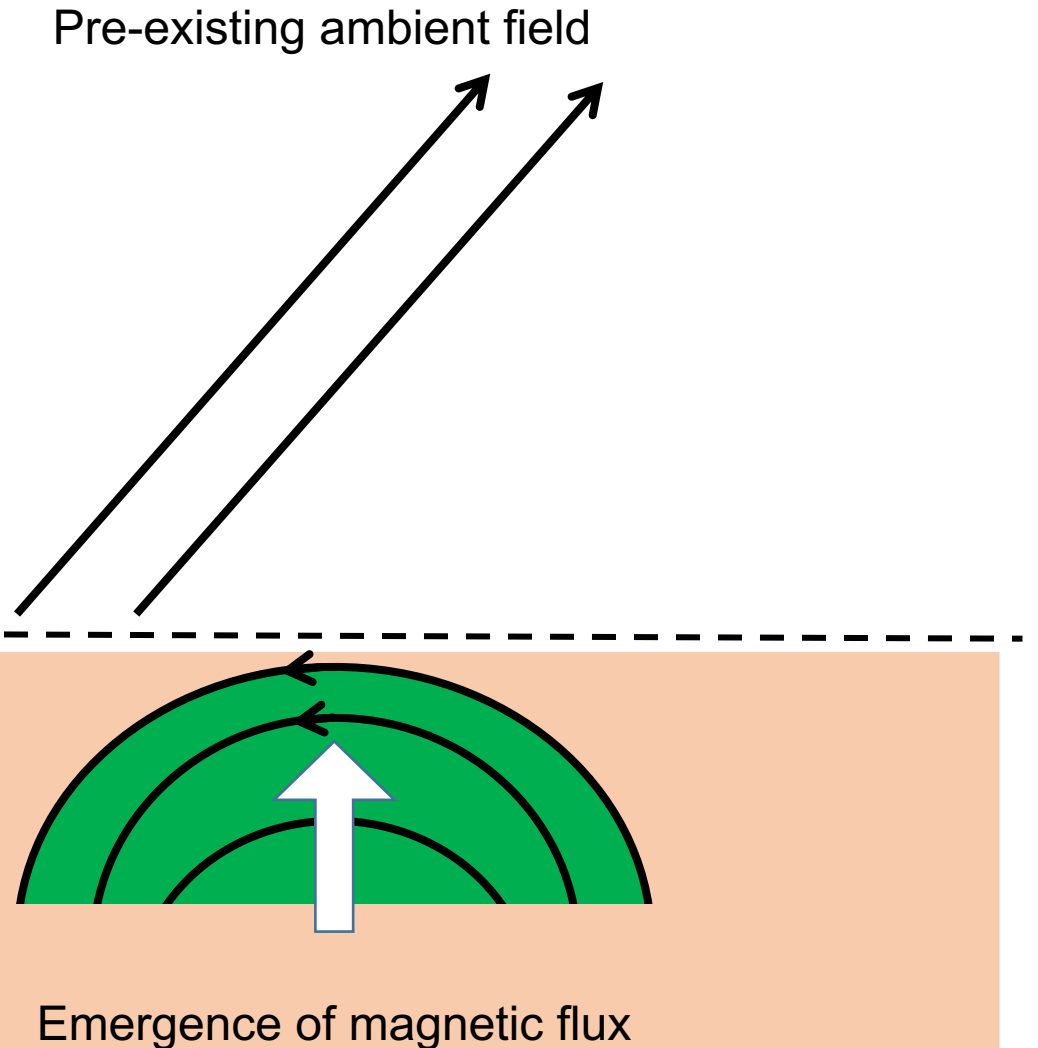
Rosseland
Centre
for Solar
Physics

In these two lectures:

- We are going to study how magnetized plasma goes from the upper layers of the convection zone up to the solar atmosphere.

Solar atmosphere

Convection zone



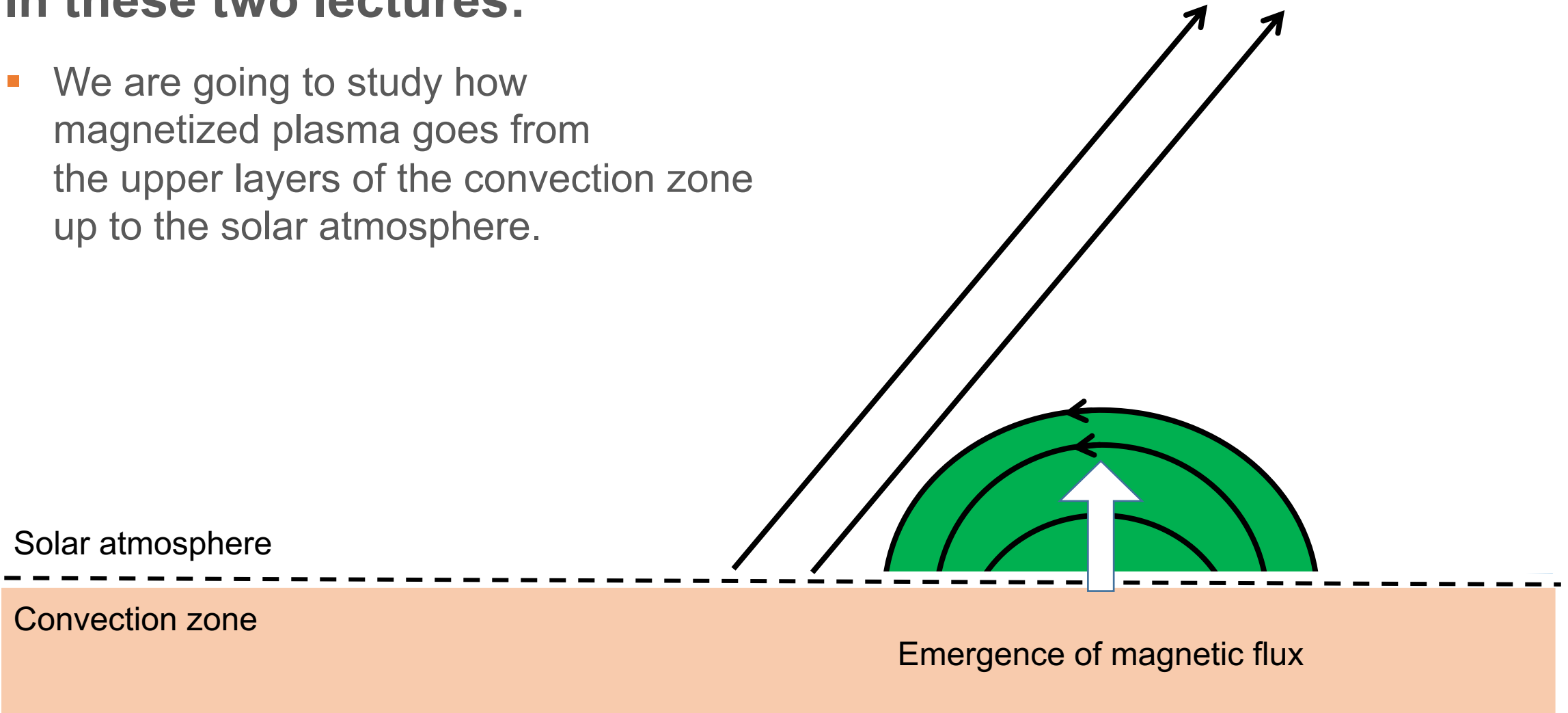
Introduction



In these two lectures:

- We are going to study how magnetized plasma goes from the upper layers of the convection zone up to the solar atmosphere.

Pre-existing ambient field



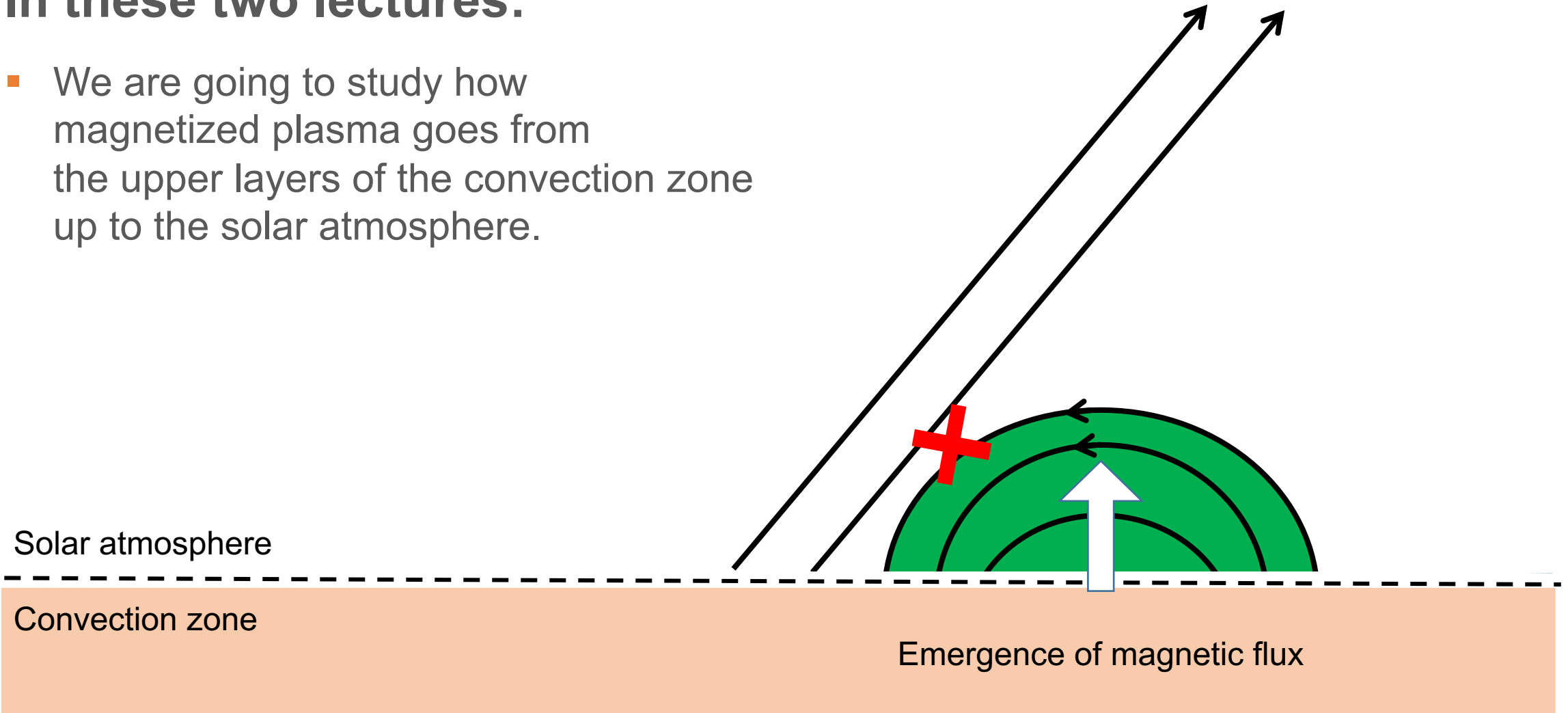
Introduction



In these two lectures:

- We are going to study how magnetized plasma goes from the upper layers of the convection zone up to the solar atmosphere.

Pre-existing ambient field

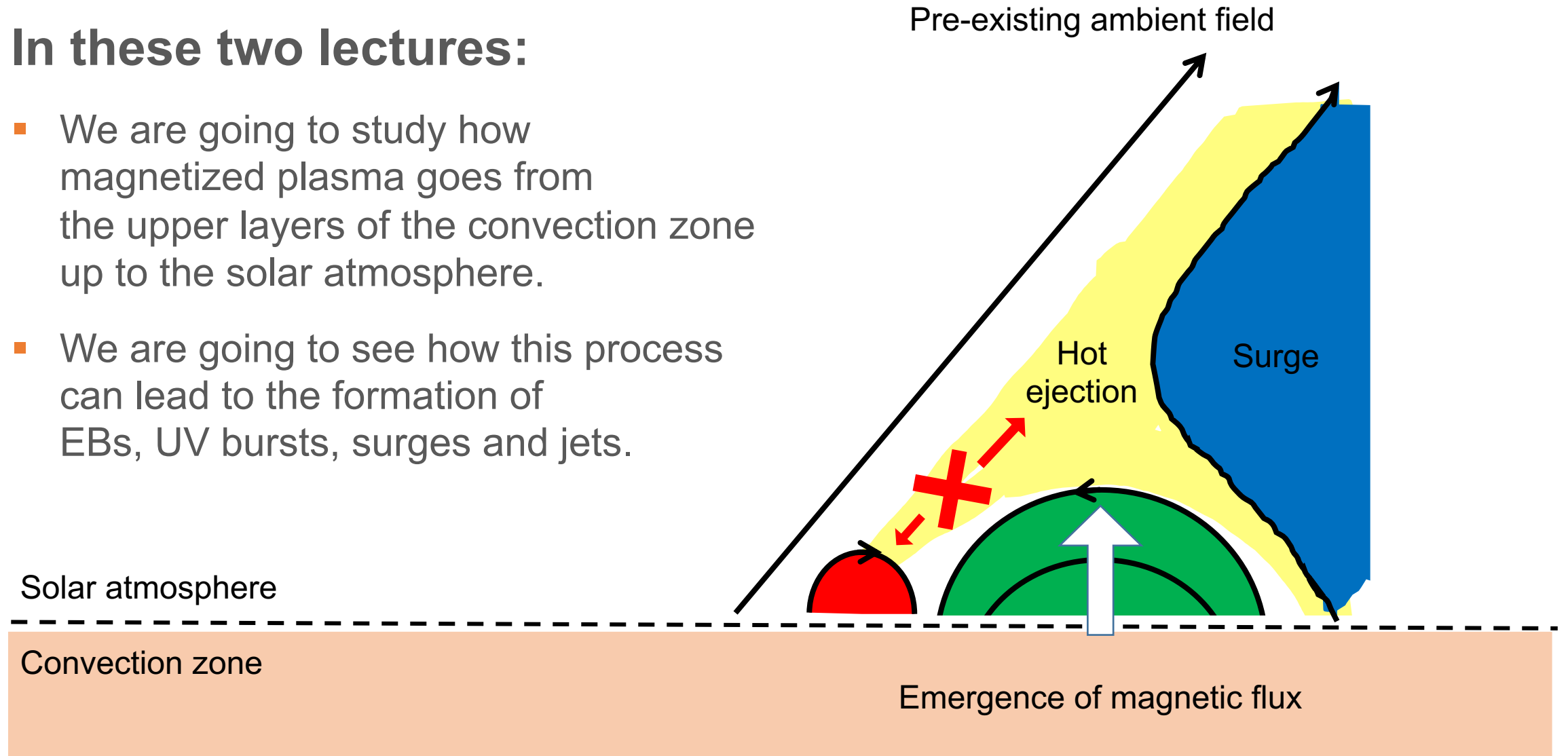


Introduction



In these two lectures:

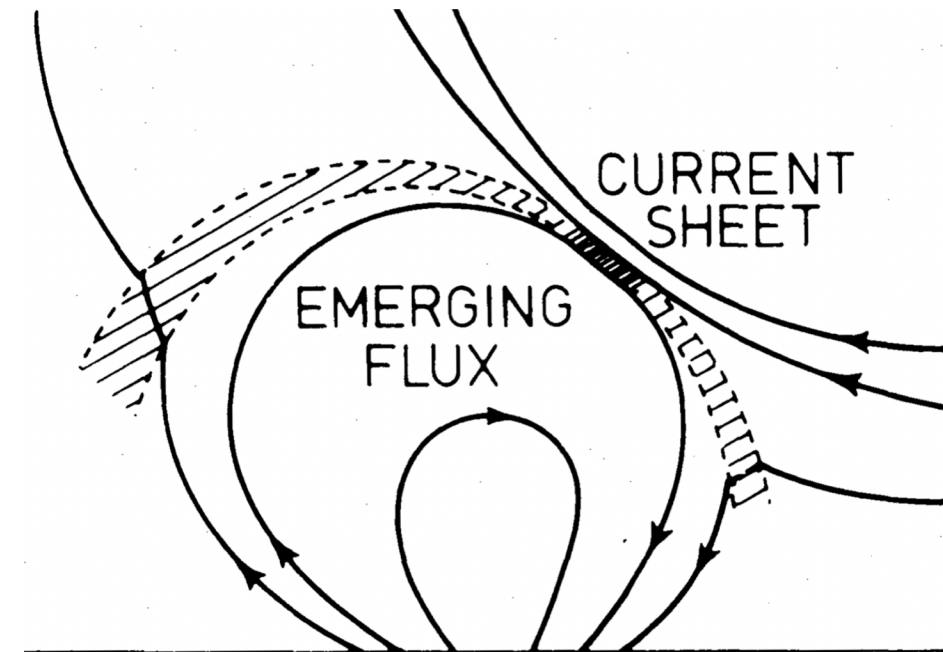
- We are going to study how magnetized plasma goes from the upper layers of the convection zone up to the solar atmosphere.
- We are going to see how this process can lead to the formation of EBs, UV bursts, surges and jets.



Flux emergence models:

- The first theoretical explanation came in the 70s by Heyvaerts et al. 1977: solar flares could be produced by magnetic reconnection between the emerging flux and the pre-existing field.
- Numerous numerical experiments have shown that **magnetic reconnection and flux emergence** are key physical mechanisms to explain eruptive/ejective events in the solar atmosphere

(e.g., Forbes & Priest 1984, Shibata et al. 1992, Yokoyama & Shibata 1995, 1996, Fan 2001, Miyagoshi & Yokoyama 2004, Archontis et al. 2004, 2005, 2007, Leake & Arber 2006, Cheung et al. 2007, 2008, Nishizuka et al. 2008, Moreno-Insertis et al. 2008, Martínez-Sykora et al. 2008, 2009, Murray et al. 2009, Archontis et al. 2010, Takasao et al. 2013, Archontis et al. 2013a,b, Moreno-Insertis & Galsgaard 2013, Fang et al. 2014, Toriumi et al. 2015, MacTaggart et al. 2015, Nóbrega-Siverio et al. 2016, 2017, 2018, Hansteen et al. 2019, Yang et al. 2018, Nóbrega-Siverio, et al. 2020, among many others)



Heyvaerts et al. 1977

Check the reviews by
Cheung & Isobe 2014, Fan 2021

The layout of the lectures is as follows:

- Interaction of magnetized plasma in the upper layers of the convection zone.
- Arrival at the solar surface: anomalous granulation.
- Rise and expansion of the magnetized plasma through the solar atmosphere.
- First stages of the reconnection between emerging and pre-existing field.
- Ellerman bombs, UV bursts, surges and jets.

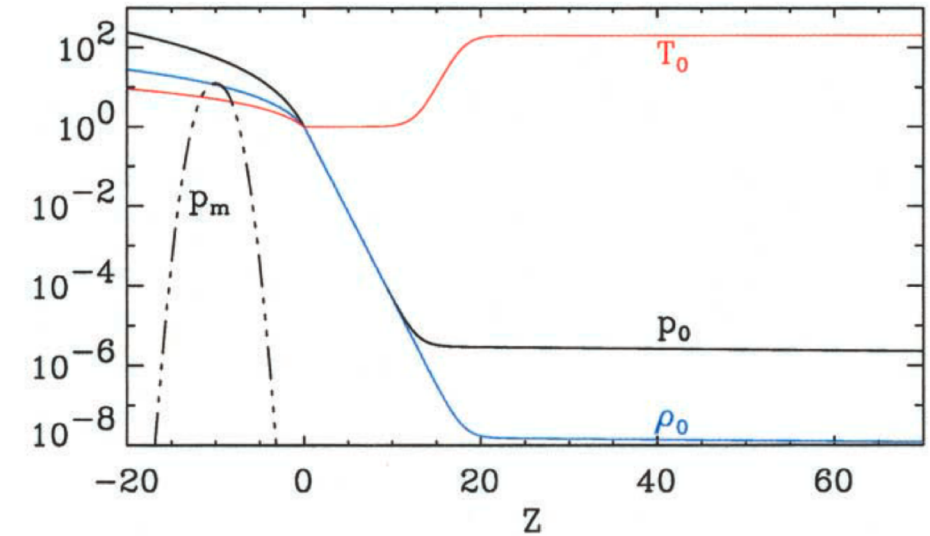
MHD modeling



Initial condition (not including convective motions): :

Plane-parallel atmosphere in hydrostatic Eq.:

- An adiabatically stratified polytropic layer (convection zone).
- An isothermal layer (photosphere and the chromosphere).
- An increasing temperature layer (transition region).
- Another isothermal layer (corona).



Fan 2001

$$T(z) = \begin{cases} 1 - \frac{\gamma-1}{\gamma}z, & z < 0, \\ 1, & 0 \leq z \leq 10, \\ T_{\text{cor}}^{[(z-10)/10]}, & 10 < z < 20, \\ T_{\text{cor}}, & z \geq 20, \end{cases}$$

McTaggart et al. 2015

Twisted magnetic tube:

- Canonical way of having flux emergence. Field components are given by:

$$B_y = B_0 \exp\left(-\frac{r^2}{R_0^2}\right),$$

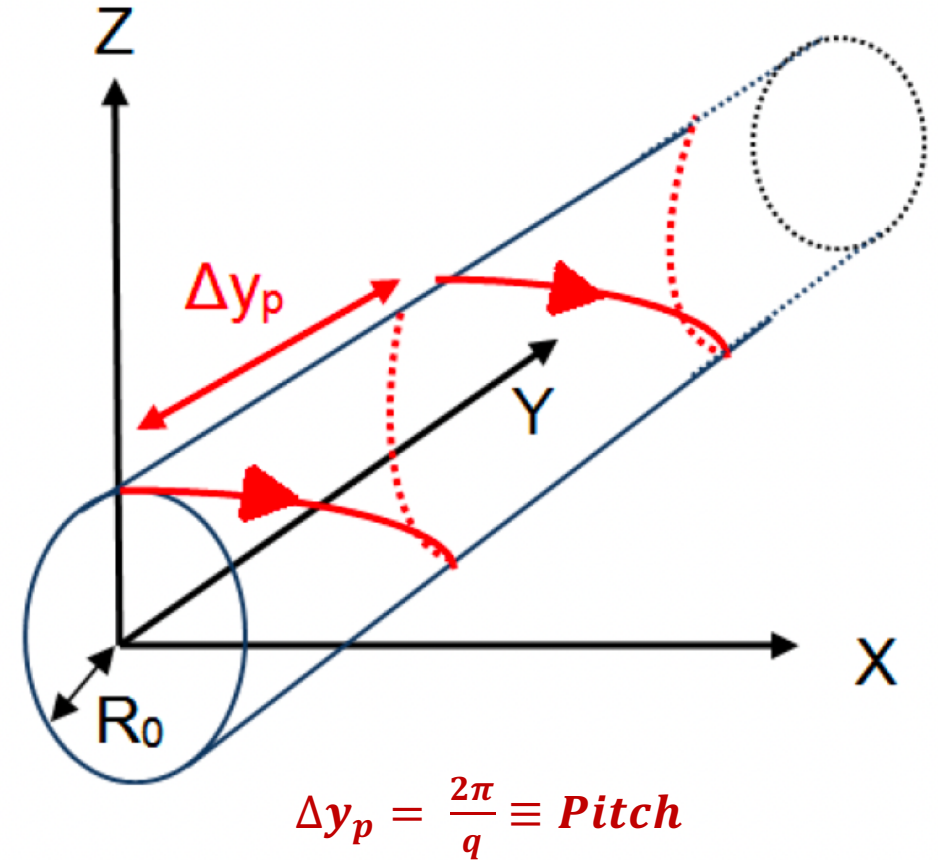
$$B_\theta = q r B_y,$$

where r and θ are the radial and azimuthal coordinates relative to the tube axis,

R_0 is the the radius of the tube,

q is a twist parameter, and

B_0 is the magnetic field in the tube axis.



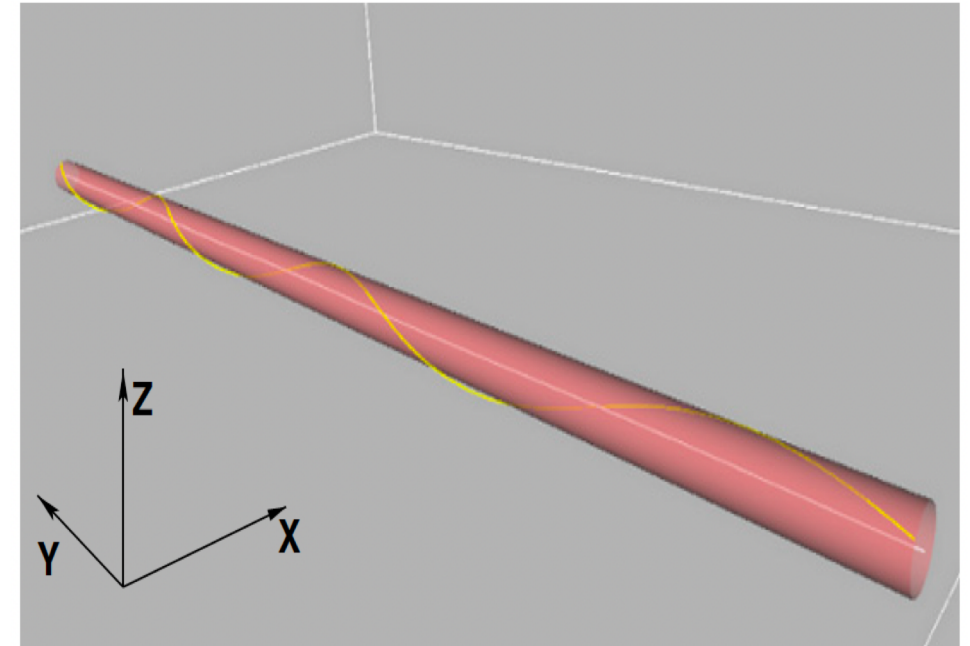
Twisted magnetic tube:

- Canonical way of having flux emergence. Field components are given by:

$$B_y = B_0 \exp\left(-\frac{r^2}{R_0^2}\right),$$

$$B_\theta = q r B_y,$$

where r and θ are the radial and azimuthal coordinates relative to the tube axis, R_0 is the the radius of the tube, q is a twist parameter, and B_0 is the magnetic field in the tube axis.

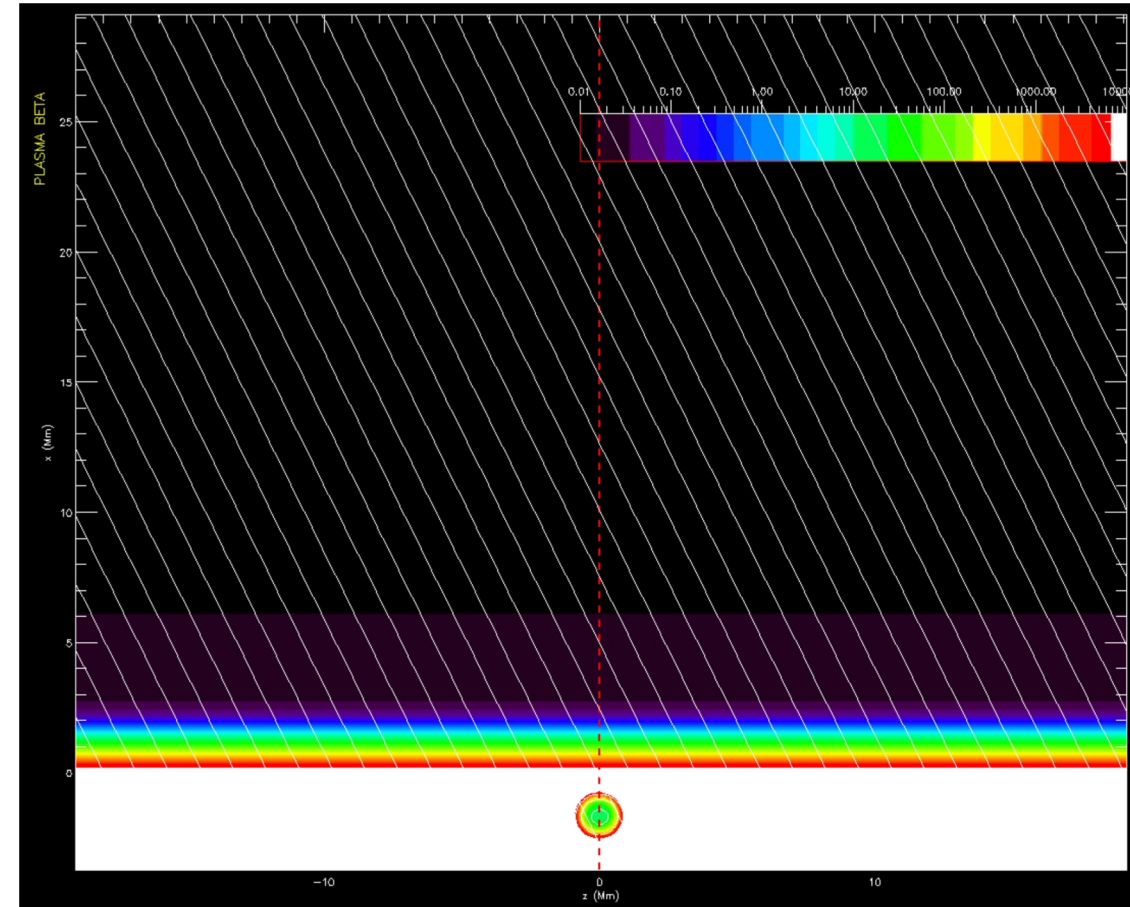


Archontis et al. 2013

Autores	z_0 (Mm)	R_0 (Mm)	q (Mm $^{-1}$)	B_0 (kG)
Moreno-Insertis & Galsgaard (2013)	-1.70	0.430	2.4	3.8
Tortosa-Andreu & Moreno-Insertis (2009)	-1.80	0.400	0.6	7.5
Cheung et al. (2008)	-3.90	0.600	0.4	14.0
Martínez-Sykora et al. (2008)	-2.60	0.450	0.6	4.5
Cheung et al. (2007)	-1.35	0.200	2.6	8.5

Twisted magnetic tube:

- To restore mechanical equilibrium, the gas pressure distribution within the tube is decreased typically via a density deficit.
- This implies that the tube is going to be buoyant, so it can go up through the convection zone.

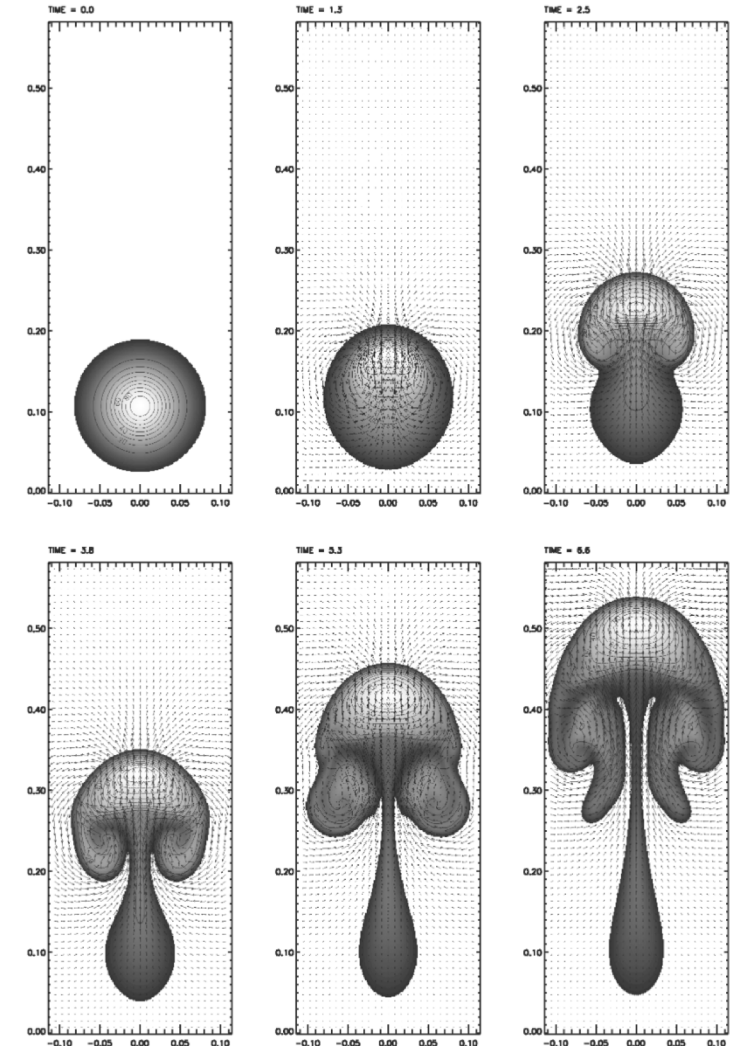




Rise of the tube:

- In a stratified environment without convective motions,
 - the tube tends to split into counter-rotating vortex rolls (dumbbell shapes).
 - the tube can retain coherence because of its transverse component counteracts the horizontal pressure gradient.
 - due to the diminishing scale heights, the tube flattens into a pancake-like structure.

(e.g Schüssler 1979, Spruit et al. 1987, Moreno-Insertis & Emonet 1996, Longcope et al. 1996, Toriumi & Yokoyaa 2001, Archontis et al. 2004, among others)



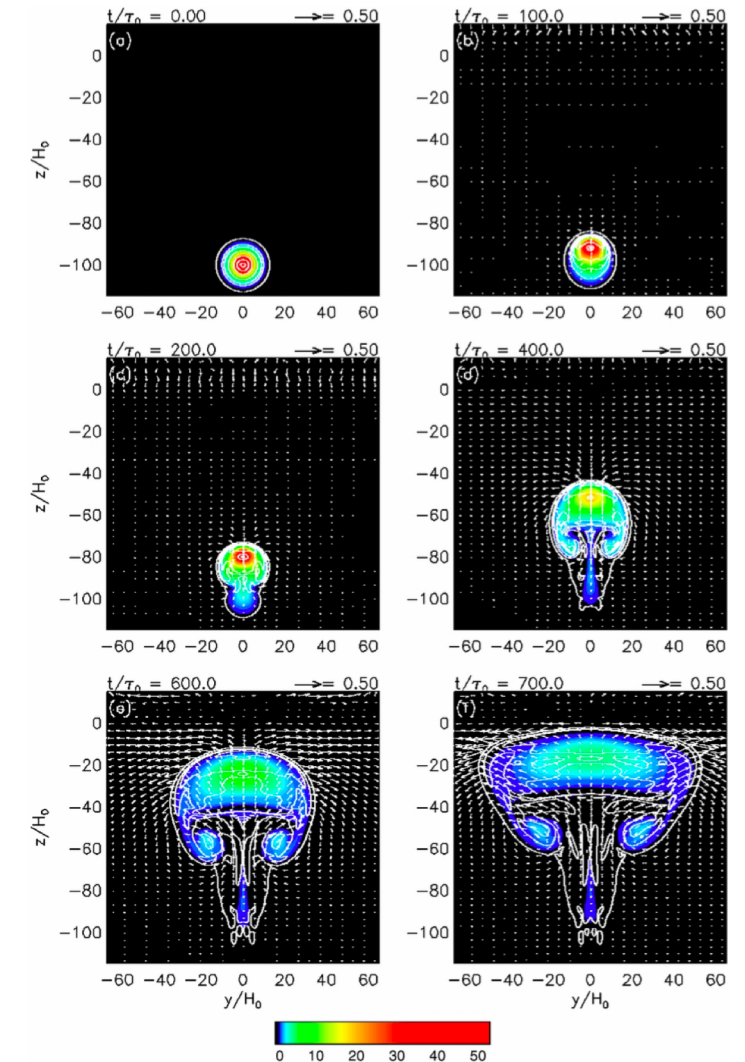
MHD modeling



Rise of the tube:

- In a stratified environment without convective motions,
 - the tube tends to split into counter-rotating vortex rolls (dumbbell shapes).
 - the tube can retain coherence because of its transverse component counteracts the horizontal pressure gradient.
 - due to the diminishing scale heights, the tube flattens into a pancake-like structure.

(e.g Schüssler 1979, Spruit et al. 1987, Moreno-Insertis & Emonet 1996, Longcope et al. 1996, Toriumi & Yokoyama 2001, Archontis et al. 2004, among others)



MHD modeling

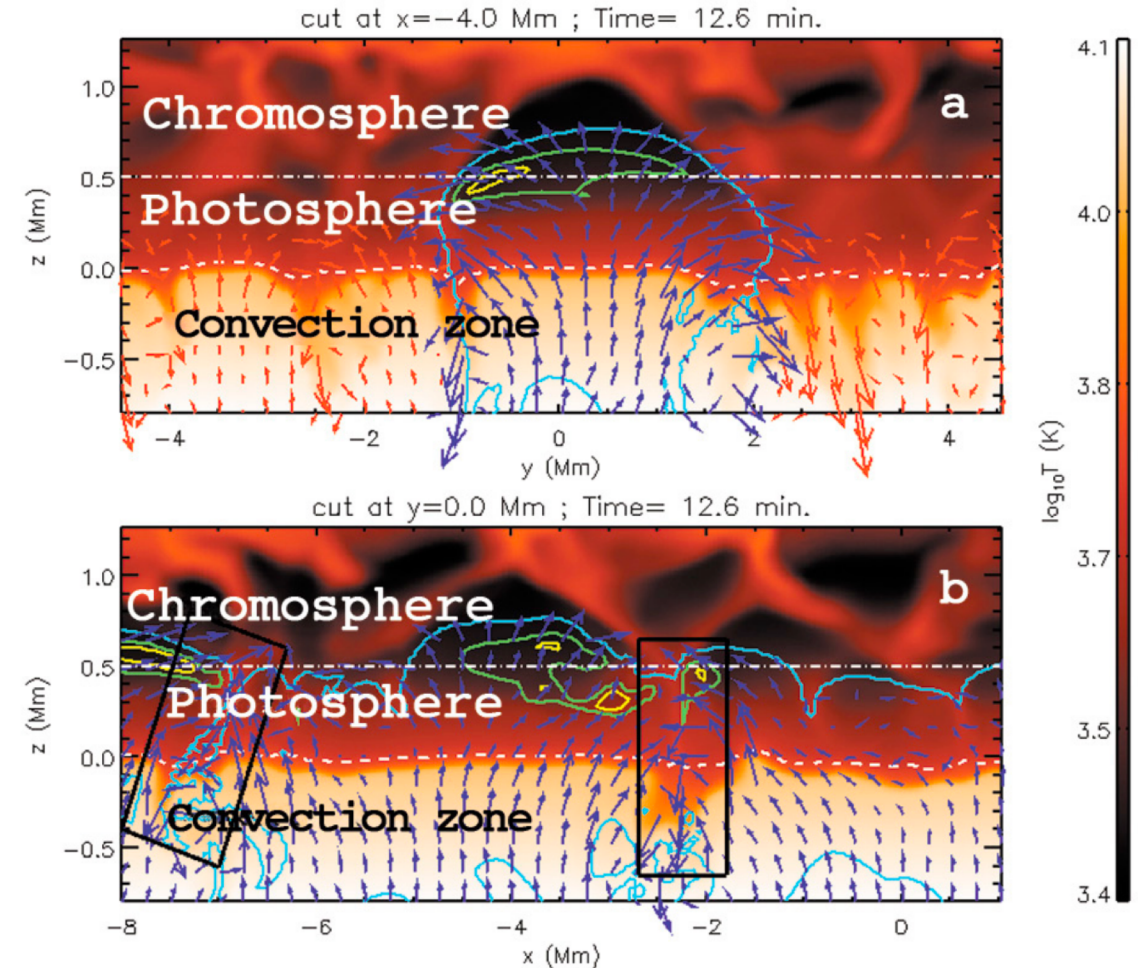


Initial condition (including convective motions):

Statistically Stationary Eq.

- Radiation transfer adequate to the photosphere to get self-consistent convection zone with granulation.

Example using the MURaM code (Vögler et al. 2005)



MHD modeling

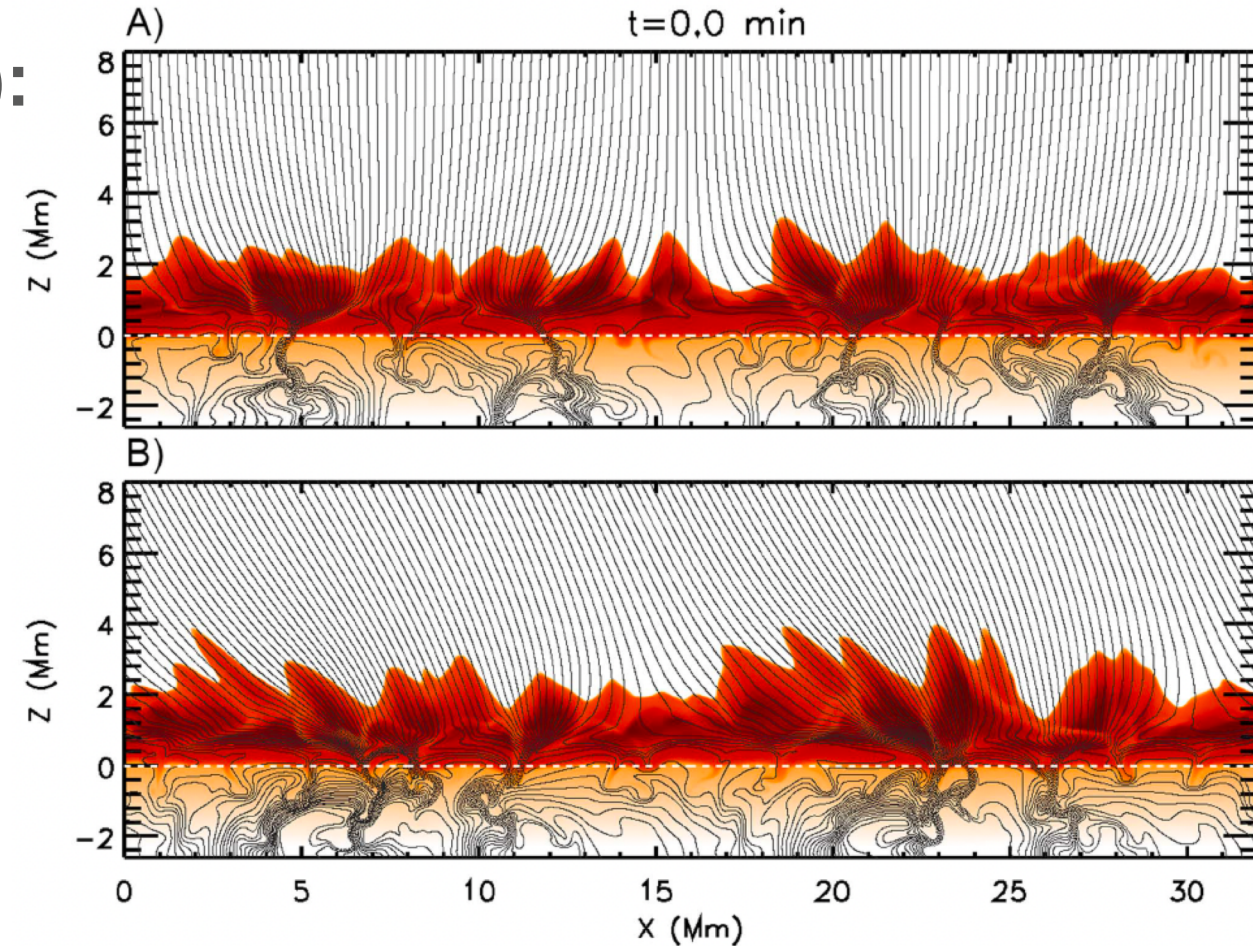


Initial condition (including convective motions):

Statistically Stationary Eq.

- Radiation transfer adequate to the photosphere to get self-consistent convection zone with granulation.

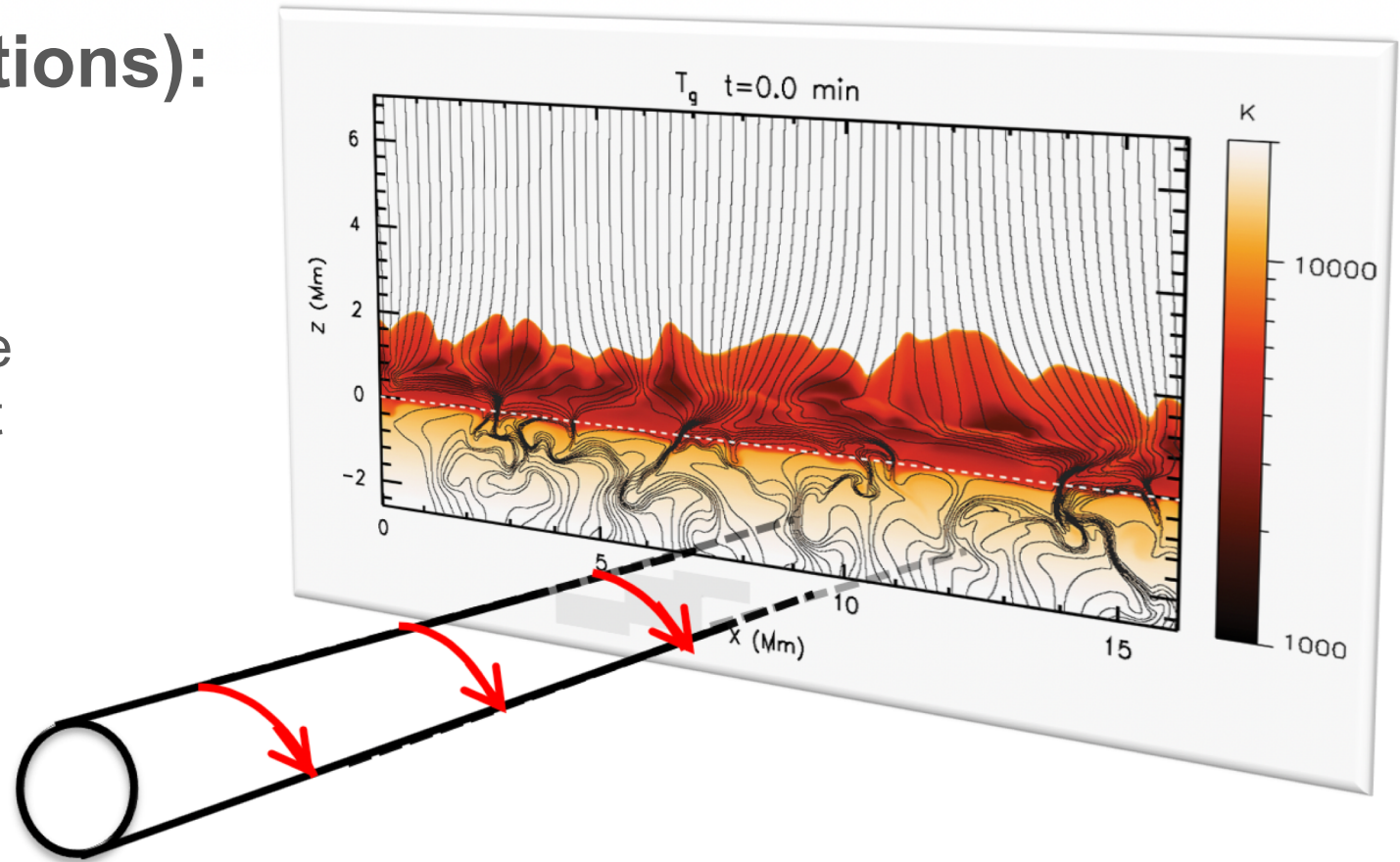
Example using the Bifrost code (Gudiksen et al. 2011)



Initial condition (including convective motions):

Statistically Stationary Eq.

- Radiation transfer adequate to the photosphere to get self-consistent convection zone with granulation.
- Magnetic flux emergence through the injection of a twisted magnetic tube →

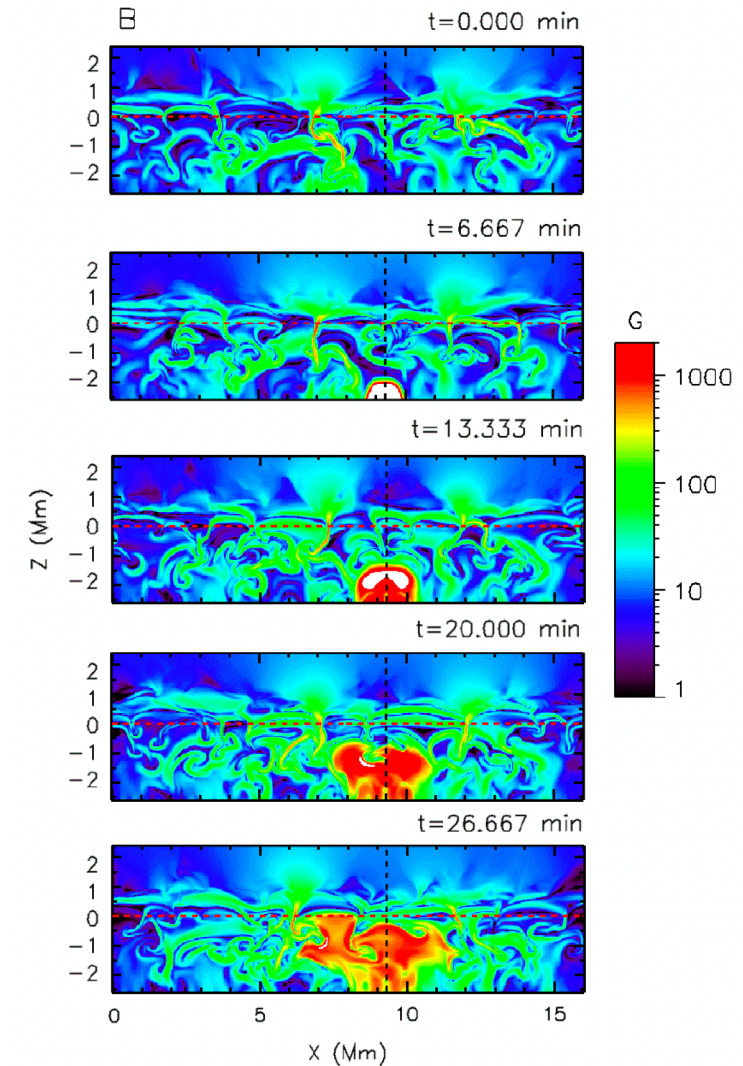


MHD modeling



Rise of the tube:

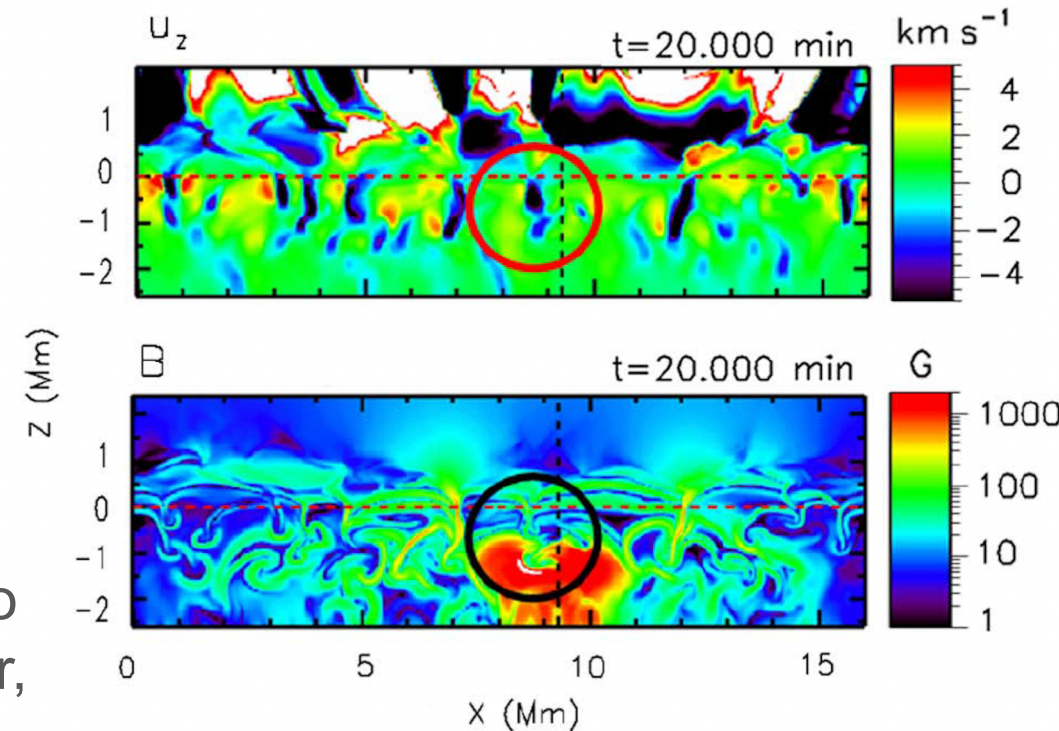
- In a stratified environment **with** convective motions,
 - similar characteristics as we saw before due to the stratification (e.g., vortex rolls),
 - the granulation perturbs the evolution of the twisted magnetic tube,



Rise of the tube:

- In a stratified environment **with** convective motions,
 - similar characteristics as we saw before due to the stratification (e.g., vortex rolls),
 - the granulation perturbs the evolution of the twisted magnetic tube,

in particular,
convection flows deform and break the tube into smaller fragments in the regions of strong shear, typically where the downflows hit the tube.

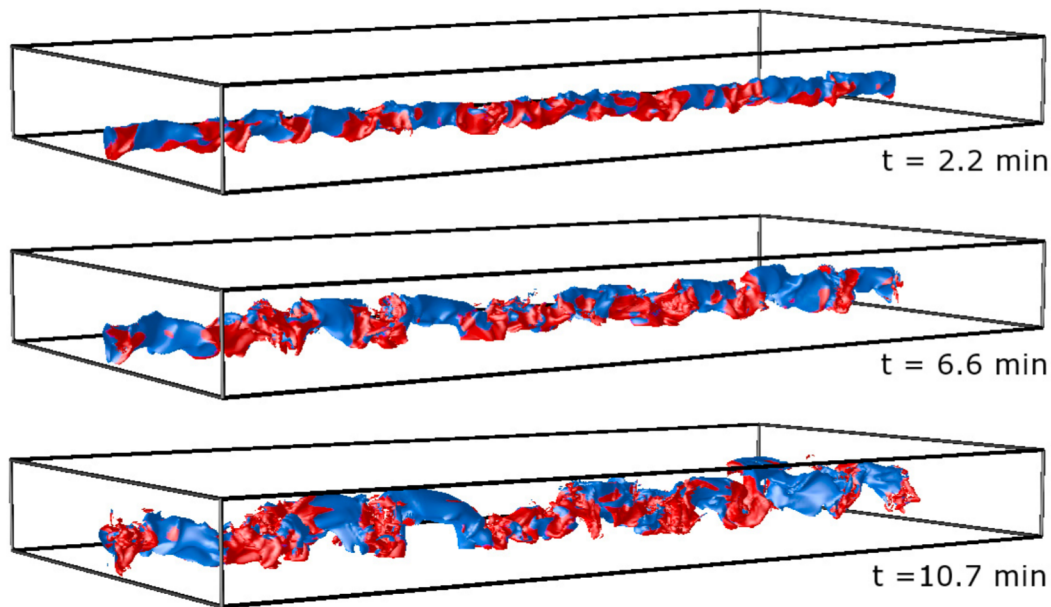


MHD modeling

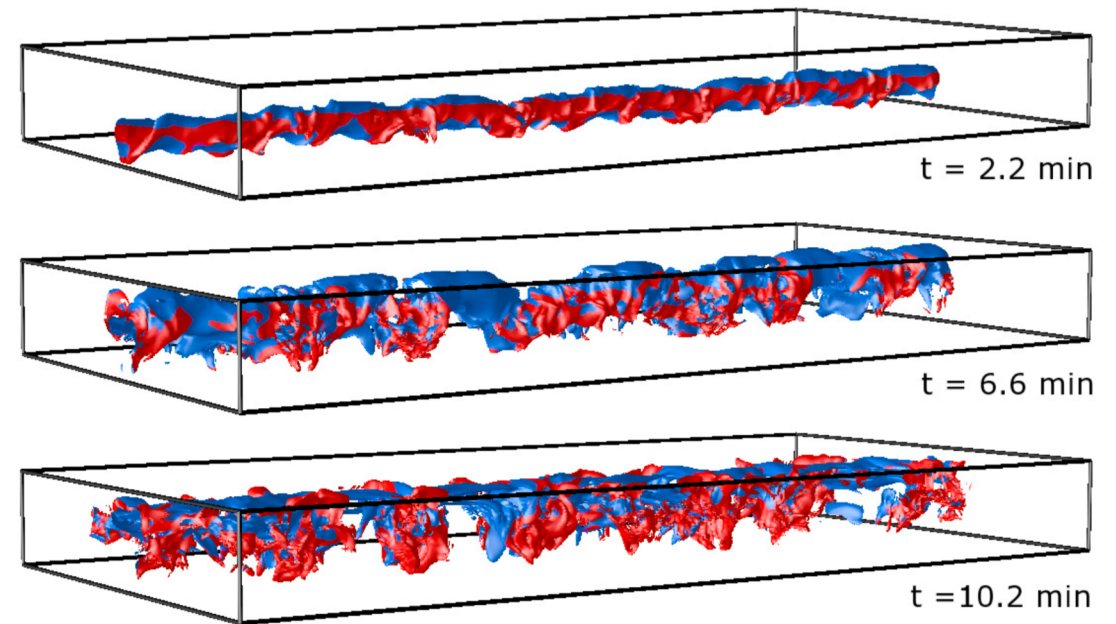


Rise of the tube:

- In a stratified environment **with** convective motions,



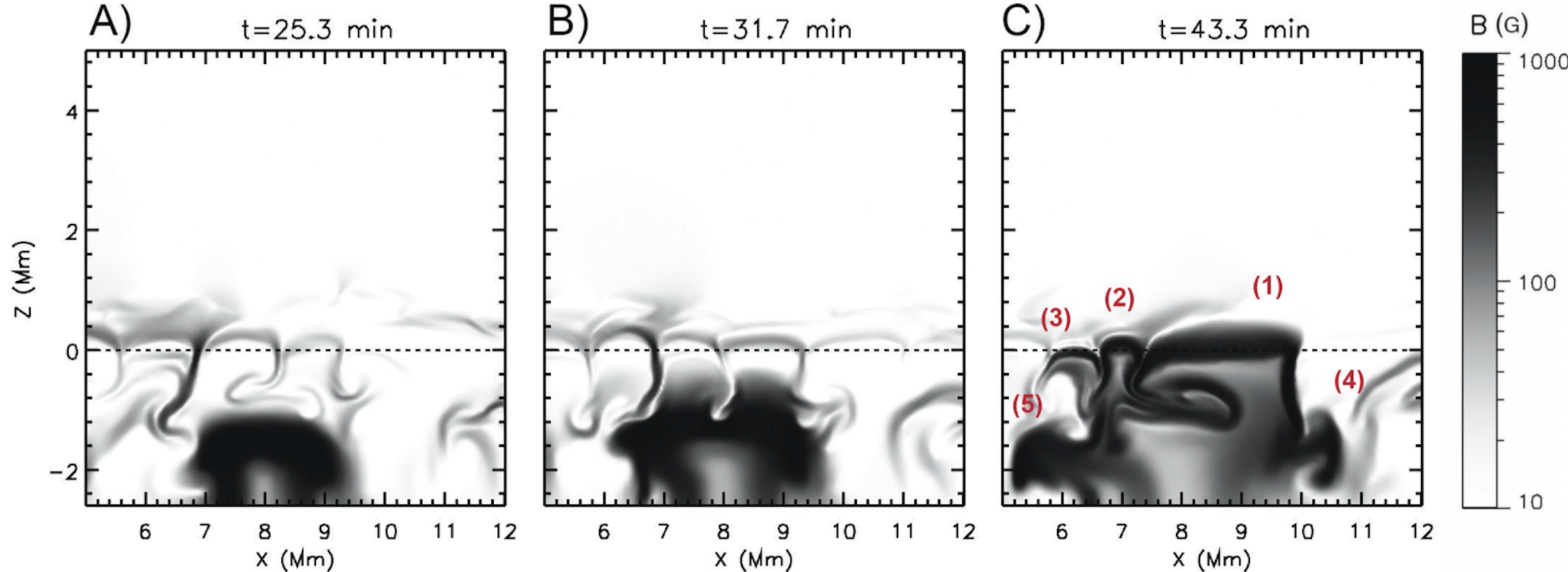
weak magnetic flux tube ($B_0 = 2500$ G)



strong magnetic flux tube ($B_0 = 8500$ G)

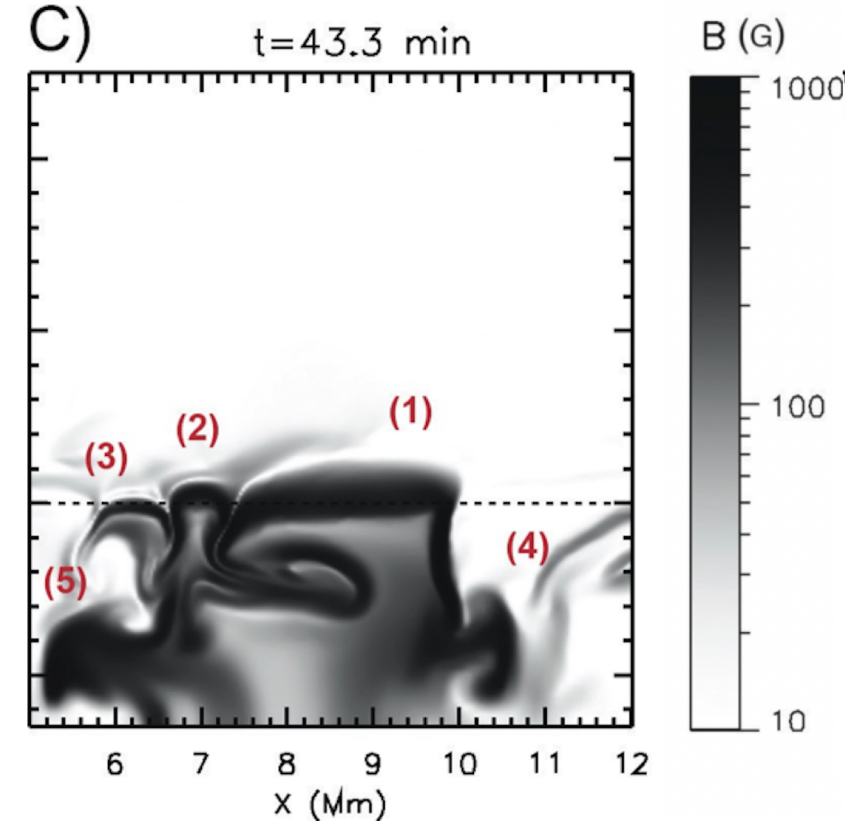
Arrival at the surface:

- Not all the parts of the fragmented tube reach the surface.



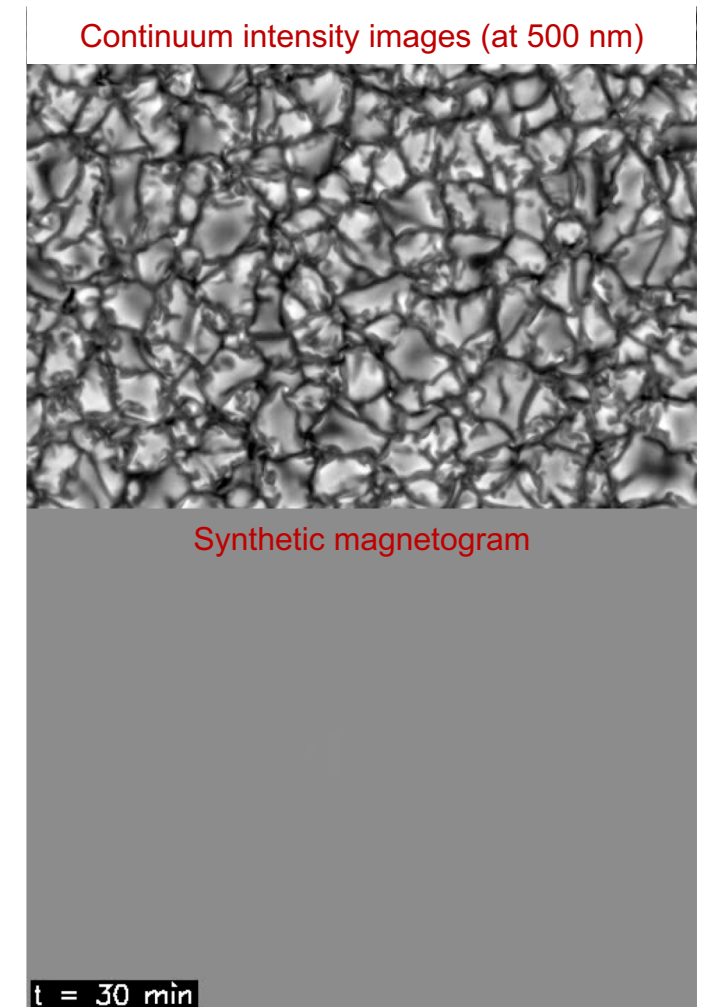
Arrival at the surface:

- In the photosphere, there is transition between super- and sub-adiabatically stratified regions.
- This implies that an emerging fluid element enters a stably stratified region that is difficult to penetrate.
- As a consequence, the magnetized plasma starts to pile up, increasing the magnetic pressure.
- The enhanced pressure produces a sideways growth of the emerging fragments, leading to anomalous granules.



Arrival at the surface:

- In the photosphere, there is transition between super- and sub-adiabatically stratified regions.
- This implies that an emerging fluid element enters a stably stratified region that is difficult to penetrate.
- As a consequence, the magnetized plasma starts to pile up, increasing the magnetic pressure.
- The enhanced pressure produces a sideways growth of the emerging fragments, leading to anomalous granules.



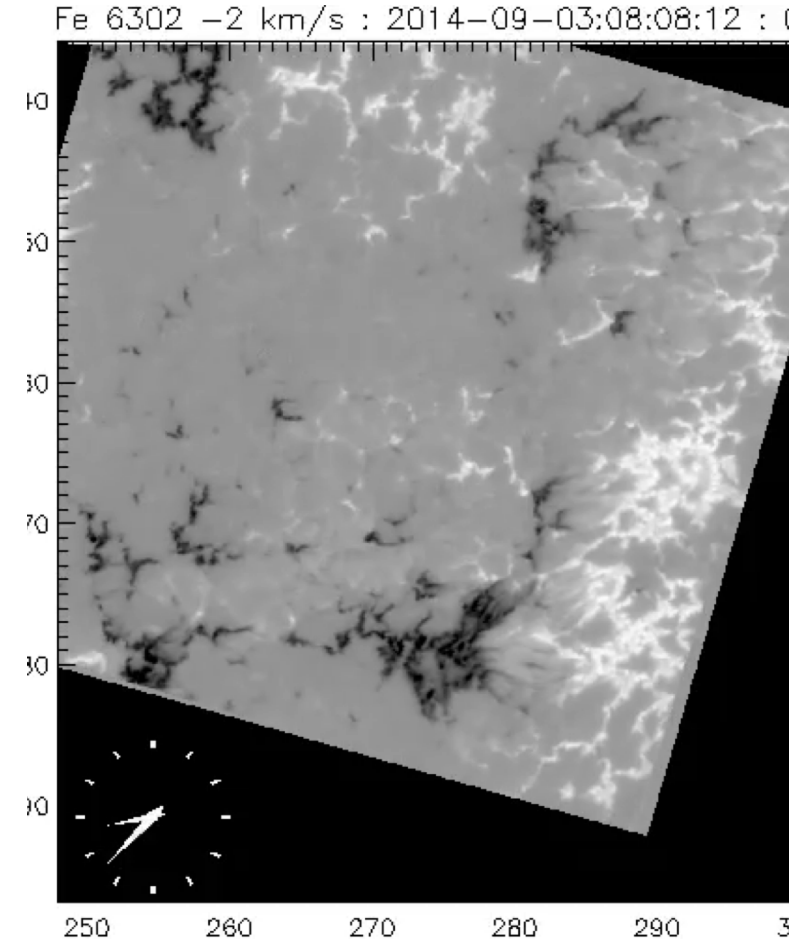
Comparison with observations



Rosseland
Centre
for Solar
Physics

Arrival at the surface:

- In the photosphere, there is transition between super- and sub-adiabatically stratified regions.
- This implies that an emerging fluid element enters a stably stratified region that is difficult to penetrate.
- As a consequence, the magnetized plasma starts to pile up, increasing the magnetic pressure.
- The enhanced pressure produces a sideways growth of the emerging fragments, leading to anomalous granules.



(observational examples of anomalous granulation: Orozco Suárez et al. 2008, Guglielmino et al. 2010, Ortiz et al. 2014, among many others)

SST Observation courtesy of
Luc Rouppe van der Voort

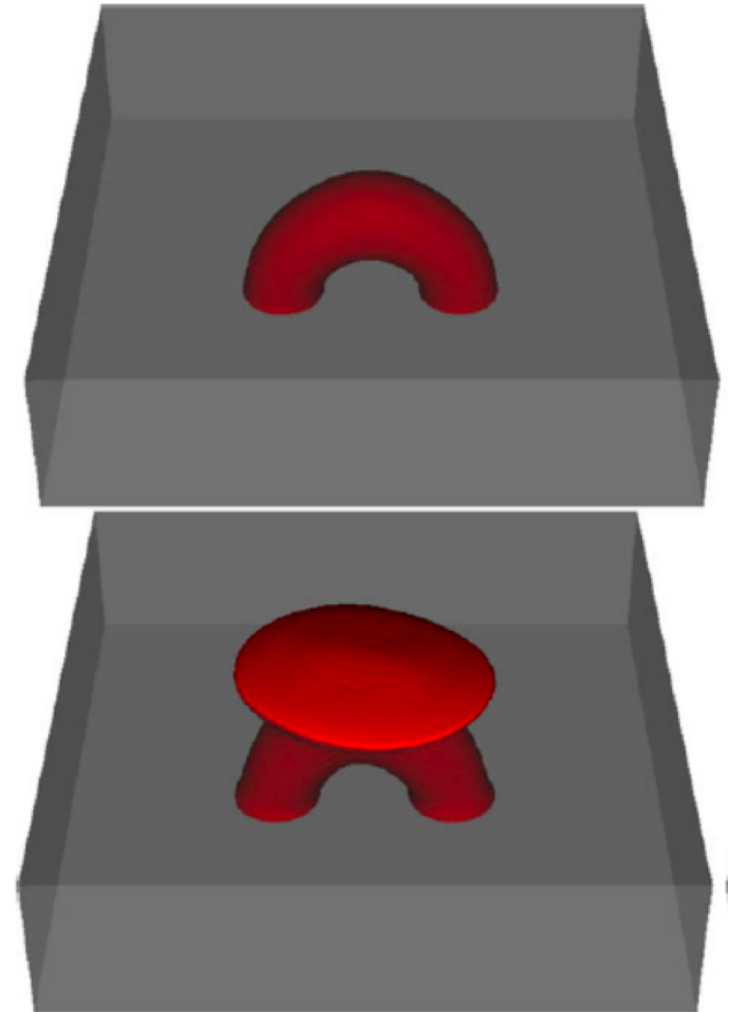
Buoyancy instability:

- The later evolution occurs in the frame of the buoyancy instability:
- This instability occurs when the vertical gradient of the magnetic field is large enough to overcome the sub-adiabatic stabilizing term.

$$-H_p \frac{\partial}{\partial z} (\log B) > -\frac{\gamma}{2} \beta \delta + (k_{\parallel} H_p)^2 \left(1 + \frac{k_{\perp}^2}{k_z^2} \right),$$

- If the field does not get strong enough, no emergence will ensue.

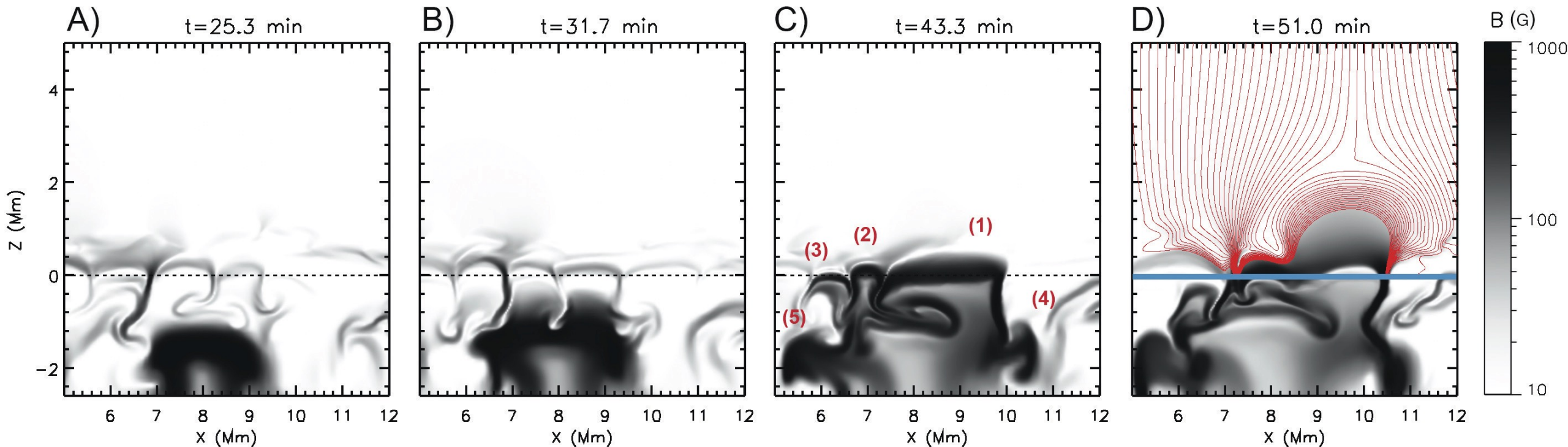
(for details about the buoyancy instability, check Newcomb 1961, Yu 1965, Thomas & Nye 1975, Acheson 1979, among others)



Buoyancy instability:

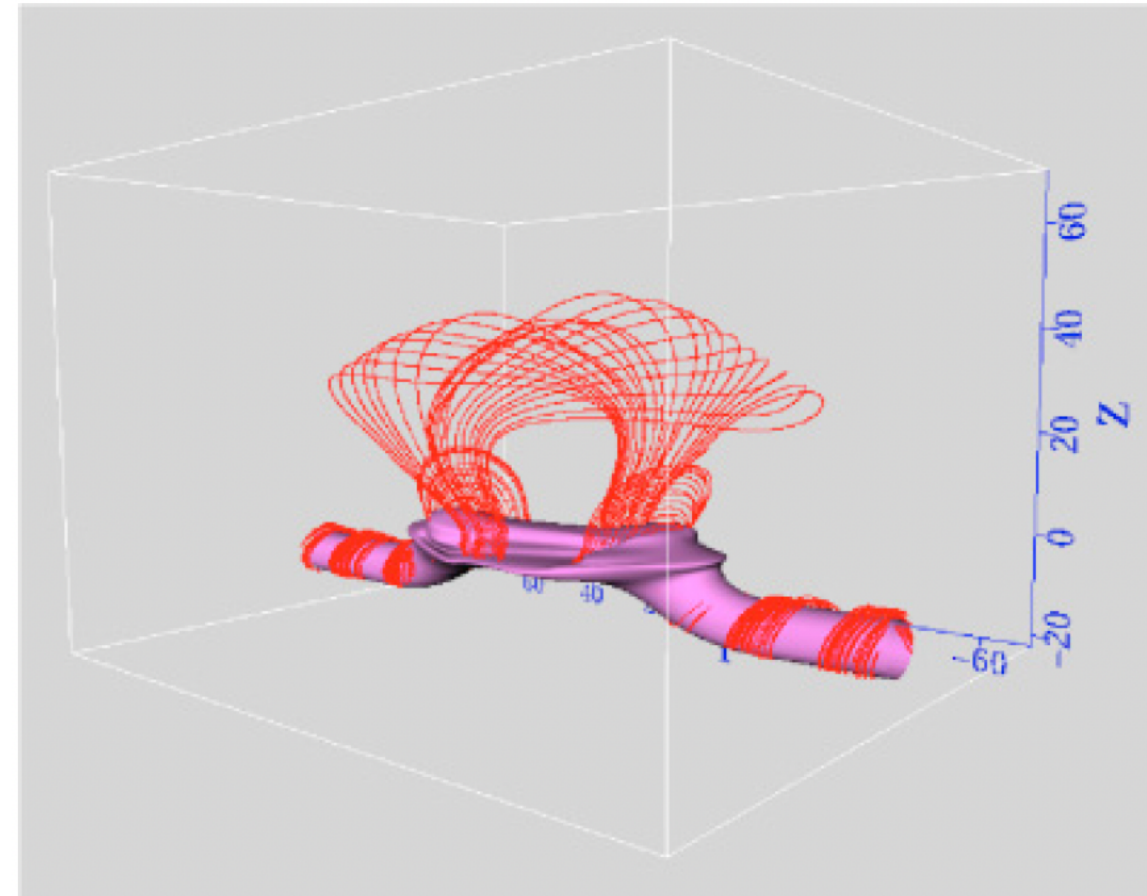
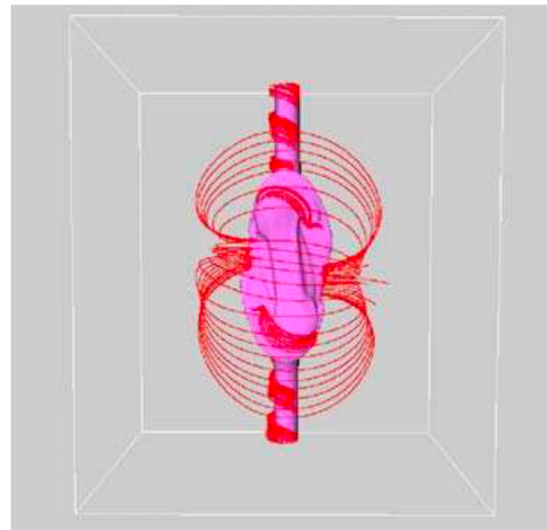
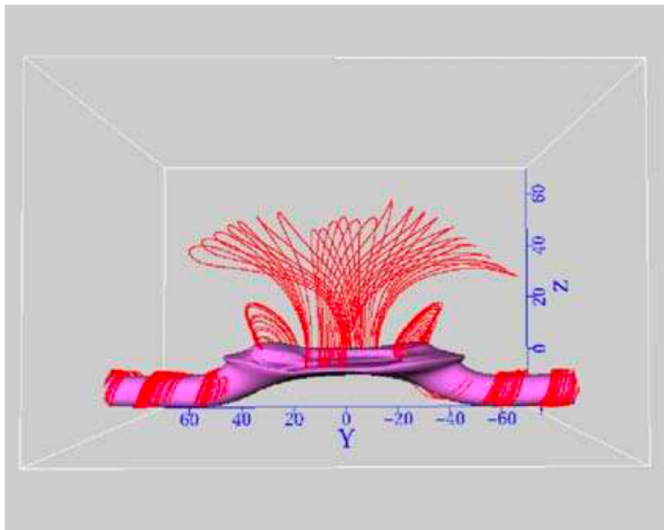
- If the field gets strong enough... this instability will allow the magnetized fluid to go through the solar atmosphere!

Nóbrega-Siverio et al. 2016



Buoyancy instability:

- If the field gets strong enough... this instability will allow the magnetized fluid to go through the solar atmosphere!

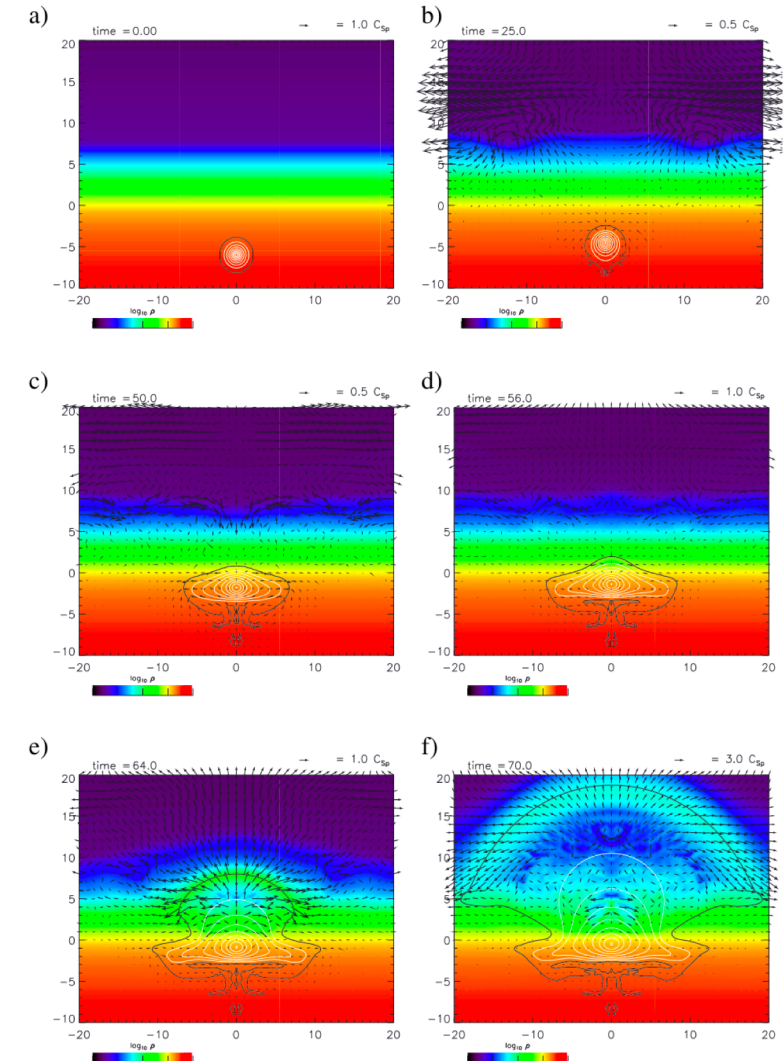


MHD modeling



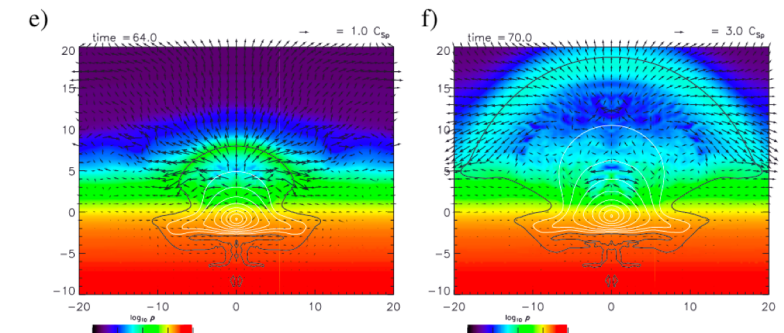
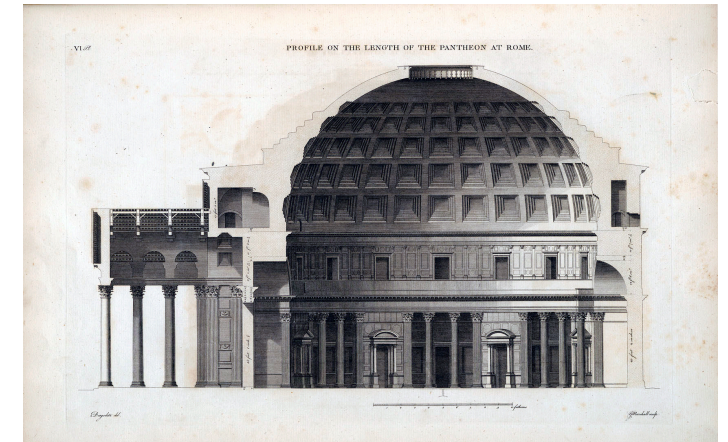
Expansion into the atmosphere:

- The expansion occurs in all directions, but the horizontal components of the velocity are larger than the vertical components.
- This expansion leads to a dome-like structure.
- The expansion is rapid and it decreases local values of temperature and density.
- The dome interior suffers a draining process owing to the gravitational flows that take place along the loop-like magnetic field lines.



Expansion into the atmosphere:

- The expansion occurs in all directions, but the horizontal components of the velocity are larger than the vertical components.
- This expansion leads to a dome-like structure.
- The expansion is rapid and it decreases local values of temperature and density.
- The dome interior suffers a draining process owing to the gravitational flows that take place along the loop-like magnetic field lines.

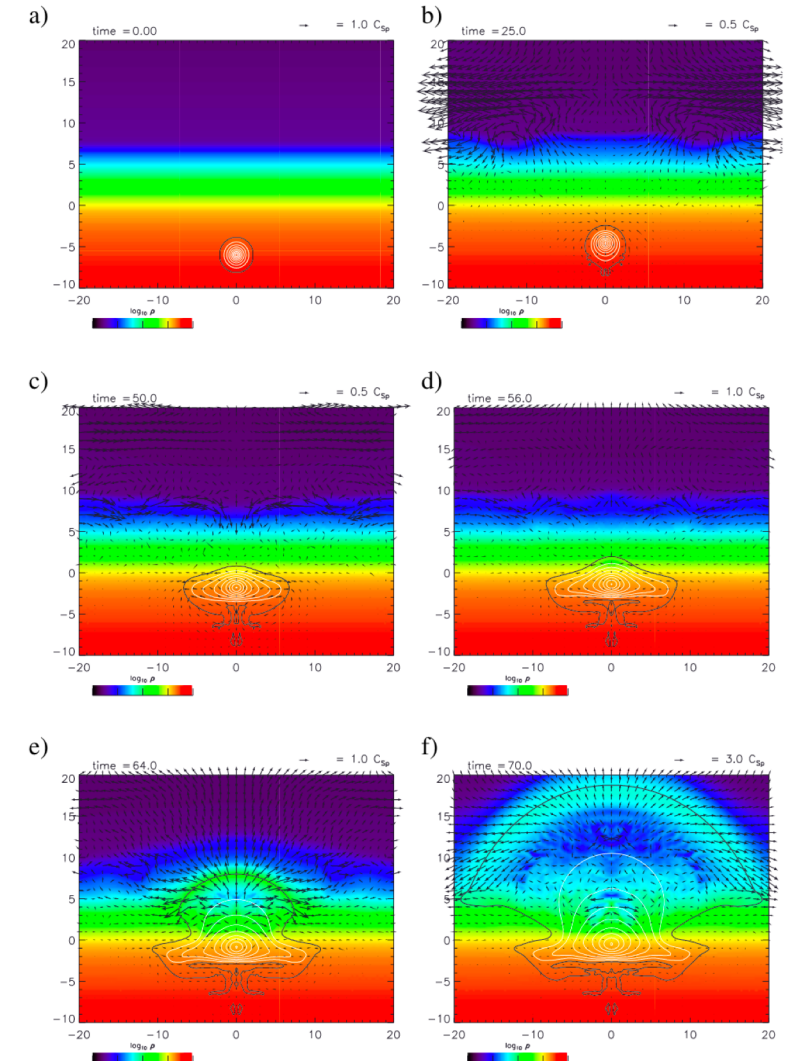


MHD modeling



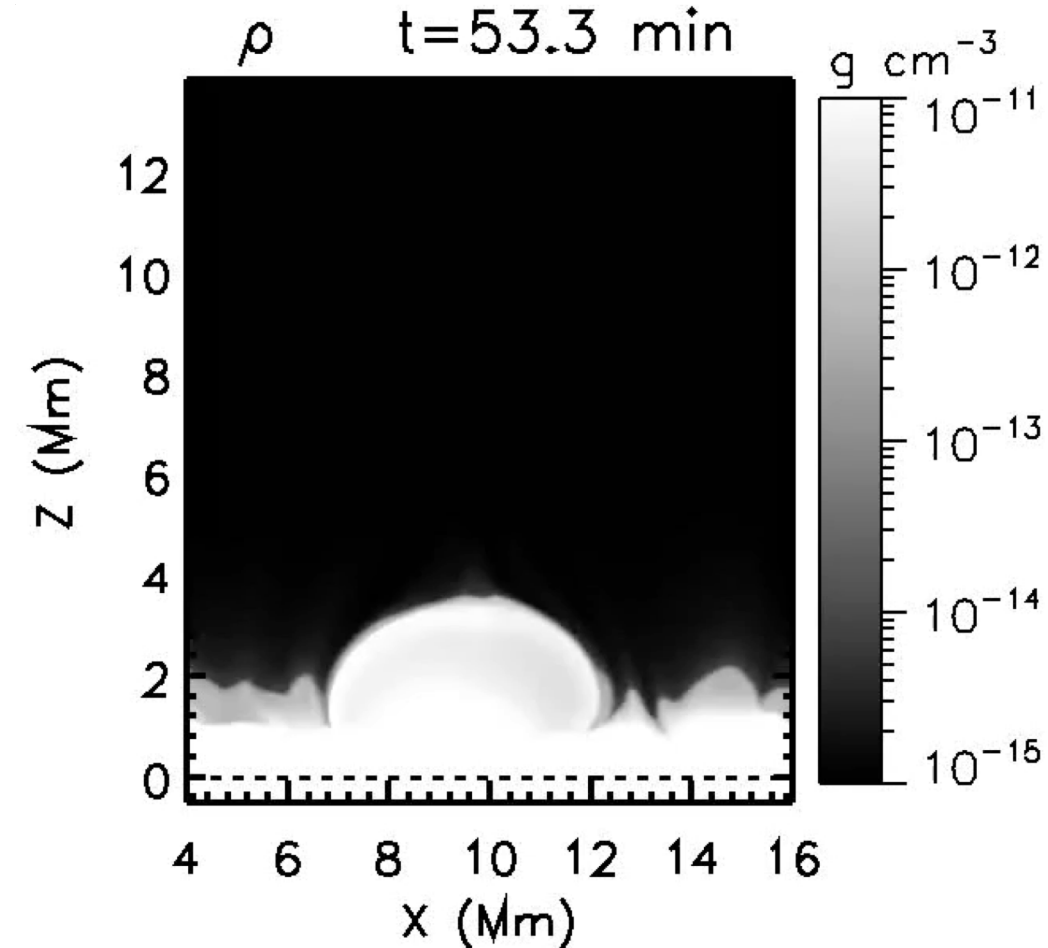
Expansion into the atmosphere:

- The expansion occurs in all directions, but the horizontal components of the velocity are larger than the vertical components.
- This expansion leads to a dome-like structure.
- The expansion is rapid and it decreases local values of temperature and density.
- The dome interior suffers a draining process owing to the gravitational flows that take place along the loop-like magnetic field lines.



Expansion into the atmosphere:

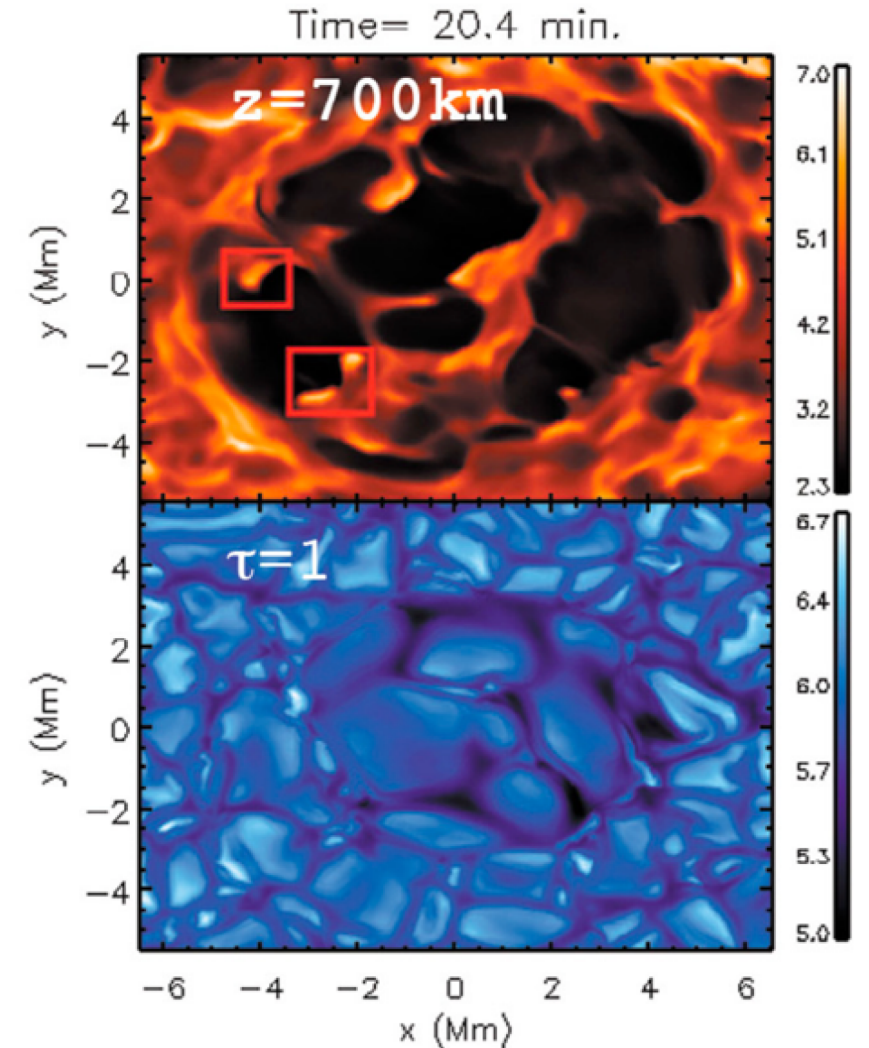
- The expansion occurs in all directions, but the horizontal components of the velocity are larger than the vertical components.
- This expansion leads to a dome-like structure.
- The expansion is rapid and it decreases local values of temperature and density.
- The dome interior suffers a draining process owing to the gravitational flows that take place along the loop-like magnetic field lines.





Expansion into the atmosphere:

- The expansion occurs in all directions, but the horizontal components of the velocity are larger than the vertical components.
- This expansion leads to a dome-like structure.
- The expansion is rapid and it decreases local values of temperature and density.
- The dome interior suffers a draining process owing to the gravitational flows that take place along the loop-like magnetic field lines.



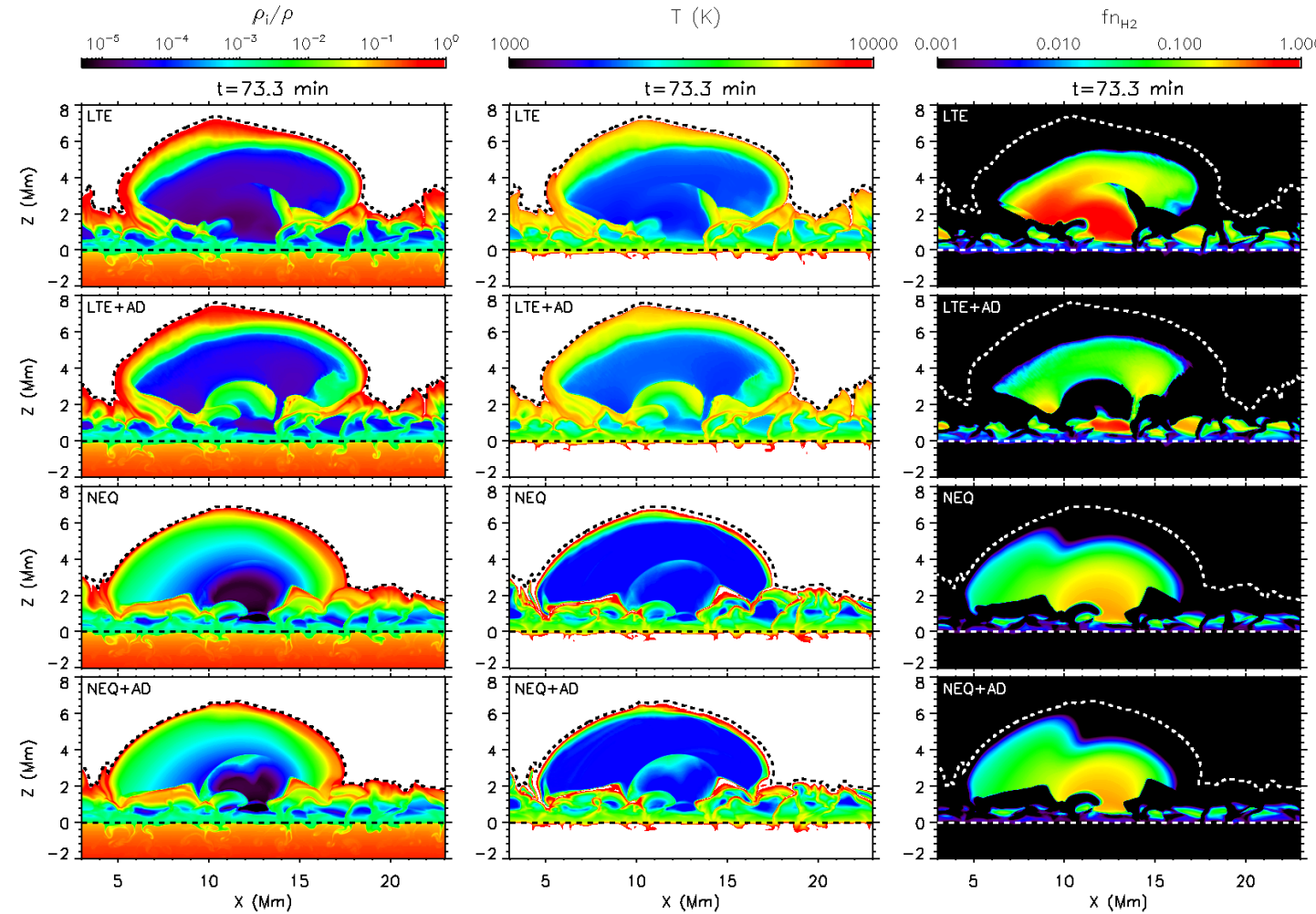
MHD modeling



Expansion into the atmosphere:

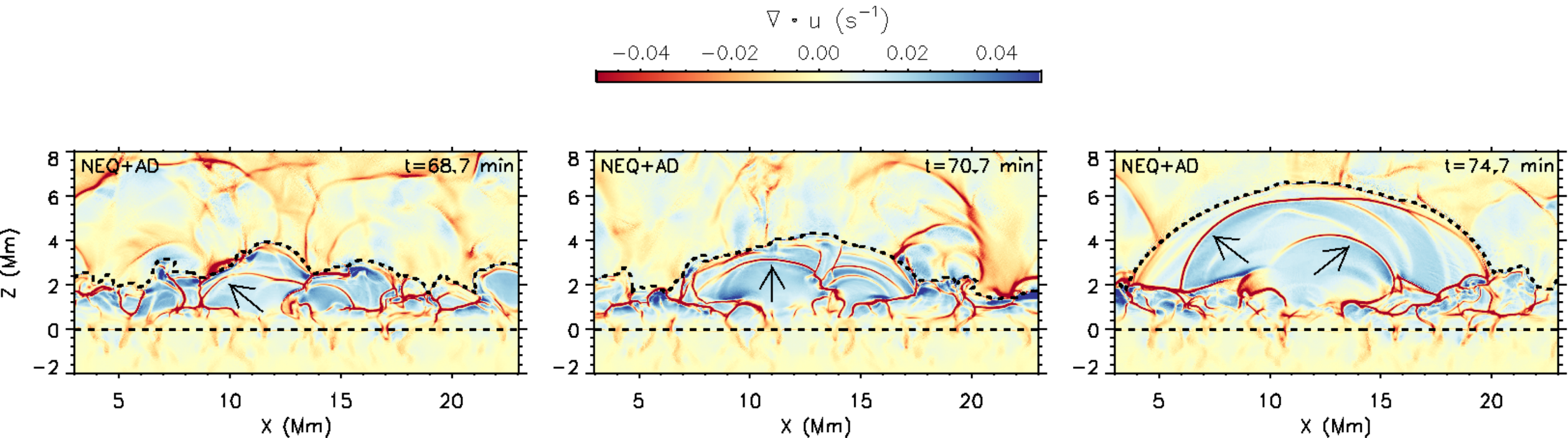
- The thermodynamics of the emerged dome are one of the main open questions.
- Studying the physical properties is difficult both from the theoretical point of view and from observations.

Nóbrega-Siverio et al. 2020



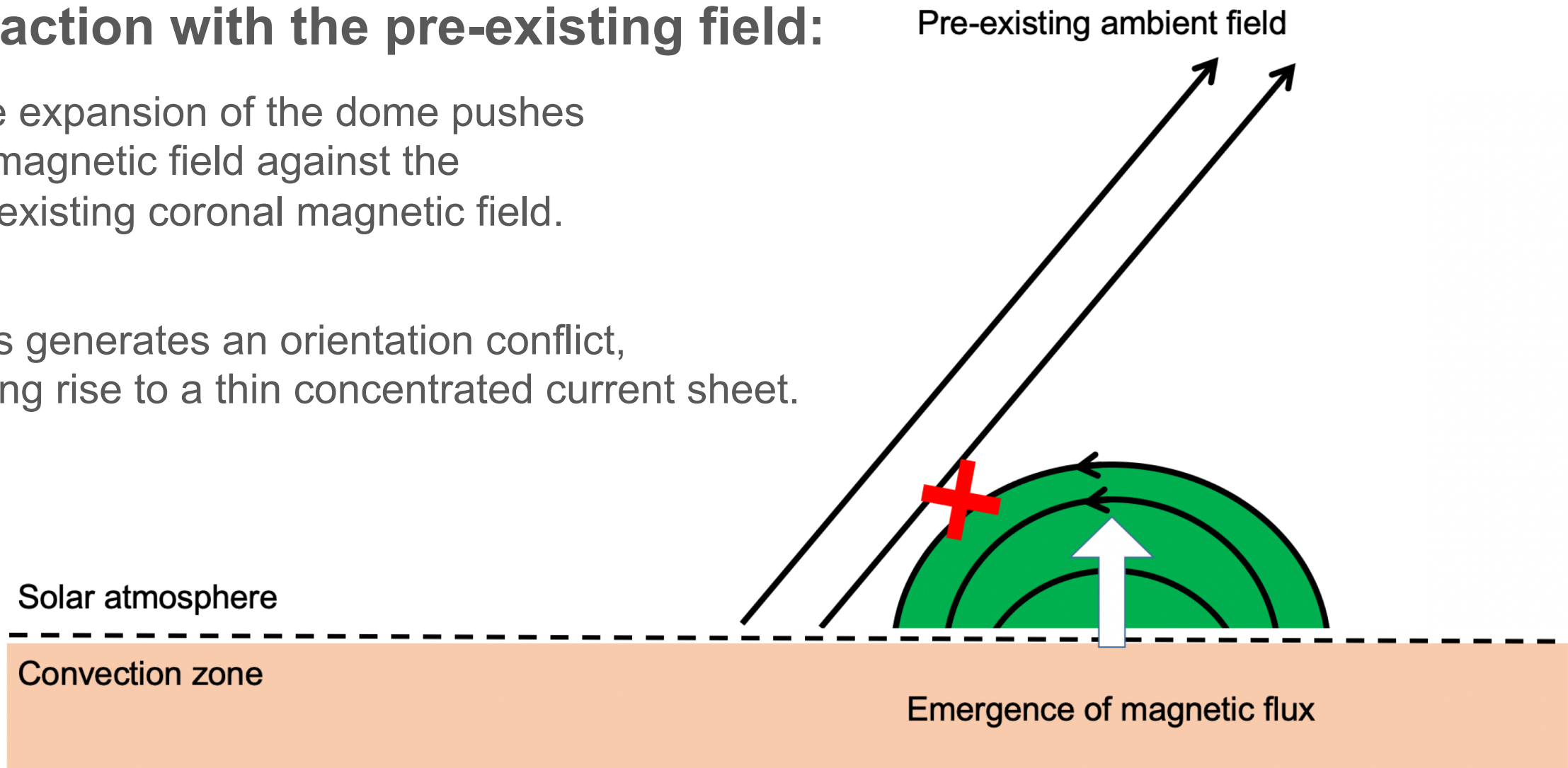
Expansion into the atmosphere:

- Important departures from LTE, and other physical mechanisms such as shocks, molecule formation, ambipolar diffusion... may play an important role.



Interaction with the pre-existing field:

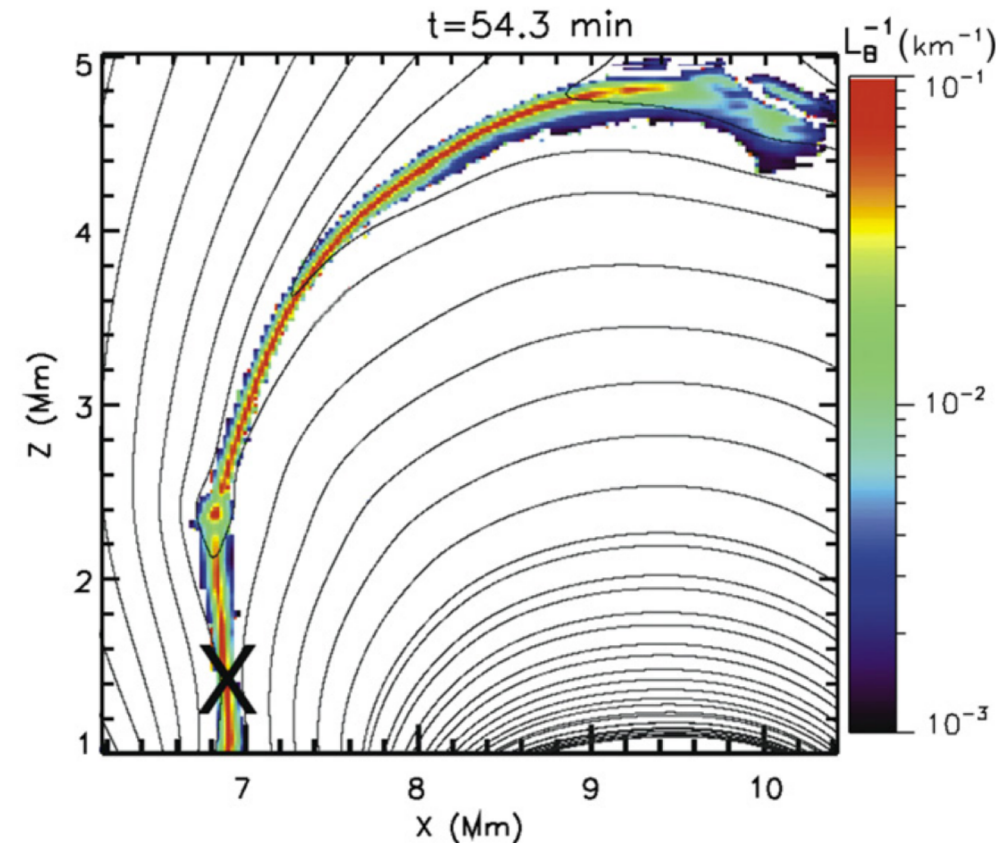
- The expansion of the dome pushes its magnetic field against the preexisting coronal magnetic field.
- This generates an orientation conflict, giving rise to a thin concentrated current sheet.



Interaction with the pre-existing field:

- The expansion of the dome pushes its magnetic field against the preexisting coronal magnetic field.
- This generates an orientation conflict, giving rise to a thin concentrated current sheet.
- This is the ideal scenario for magnetic reconnection.

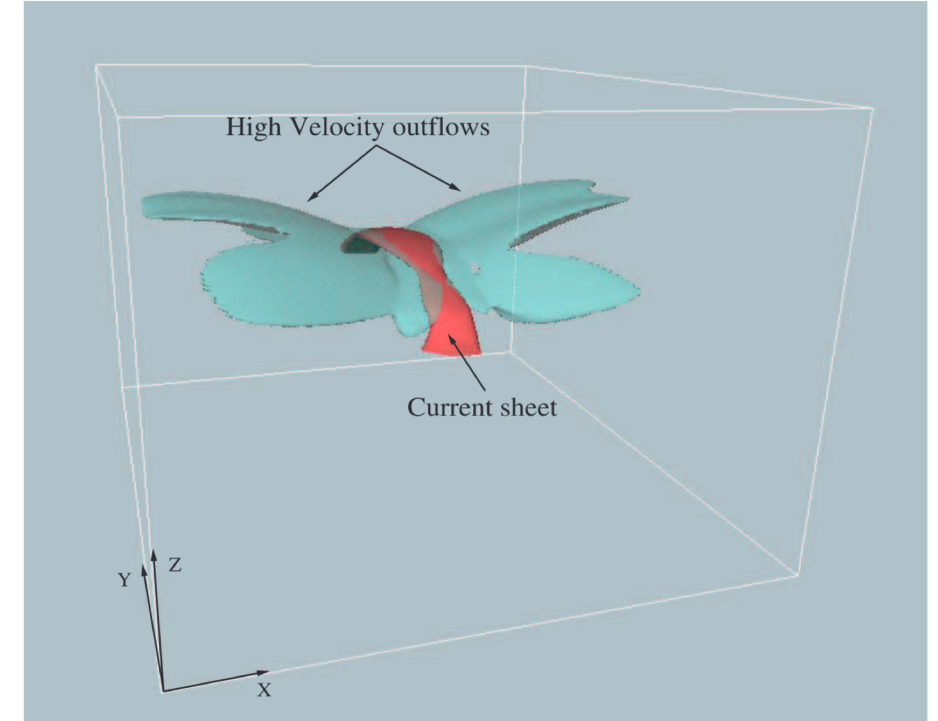
$$L_B^{-1} = \frac{|\nabla \times \mathbf{B}|}{|\mathbf{B}|}$$



Interaction with the pre-existing field:

- The expansion of the dome pushes its magnetic field against the preexisting coronal magnetic field.
- This generates an orientation conflict, giving rise to a thin concentrated current sheet.
- This is the ideal scenario for magnetic reconnection.

$$L_B^{-1} = \frac{|\nabla \times \mathbf{B}|}{|\mathbf{B}|}.$$

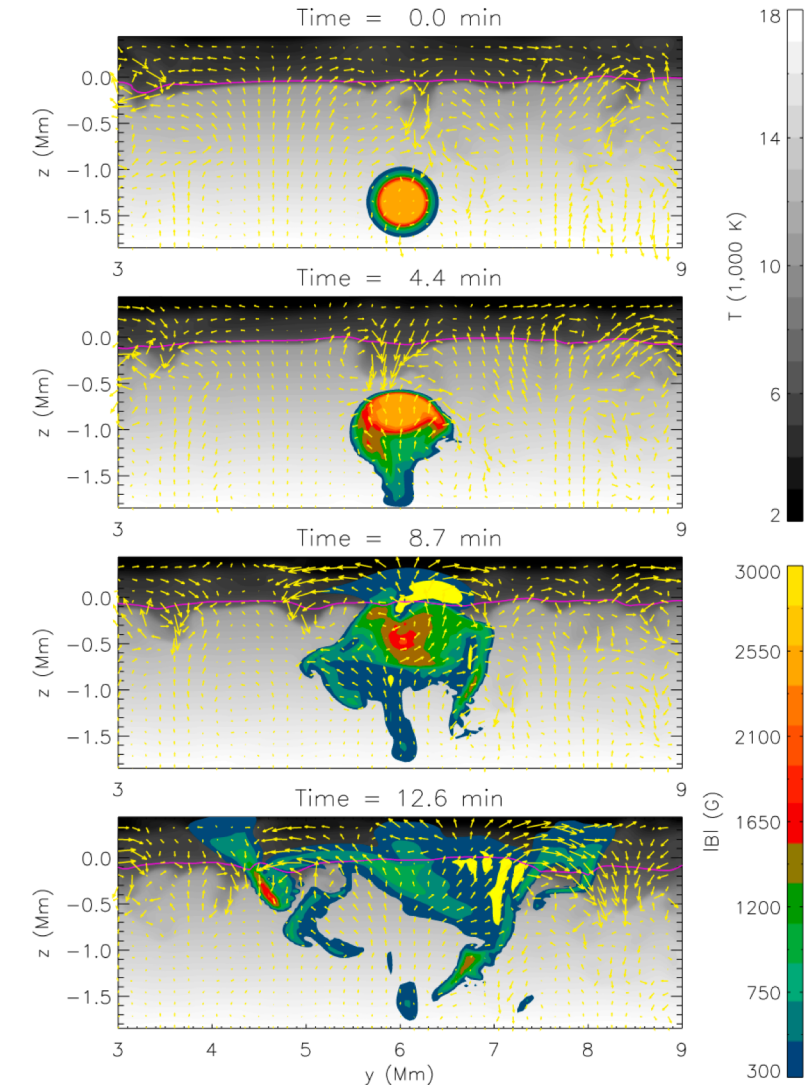


Take away messages



Rise of the tube:

- Because of the stratification, the tube tends to split into counter-rotating vortex rolls and flattens into a pancake-like structure.
- Convection flows can deform and break the tube.
- Strong magnetic field values in the tube diminish the role of the convection flows.
- Higher values of the twist parameter q allow the tube to keep a more coherent structure.

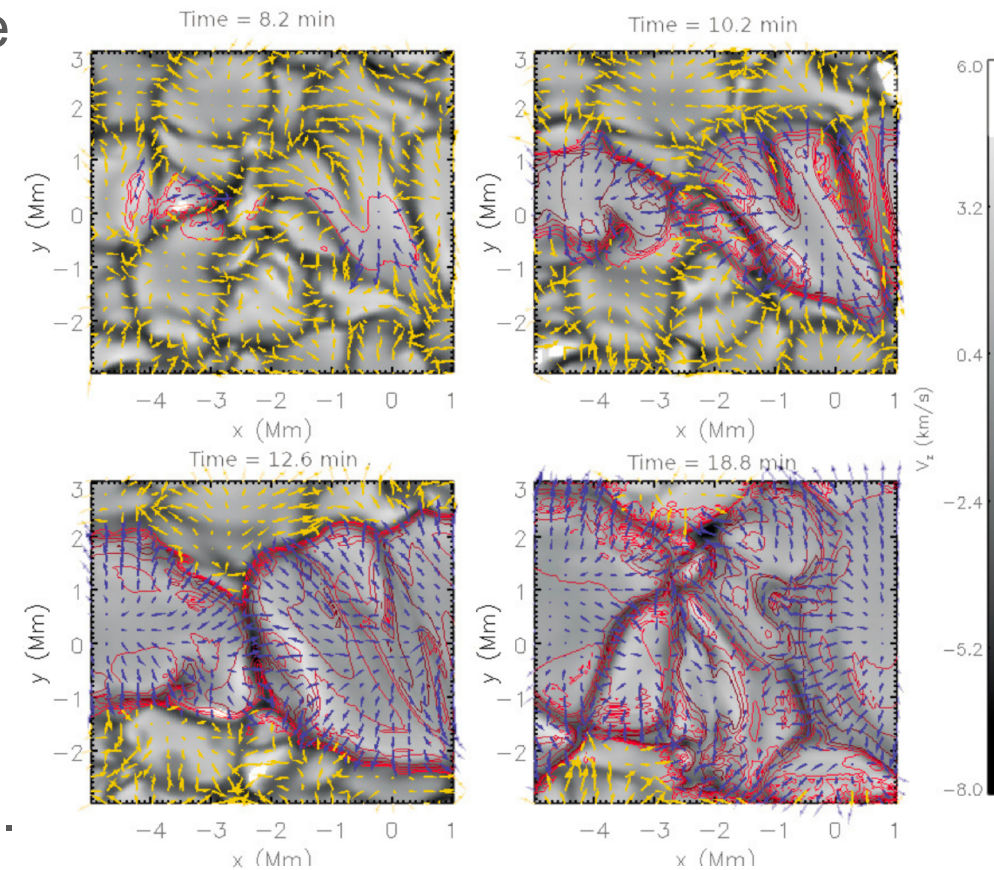


Take away messages



Arrival at the surface:

- Strong subadiabaticity of the photosphere acts like a barrier for passage of flux from the solar interior to the chromosphere and beyond.
- The magnetized plasma piles up and deforms the granulation pattern (anomalous granulations) due to the increasing magnetic pressure.
- When the vertical magnetic field gradient is large enough, buoyancy instability occurs, allowing the magnetized plasma to rise up into the atmosphere.



Take away messages



Expansion into the atmosphere:

- Rapid expansion locally decreases temperature and density, and it leads to the formation of dome-like structures.
- The dome also suffers a rarefaction process due to gravitational downflows
- Because of the expansion, the emerging magnetic field is pushed against the pre-existing coronal ambient field and magnetic reconnection can be triggered.

